# Response of an oscillatory differential delay equation to a single stimulus 

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# Response of an oscillatory differential delay equation to a single stimulus 

Michael C. Mackey ${ }^{1}$ • Marta Tyran-Kamińska ${ }^{2}$. Hans-Otto Walther ${ }^{3}$

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#### Abstract

Here we analytically examine the response of a limit cycle solution to a simple differential delay equation to a single pulse perturbation of the piecewise linear nonlinearity. We construct the unperturbed limit cycle analytically, and are able to completely characterize the perturbed response to a pulse of positive amplitude and duration with onset at different points in the limit cycle. We determine the perturbed minima and maxima and period of the limit cycle and show how the pulse modifies these from the unperturbed case.


Mathematics Subject Classification 92 B $25 \cdot 92$ C $45 \cdot 34$ K $27 \cdot 34$ K 99

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## 1 Introduction

Mammalian hematopoietic systems have complex and complicated regulatory processes that control the production of red blood cells, white blood cells and platelets. However, boiled down to their essence, each is a negative feedback system with a time delay that is controlling the production of primitive cells entering from the hematopoietic stem cell compartment.

Often the numbers of circulating blood cells will display oscillations that are more or less regular. This may occur (Foley and Mackey 2009) because of the existence of a spontaneously occurring disorder like cyclical neutropenia (Dale and Hammond 1988; Haurie et al. 1998; Colijn et al. 2006; Dale and Mackey 2015), periodic thrombocytopenia (Apostu and Mackey 2008; Swinburne and Mackey 2000), periodic leukemia (Colijn and Mackey 2005; Fortin and Mackey 1999), or periodic autoimmune hemolytic anemia (Mackey 1979; Milton and Mackey 1989). Or, it may occur because of the intrusive administration of chemotherapy in a periodic fashion (Krinner et al. 2013) which has the unfortunate side effect of killing both malignant and normal cells.

In either case (spontaneously occurring oscillations due to disease or induced oscillations due to the side effects of chemotherapy) a clinical intervention often consists of trying to administer a recombinant cytokine of the appropriate type to alleviate the more serious symptoms of the oscillation. In the case of cyclical neutropenia this is granulocyte colony stimulating factor (G-CSF) (Dale and Hammond 1988; Dale et al. 1993, 2003), and the same is true during chemotherapy induced neutropenia (Bennett et al. 1999; Clark et al. 2005). G-CSF has among its effects the ability to interfere with apoptosis (pre-programmed death) of cells (Colijn et al. 2007), be this cell death naturally present or induced by an external agent like chemotherapy. Unfortunately the issue of when to administer this (or other) cytokines is hotly debated and this is, without a doubt, because the cytokines in question have an effect on the dynamics of the affected system many days before the desired (or undesired) effect is manifested in the peripheral blood.

It has been noted that the timing of the administration of G-CSF can have profound consequences on the neutropenia. Given at some points in the cycle it can dramatically
reduce the neutropenia (increasing the nadir of the cycle) while at other times it can actually make the neutropenia worse by deepening the nadir (Aapro et al. 2011; Barni et al. 2014; Munoz Langa et al. 2012; Palumbo et al. 2012). Thus, from a mathematical perspective the problem is simply "How and when do we deliver a perturbation to a delayed dynamical system in order to achieve some desired objective?"

The problem outlined above can, from a mathematical point of view, be viewed within the context of 'phase resetting of an oscillator' and as such has received widespread attention especially within the biological community. This field has a large and varied literature (see Glass and Mackey 1988 for an elementary introduction and Winfree 1980 for an exhaustive treatment of the subject from a historical perspective) which is almost exclusively devoted to the interaction of oscillatory systems in a finite dimensional space (i.e. limit cycles in ordinary differential equations) with a single perturbation or periodic perturbation. Surprisingly, however, there is little that has been done on such interactions when the limit cycle is in an infinite dimensional phase space (e.g a differential delay equation). There are, however, a few authors who have considered such situations. For example, Bodnar et al. (2013a, c), Piotrowska and Bodnar (2014) and Foryś et al. (2014) studied simple models of tumour growth where the delayed model equation has an additional term describing an external influence and reflecting a treatment. There have been a number of both experimental and theoretical papers (Israelsson and Johnsson 1967; Johnsson and Israelsson 1968; Johnsson 1971; Andersen and Johnsson 1972a, b) devoted to the autonomous growth of the tip of Helianthus annuus which describes a variety of patterns as a function of time and which is thought to involve a delay between the sensing of a gravitational stimulus and the bending of the plant (c.f Israelsson and Johnsson 1967 for a very nice historical review of this problem). Another class of problems involving delayed dynamics is related to pulse coupled oscillators which have been treated recently by Canavier and Achuthan (2010), Klinshov and Nekorkin (2011) and Klinshov et al. (2015). Kotani et al. (2012) and Novicenko and Pyragas (2012) have developed phase reduction methods appropriate for delayed dynamics. Finally we should note the recent numerical work on several gene regulatory circuits by Lewis (2003) and Horikawa et al. (2006) for the segmentation clock in zebrafish as well as the work of Doi et al. (2011) on circadian regulation of G-protein signaling. However, none of these papers have addressed the problem that we study here from an analytic point of view. This paper offers a partial study of the problem.

The regulation of the production of blood cells, denoted by $x(t)$ (and typically measured in units of cells $/ \mu \mathrm{L}$ of blood or alternately in units of cells $/ \mathrm{kG}$ body weight), reduced to the barest of descriptions, can be described most simply by a differential delay equation of the form

$$
\begin{equation*}
\frac{d x}{d t}(t)=-\gamma x(t)+f(x(t-\tau)) \quad \text { with constant } \gamma>0 \tag{1.1}
\end{equation*}
$$

in which $f: \mathbb{R} \rightarrow \mathbb{R}$ is monotone decreasing such that $\xi_{1} \leq \xi_{2}$ implies that $f\left(\xi_{1}\right) \geq$ $f\left(\xi_{2}\right)$. In Eq. (1.1) we must also specify an initial function $\varphi:[-\tau, 0] \rightarrow \mathbb{R}$, in order to obtain a solution. Here we replace the nonlinearity $f$ with a piecewise constant function. This permits us to compute solutions explicitly, so we may analytically study
their behaviour, and the response of the solutions to perturbation meant to represent the effect of cytokine administration.

This generic model captures the essence, if not the subtleties, of peripheral blood production. The monotone nature of $f$ is mediated via the effects of the important regulatory cytokines, e.g. G-CSF for the white blood cells, erythropoietin for the red blood cells, and thrombopoietin for the platelets. The administration of exogenous cytokines in an attempt to control the dynamics of (1.1) will typically have an effect that be interpreted as increasing $f$ over some portion of time, and the goal of this paper is to study the effect of such a perturbation on the solution of (1.1).

This paper is organized as follows. Section 2 describes the model and formulates it in a mathematically convenient form. Section 3 provides basic facts about continuous, piecewise differentiable solutions. On a state space of simple initial functions these solutions yield a continuous semiflow. There is a periodic solution whose orbit in state space is stable with strong attraction properties. Section 4 introduces the pulse-like perturbations (a perturbation of constant amplitude $a$ lasting for a finite period of time $\sigma$ ) of the model which correspond to the effect of medication in the sense that during a finite time interval the production of blood cells is increased. It is shown that the response of the system to such perturbations is continuous provided the latter are not too large. This includes a continuity result for the cycle length map, which assigns to each onset time of increased production a time of return to the periodic orbit, after the end of increased production. The bulk of our results are presented in Sect. 5 where we examine the effect of cytokine perturbation when the perturbation away from the stable periodic orbit begins at different points in the cycle. In particular we look at phase resetting properties of the system-in terms of the cycle length map-and at the minima and maxima compared to the amplitudes of the periodic solution. Section 6 examines the various forms assumed by the cycle length map for different values of the parameters. Section 7 shows how the results of the previous sections may be potentially used to tailor therapy to achieve certain results. The paper concludes with a brief discussion in Sect. 8. There we consider a simple extension in which a pulse-like perturbation may decrease the nadir of the limit cycle as is noted clinically. The proofs of many of the results are given in the two appendices.

## 2 The model

### 2.1 Scalar delay differential equations

Consider the delay differential equation (1.1). If $f$ is continuous and monotone decreasing then there is a unique constant solution $t \mapsto x_{*}$ given by $f\left(x_{*}\right)=\gamma x_{*}$. If in addition $f$ is, say, continuously differentiable then this constant solution may be stable or not, depending on $\gamma$ and $f^{\prime}\left(x_{*}\right)$. In case $t \mapsto x_{*}$ is linearly unstable, that is, the linearized equation

$$
y^{\prime}(t)=-\gamma y(t)+f^{\prime}\left(x_{*}\right) y(t-\tau)
$$

is unstable then also $t \mapsto x_{*}$ is unstable as a solution of (1.1). In case $f$ also satisfies a one-sided boundedness condition there exists a periodic solution which is slowly oscillating in the sense that the intersections with the equilibrium level $\xi=x_{*}$ are spaced at distances larger than the delay $\tau$, and the minimal period is given by three consecutive such intersections. In general slowly oscillating periodic solutions are not unique in the sense that they are not all translates of each other. Depending on $\gamma$ and $f^{\prime}\left(x_{*}\right)$ there may also exist rapidly oscillating periodic solutions about $\xi=x_{*}$. These have all consecutive zeros spaced at distances strictly less than the delay $\tau$. In case $f^{\prime}(\xi)<0$ for all $\xi \in \mathbb{R}$ every rapidly oscillating periodic solution is unstable. For details and for more about (1.1), see e.g. the recent survey by Walther (2014) and the references given there.

### 2.2 A piecewise constant approximation of the nonlinearity

To obtain solutions which can be computed in terms of elementary functions consider the situation in which the function $f$ is piecewise constant and given by

$$
f(x)=\left\{\begin{array}{ll}
b_{L} & \text { for } \quad x<\theta  \tag{2.1}\\
b_{U} & \text { for } \quad x \geq \theta,
\end{array} \quad 0<\theta, \quad 0<b_{U}<b_{L}\right.
$$

We exclude the special cases $b_{L}=\gamma \theta$ and $b_{U}=\gamma \theta$, in order to facilitate the construction of the solution semi-flow in Sect. 3 below, see the skewed dashed lines in Fig. 1 and Remark 3.1.

Then we obtain the delay differential equation

$$
\frac{d x}{d t}(t)=-\gamma x(t)+ \begin{cases}b_{L} & \text { for } \quad x(t-\tau)<\theta  \tag{2.2}\\ b_{U} & \text { for } \quad x(t-\tau) \geq \theta\end{cases}
$$

with $\gamma>0, \theta>0, b_{L}>b_{U}>0, \gamma \theta \neq b_{U}, b_{L}$, whose solutions satisfy linear, inhomogeneous ordinary differential equations in intervals on which their delayed values remain either below or above the level $\xi=\theta$, resulting in increasing and decreasing exponentials on such intervals. In the situation illustrated in Fig. 1 when the graph of $x \mapsto \gamma x$ passes through the gap of the nonlinearity in Eq. (2.2), so that $b_{U}<\gamma \theta<b_{L}$, there is no steady state of Eq. (2.2). Otherwise, if either $\gamma \theta<b_{U}$ or $b_{L}<\gamma \theta$ then $x(t)=b_{U} / \gamma$ or $x(t)=b_{L} / \gamma$ is the steady state of Eq. (2.2).

Fig. 1 The graph of the function $f$ as given in (2.1). The skewed solid line is the graph of $\gamma x$ in general, while the skewed dashed lines are the graphs of $\gamma x$ in the special excluded cases that $b_{L}=\gamma \theta$ and $b_{U}=\gamma \theta$


Note that in (2.2) there are five parameters $\left(\gamma, b_{U}, b_{L}, \theta, \tau\right)$ and we can reduce these to three by a change of variables, since with

$$
\left\{\begin{array}{l}
\hat{x}(t)=x(t / \gamma)-\theta  \tag{2.3}\\
\hat{\tau}=\gamma \tau \\
\beta_{L}=-\theta+b_{L} / \gamma \\
\beta_{U}=\theta-b_{U} / \gamma
\end{array}\right.
$$

we can rewrite (2.2) in the form

$$
\frac{d \hat{x}}{d t}(t)= \begin{cases}-\hat{x}(t)+\beta_{L} & \text { for } \quad \hat{x}(t-\hat{\tau})<0  \tag{2.4}\\ -\hat{x}(t)-\beta_{U} & \text { for } \quad \hat{x}(t-\hat{\tau}) \geq 0\end{cases}
$$

where $\beta_{L}+\beta_{U}>0$.
Now we have only a three parameter $\left(\beta_{U}, \beta_{L}, \tau\right)$ system to consider, having reduced (2.2) to

$$
\begin{equation*}
\frac{d x}{d t}(t)=-x(t)+f(x(t-\tau)) \tag{2.5}
\end{equation*}
$$

where the function $f$ is of the form

$$
f(\xi)=\left\{\begin{array}{ll}
\beta_{L} & \text { for } \xi<0  \tag{2.6}\\
-\beta_{U} & \text { for } \xi \geq 0,
\end{array} \quad-\beta_{U}<\beta_{L}, \quad \beta_{L}, \beta_{U} \neq 0 .\right.
$$

The discontinuity of $f$ requires a moment of reflection about what a solution of Eq. (2.5) should be-certainly not a continuous function $x:[-\tau, \infty) \rightarrow \mathbb{R}$ which is differentiable and satisfies Eq. (2.5) for all $t>0$, as it is familiar from delay differential equations with a functional on the right hand side which is at least continuous. We shall come back to this in Sect. 3.

## 3 The semiflow of the unperturbed system

For continuity properties, e.g., continuity of the reset map which is to be introduced in Sect. 4, we need to develop a framework.

Consider the initial value problem of Eq. (2.5) for $t>0$, with initial condition $x(t)=\phi(t)$ for $-\tau \leq t \leq 0$, and where the function $\phi:[-\tau, 0] \rightarrow \mathbb{R}$ is continuous and has at most a finite number of zeros. Let $C=C([-\tau, 0], \mathbb{R})$ denote the Banach space of continuous functions $[-\tau, 0] \rightarrow \mathbb{R}$ equipped with the maximumnorm, $|\phi|_{C}=\max _{-\tau \leq t \leq 0}|\phi(t)|$, and set

$$
Z=\left\{\phi \in C: \phi^{-1}(0) \text { finite or empty }\right\} .
$$

For initial data $\phi \in Z$ we construct continuous solutions of Eq. (2.5) by means of the variation-of-constants formula as follows. Suppose $z_{1}<z_{2}<\cdots<z_{J}$ are the zeros of $\phi$ in $(-\tau, 0)$. On $\left(0, z_{1}+\tau\right]$ we define

$$
x(t)=e^{-t} \phi(0)+\int_{0}^{t} e^{-(t-s)} f(x(t-s)) d s
$$

or,

$$
\begin{equation*}
x(t)=e^{-t} x(0)-\beta_{U}\left(1-e^{-t}\right)=-\beta_{U}+\left(x(0)+\beta_{U}\right) e^{-t} \tag{3.1}
\end{equation*}
$$

in case $0<\phi(v)$ in $\left(-\tau, z_{1}\right)$,

$$
\begin{equation*}
x(t)=e^{-t} x(0)+\beta_{L}\left(1-e^{-t}\right)=\beta_{L}+\left(x(0)-\beta_{L}\right) e^{-t} \tag{3.2}
\end{equation*}
$$

in case $\phi(v)<0$ in $\left(-\tau, z_{1}\right)$. Notice that on the interval $\left[0, z_{1}+\tau\right]$ the solution $x$ is either constant with value $-\beta_{U} \neq 0$ or $\beta_{L} \neq 0$, or strictly monotone. We conclude that $x$ has at most one zero in $\left(0, z_{1}+\tau\right)$, and if there is a zero then $x$ changes sign at the zero. Let us call such zeros transversal. In case $\phi$ has no zero in $(-\tau, 0)$ we define $x(t)$ analogously, for $0<t \leq \tau$. The procedure just described can be iterated and yields a continuous function $x:[-\tau, \infty) \rightarrow \mathbb{R}$ which we define to be the solution of the initial value problem above.

Notice that all segments, or histories, $x_{t}$ given by

$$
x_{t}(s)=x(t+s) \text { for } t \geq 0 \text { and }-\tau \leq s \leq 0,
$$

belong to the set $Z$. We assume that $\beta_{L}, \beta_{U} \neq 0$ and we write $x^{\phi}$ instead of $x$ when convenient, and define the semiflow $S:[0, \infty) \times Z \rightarrow Z$ of Eq. (2.5) by

$$
S(t, \phi)=x_{t}^{\phi}
$$

The proof of the following result is given in the appendix.
Proposition 3.1 The semiflow $S$ is continuous.
Remark 3.1 Incidentally, let us see what goes wrong in the excluded cases $\beta_{U}=0$ and $\beta_{L}=0$. If $\beta_{U}=0$ then for each $\phi \in Z$ with $\phi(t)>0$ on $[-\tau, 0)$ and $\phi(0)=0$ the formula (3.1)—or, the equation $x^{\prime}=-x$-yields $x(t)=0$ for all $t \geq 0$, and $Z$ is not positively invariant. On a set $\tilde{Z} \subset C$ of initial data which contains $0 \in C$ and negative data $\psi$ with arbitrarily small norm continuous dependence on initial data would be violated for $\beta_{U}=0$ because (3.2) yields

$$
\begin{aligned}
|x(\tau)-0| & =|x(\tau)|=\left|\beta_{L}+\left(x(0)-\beta_{L}\right) e^{-\tau}\right| \\
& \geq \beta_{L}\left(1-e^{-\tau}\right)-\psi(0) e^{-\tau} \geq \frac{1}{2} \beta_{L}\left(1-e^{-\tau}\right)
\end{aligned}
$$

for the solution starting from negative and sufficiently small $\psi \in \tilde{Z}$.
If $\beta_{L}=0$ then for each $\phi \in Z$ with $\phi(t)<0$ on $[-\tau, 0)$ and $\phi(0)=0$ we have $x(t)=0$ for $t \in[0, \tau]$, which implies that $Z$ is not positively invariant in this case as well.

For later use we show next that transversal zeros depend continuously on the initial data $\phi \in Z$.

Proposition 3.2 For $\phi \in Z$ and $z>0$ with $x^{\phi}(z)=0 \neq x^{\phi}(z-\tau)$ and $\epsilon>0$, there exists $\delta>0$ such that for each $\psi \in Z$ with $|\psi-\phi|_{C}<\delta$ there is $z^{\prime} \in(z-\epsilon, z+\epsilon)$ with $x^{\psi}\left(z^{\prime}\right)=0$. Moreover, $x^{\psi}\left(z^{\prime}-\tau\right) \neq 0$.

Proof By continuity there exists $\eta \in(0, \epsilon)$ so that $x^{\phi}(s) \neq 0$ on $[z-\tau-\eta, z-\tau+$ $\eta] \cap[-\tau, \infty)$. Using (3.1) and (3.2) we infer that on $[z-\eta, z+\eta]$ the solution $x^{\phi}$ either equals a nonzero constant or is strictly monotone. As $x^{\phi}(z)=0$ the solution $x^{\phi}$ must be strictly monotone on $[z-\eta, z+\eta]$, with $\operatorname{sign}\left(x^{\phi}(z-\eta)\right) \neq \operatorname{sign}\left(x^{\phi}(z+\eta)\right) \neq 0$. By continuous dependence on initial data, there exists $\delta>0$ such that for each $\psi \in Z$ with $|\psi-\phi|_{C}<\delta$ we have

$$
x^{\psi}(t) \neq 0 \quad \text { on }[z-\tau-\eta, z-\tau+\eta] \cap[-\tau, \infty)
$$

and

$$
\operatorname{sign}\left(x^{\psi}(z-\eta)\right)=\operatorname{sign}\left(x^{\phi}(z-\eta)\right), \quad \operatorname{sign}\left(x^{\psi}(z+\eta)\right)=\operatorname{sign}\left(x^{\phi}(z+\eta)\right) .
$$

Hence $x^{\psi}$ changes sign in $[z-\eta, z+\eta]$, so it has a zero $z^{\prime}$ in this interval.
The condition $\beta_{L}<0\left(\beta_{U}<0\right)$ in the next result means that in the original model given by Eq. (2.2) the positive constant solution given by $\gamma x^{*}=f\left(x^{*}\right)$ has its value $x^{*}$ beyond (above) the discontinuity $\theta$ of $f$. In this case one may interpret $x^{*}=\theta$ as an equilibrium position for (2.2).

Theorem 3.1 If $\beta_{U}<0$ or $\beta_{L}<0$ the equilibrium state of the semi-flow $S$, which is respectively given by $x^{\phi}(t)=-\beta_{U}$ or $x^{\phi}(t)=\beta_{L}$ with $\phi \in Z$, is globally asymptotically stable.

Proof In case $\beta_{U}<0$ the constant function $\mathbb{R} \ni t \mapsto-\beta_{U} \in \mathbb{R}$ is a positive solution. Notice that due to (3.1) every solution $x^{\phi}$ with $0<\phi(t)$ on $[-\tau, 0]$ has its values $x^{\phi}(t), t \geq 0$, between $\phi(0)$ and $-\beta_{U}$ and converges to $-\beta_{U}$ as $t \rightarrow \infty$. This implies local asymptotic stability of the positive steady state. Moreover, using (3.1) and (3.2) one can show that every solution $x=x^{\phi}, \phi \in Z$, becomes positive on some interval $[\tilde{T}-\tau, \tilde{T}], \tilde{T} \geq 0$, and is given by

$$
x(t)=-\beta_{U}+\left(x(\tilde{T})+\beta_{U}\right) e^{-(t-\tilde{T})}
$$

for $t \geq \tilde{T}$, so it tends to $-\beta_{U}$ as $t \rightarrow \infty$. The proof in case $\beta_{L}<0$ is analogous.
The situation becomes more interesting when $\beta_{L}, \beta_{U}>0$ which gives $b_{U}<\gamma \theta<$ $b_{L}$ for the original parameters, so that the model Eq. (2.2) does not have a steady state, being a constant function. We will show in the next two theorems that in this case there exists a periodic solution and that it is stable, when the semiflow is restricted to the smaller set $Z_{0} \subset Z$ of all $\phi \in Z$ which have at most one zero and change sign at this zero $z$ in case $-\tau<z<0$.

We first show that $Z_{0}$ is positively invariant for the semiflow $S$, i.e., $S(t, \phi) \in Z_{0}$ for all $t \geq 0$ and $\phi \in Z_{0}$. Suppose that $\beta_{L}>0$ and $\beta_{U}>0$. For $\phi \in Z_{0}$ and
$x=x^{\phi}$ we make the following observations. If $\phi(z)=0$ and $-\tau<z<0$, and $\phi(t)<0$ for $-\tau \leq t<z$ and $0<\phi(t)$ for $z<t \leq 0$ then by (3.2), $0<x(t)$ on $(z, z+\tau]=(z, 0] \cup(0, z+\tau]$. In particular, $x_{z+\tau} \in Z_{0}$; moreover, $x_{t} \in Z_{0}$ for all $t \in[0, z+\tau]$. Using (3.1) there is a smallest $z^{\prime}$ in $(0, \infty)$ with $x\left(z^{\prime}\right)=0, x(t) \neq 0$ on $\left[z^{\prime}-\tau, z^{\prime}\right)$, and $x$ changes sign at $t=z^{\prime}$, and we can iterate. Thus we obtain $S(t, \phi) \in Z_{0}$ for all $t \geq 0$. The same holds for arbitrary $\phi \in Z_{0}$.

Note that the zeros of $x^{\phi}, \phi \in Z_{0}$ arbitrary, in $(0, \infty)$ are all transversal and form a strictly increasing sequence of times $z_{j}=z_{j}(\phi), j \in \mathbb{N}$, with

$$
z_{j}+\tau<z_{j+1} \text { and } x\left(z_{j}-\tau\right) \neq 0 \text { for all } j \in \mathbb{N} .
$$

From Proposition 3.2 we conclude the following concerning solutions $x:[-\tau, \infty) \rightarrow$ $\mathbb{R}$ starting from initial data $x_{0}=\phi \in Z_{0}$.

Corollary 3.1 Let $\beta_{L}, \beta_{U}>0$. Then each map

$$
Z_{0} \ni \phi \mapsto z_{j}(\phi) \in(0, \infty), \quad j \in \mathbb{N},
$$

is continuous at every point $\phi \in Z_{0}$ with $\phi(0) \neq 0$.
We now show the existence of a periodic solution.
Theorem 3.2 Let $\beta_{L}, \beta_{U}>0$. Then there is a periodic solution $\tilde{x}: \mathbb{R} \rightarrow \mathbb{R}$ of Eq. (2.5) with

$$
\tilde{x}(t)=-\beta_{U}+\beta_{U} e^{-(t+\tau)} \quad \text { for }-\tau \leq t \leq 0 .
$$

We have $\tilde{x}(-\tau)=0$, and with $\tilde{z}_{j}=z_{j}\left(\tilde{x}_{0}\right)$ for all $j \in \mathbb{N}$,

$$
\begin{aligned}
& \tilde{x}(t)=\beta_{L}+\left(\tilde{x}(0)-\beta_{L}\right) e^{-t} \quad \text { for } 0 \leq t \leq \tilde{z}_{1}+\tau \\
& \tilde{x}(t)=-\beta_{U}+\left(\tilde{x}\left(\tilde{z}_{1}+\tau\right)+\beta_{U}\right) e^{-\left(t-\left(\tilde{z}_{1}+\tau\right)\right)} \quad \text { for } \quad \tilde{z}_{1}+\tau \leq t \leq \tilde{z}_{2}+\tau .
\end{aligned}
$$

The minimal period of $\tilde{x}$ is

$$
\tilde{T}=\tilde{z}_{2}+\tau
$$

Proof Compute the solution (see Fig. 2) starting from $\phi \in Z_{0}$ given by

$$
\phi(t)=-\beta_{U}+\beta_{U} e^{-(t+\tau)}
$$

It is convenient to set $\tilde{z}_{0}=-\tau$.
Corollary 3.2 Let $\beta_{L}, \beta_{U}>0$. The minimum of $\tilde{x}$ is given by

$$
\underline{x}=\tilde{x}(0)=-\beta_{U}\left(1-e^{-\tau}\right)<0
$$



Fig. 2 A graph of the periodic solution on the interval $[-\tau, \tilde{T}]$
the maximum of $\tilde{x}$ is given by

$$
\bar{x}=\tilde{x}\left(\tilde{z}_{1}+\tau\right)=\beta_{L}\left(1-e^{-\tau}\right)>0
$$

and

$$
\begin{gathered}
\tilde{z}_{1}=\ln \frac{\beta_{L}-\underline{x}}{\beta_{L}}, \quad \tilde{z}_{2}=\tilde{z}_{1}+\tau+\ln \frac{\bar{x}+\beta_{U}}{\beta_{U}} \\
\tilde{T}=\tilde{z}_{2}+\tau=2 \tau+\ln \left(\frac{\beta_{L}-\underline{x}}{\beta_{L}} \cdot \frac{\bar{x}+\beta_{U}}{\beta_{U}}\right)
\end{gathered}
$$

Occasionally we shall abbreviate

$$
t_{\max }=\tilde{z}_{1}+\tau
$$

for the first non-negative time where $\tilde{x}$ achieves its maximum. Accordingly,

$$
t_{\min }=0
$$

Remark 3.2 For $\beta_{U}>0$ there are infinitely many other periodic orbits in $Z$, given by periodic solutions of higher oscillation frequencies, compare (Diekmann et al. 1995, Chapter XVI) which deals with the simpler equation

$$
x^{\prime}(t)=-\operatorname{sign}(x(t-1))
$$

Therefore the periodic orbit

$$
\tilde{O}=\left\{\tilde{x}_{t} \in Z: t \in \mathbb{R}\right\} \subset Z_{0}
$$

is not globally attracting on $Z$.
Theorem 3.3 Let $\beta_{L}, \beta_{U}>0$. Then for every $\phi \in Z_{0}$ we have either

$$
x^{\phi}\left(t+z_{1}(\phi)+\tau\right)=\tilde{x}(t) \quad \text { for all } t \geq-\tau
$$

or

$$
x^{\phi}\left(t+z_{1}(\phi)+\tau\right)=\tilde{x}\left(t+\tilde{z}_{1}+\tau\right) \quad \text { for all } t \geq-\tau
$$

and the periodic orbit $\tilde{O} \subset Z_{0}$ is stable.

Proof 1. For $\phi \in Z_{0}, x=x^{\phi}$, and $z_{1}=z_{1}(\phi)$ we infer from (3.1) and (3.2) that either $x$ is strictly decreasing on $\left[z_{1}, z_{1}+\tau\right]$, or that $x$ is strictly increasing on $\left[z_{1}, z_{1}+\tau\right]$. In the first case we obtain

$$
x\left(t+z_{1}+\tau\right)=\tilde{x}(t) \quad \text { for all } t \geq-\tau
$$

while in the second case,

$$
x\left(t+z_{1}+\tau\right)=\tilde{x}\left(t+\tilde{z}_{1}+\tau\right) \quad \text { for all } \quad t \geq-\tau
$$

In both cases, $x_{t}^{\phi}$ is on the periodic orbit $\tilde{O}$ for $t \geq z_{1}(\phi)+\tau$.
2. There exist $r>0$ and $\rho \in\left(0, \tilde{z}_{1}\right)$ with $\tilde{x}(t) \leq-r$ for $\tilde{z}_{1}-\rho-\tau \leq t \leq \tilde{z}_{1}-\rho$ and $\tilde{x}\left(\tilde{z}_{1}+\rho\right) \geq r$. It follows that for each $\psi \in \tilde{O}$ the shifted copy $x^{\psi}$ of $\tilde{x}$ has the property that there exists $u=u(\psi) \in[0, \tilde{T}]$ with $x^{\psi}(t) \leq-r$ for $u \leq t \leq u+\tau$ and $x^{\psi}(u+\tau+2 \rho) \geq r$. Observe that for every $\phi \in Z$ with $x^{\phi}(t)<0$ for $u \leq t \leq u+\tau$ and $x^{\phi}(u+\tau+2 \rho)>0$ the solution $x^{\phi}$ has a first zero $z$ in $(u+\tau, u+\tau+2 \rho)$, and $x_{z+\tau}^{\phi} \in \tilde{O}$, which implies that all segments $x_{t}^{\phi}$ with $t \geq \tilde{T}+2 \tau+2 \rho$ belong to the orbit $\tilde{O}$.
3. Let $\epsilon>0$. Continuous dependence on initial data yields that the map

$$
\left.Z \ni \phi \mapsto x^{\phi}\right|_{[-\tau, \tilde{T}+2 \tau+2 \rho]} \in C([-\tau, \tilde{T}+2 \tau+2 \rho], \mathbb{R}),
$$

where $\left.x^{\phi}\right|_{I}$ denotes the restriction of the function $x^{\phi}$ to the interval $I$, is continuous (with respect to the maximum-norm on the target space). Using uniform continuity on the compact orbit $\tilde{O} \subset Z$ we find $\delta>0$ so that for all $\phi \in Z$ and all $\psi \in \tilde{O}$ with $|\phi-\psi|_{C}<\delta$ we have

$$
\left|x^{\phi}(t)-x^{\psi}(t)\right|<\min \left\{\epsilon, \frac{r}{2}\right\} \quad \text { for all } t \in[-\tau, \tilde{T}+2 \tau+2 \rho] .
$$

Let $\phi \in Z$ with $\operatorname{dist}(\phi, \tilde{O})=\inf _{\psi \in \tilde{O}}|\phi-\psi|_{C}<\delta$ be given. Then for some $\psi \in \tilde{O},|\phi-\psi|_{C}<\delta$. Choose $u=u(\psi) \in[0, \tilde{T}]$ according to Part 2. The previous estimate of $\left|x^{\phi}(t)-x^{\psi}(t)\right|$ yields $x^{\phi}(t)<0$ for $u \leq t \leq u+\tau$ and $x^{\phi}(u+\tau+2 \rho)>0$. Using part 2 we infer $x_{t}^{\phi} \in \tilde{O}$ for $t \geq \tilde{T}+2 \tau+2 \rho$. Altogether,

$$
\operatorname{dist}\left(x_{t}^{\phi}, \tilde{O}\right)<\epsilon \quad \text { for } \quad 0 \leq t \leq \tilde{T}+2 \tau+2 \rho
$$

and $\operatorname{dist}\left(x_{t}^{\phi}, \tilde{O}\right)=0$ for $t \geq \tilde{T}+2 \tau+2 \rho$.

Remark 3.3 One can show that the other periodic orbits are all unstable, and that the domain of attraction of the periodic orbit $\tilde{O}$ is open and dense in $Z$, compare (Diekmann et al. 1995, Chapter XVI), and the main result of Mallet-Paret and Walther (1994) about Eq. (1.1) with a smooth and strictly monotone function $f$.

## 4 Pulse-like perturbations

In the following we assume

$$
-\beta_{U}<0<\beta_{L}
$$

and study a particular, simple deviation from the periodic solution $\tilde{x}$ and the subsequent return to the stable and attracting periodic orbit $\tilde{O}:$ We consider a function $x^{(\Delta)}: \mathbb{R} \rightarrow$ $\mathbb{R}$ which up to $t=\Delta \in[0, \tilde{T})$ equals the periodic solution $\tilde{x}$ of Eq. (2.5). Then for $\Delta \leq t \leq \Delta+\sigma$ the function $x^{(\Delta)}$ is defined by the equation

$$
x^{\prime}(t)=-x(t)+f(x(t-\tau))+a
$$

with a constant $a>0$. This results in a deviation from the periodic solution $\tilde{x}$ which begins at time $t=\Delta$ and lasts until the time $t=\Delta+\sigma$. Informally we speak of a pulse of amplitude $a$, with $\Delta$ the onset time of the pulse and $\sigma$ the duration of the pulse. For $t \geq \Delta+\sigma$, the function $x^{(\Delta)}$ is given again by Eq. (2.5). For perturbations $a>0$ not too large it will merge into the periodic solution in finite time, compare Theorem 3.3.

An interpretation of this is as follows. Eq. (2.5) is a mathematically convenient form of a (very simple) model for the production and decay of blood cells of a certain type, e.g. neutrophils. The periodic solution $\tilde{x}$ stands for the density of neutrophils in a patient as a function of time, perhaps induced by chemotherapy or as a consequence of cyclical neutropenia but in the absence of any further medical intervention. The function $x^{(\Delta)}$ describes the evolution of the neutrophil density for the case that at time $\Delta-\tau$ some medication has been administered which increases the production of cells in the bone marrow during the time interval [ $\Delta-\tau, \Delta-\tau+\sigma$ ]. The constant $a>0$ stands for the increase in production occasioned, for example, by the administration of G-CSF. After the time $\tau>0$ needed for production (and differentiation) of cells, that is, during the time interval $[\Delta, \Delta+\sigma]$ the neutrophils are released into the blood stream. Later on production and decay of neutrophils is again governed by the patient's feedback system alone.

For simplicity we assume

$$
0<\sigma \leq \tau,
$$

from here on, that is, the effect of intervention lasts for a time interval $\sigma$ less than the (production) delay $\tau$. The quantities we are interested in are the local extrema of $x^{(\Delta)}$ and the time required to return to the periodic orbit. The latter is captured by the cycle length map

$$
T:[0, \tilde{T}) \ni \Delta \mapsto T(\Delta) \in \mathbb{R} \cup\{\infty\}
$$

which is defined formally as follows: The zeros of $x^{(\Delta)}$ and of $\tilde{x}$ in $(-\infty, \Delta]$ coincide. Suppose $\tilde{z}_{J}, J=j(\Delta) \in\{0,1,2\}$, is the largest one of these zeros, and there exists a smallest zero $z>\tilde{z}_{J}$ of $x^{(\Delta)}$ with $x^{(\Delta)}(z+t)=\tilde{x}\left(\tilde{z}_{J}+t\right)$ for all $t \geq 0$. Then

$$
T(\Delta)=z-\tilde{z}_{J}, \quad \text { and } \quad T(\Delta)=\infty \quad \text { otherwise. }
$$

In case $T(\Delta)<\infty$ let

$$
\underline{x}_{\Delta}=\min _{\tilde{z}_{J \leq t \leq z}} x^{(\Delta)}(t)
$$

and

$$
\bar{x}_{\Delta}=\max _{\tilde{z}_{J \leq t \leq z}} x^{(\Delta)}(t)
$$

Remark 4.1 Observe that in our model situation treatment is considered successful if the minimal value $\underline{x}_{\Delta}$ is above the minimal value

$$
\underline{x}=\tilde{x}(0)=-\beta_{U}\left(1-e^{-\tau}\right)<0 \quad(\text { see Corollary 3.2) }
$$

of $\tilde{x}$ while in the opposite case medication actually increases the risk for the patient in the sense that the nadir of the oscillation is lower and would thus lead to more severe cytopenia which is one of the major clinical problems.

The local minima and maxima of the function $x^{(\Delta)}$ and the cycle length $T(\Delta)$ depend on the parameters

$$
\tau>0, \quad \beta_{L}>0>-\beta_{U}, \quad a>0, \quad \sigma \in(0, \tau], \quad \text { and } \quad \Delta \geq 0 .
$$

We assume that

$$
-\beta_{U}+a<0
$$

which will be instrumental in showing that after perturbation solutions do return to the periodic solution $\tilde{x}$, with the consequence that cycle lengths are finite. (For larger parameter $a$ solutions after perturbation may settle down on other, unstable periodic solutions of Eq. (2.5), with higher oscillation frequencies. For more on this, see Remark 5.4.) In the remainder of this section we keep the parameters $\tau, \beta_{L}, \beta_{U}, a, \sigma$ fixed. In addition to finiteness of cycle lengths we shall see that the cycle length map and the maps

$$
[0, \tilde{T}) \ni \Delta \mapsto \bar{x}_{\Delta} \in \mathbb{R} \text { and }[0, \tilde{T}) \ni \Delta \mapsto \underline{x}_{\Delta} \in \mathbb{R}
$$

are continuous.
For $a>0$ with $-\beta_{U}+a<0$ the solutions $x=x^{a, \phi}$ of the initial value problem

$$
x^{\prime}(t)=-x(t)+f(x(t-\tau))+a \text { for } t>0, x_{0}=\phi
$$

define a continuous semiflow $S_{a}:[0, \infty) \times Z \rightarrow Z$ by $S_{a}(t, \phi)=x_{t}^{a, \phi}$, see Sect. 3 . The set $Z_{0}$ is positively invariant under $S_{a}$, and for each $\phi \in Z_{0}$ the zeros of $x^{a, \phi}$ are all transversal and spaced at distances larger than the delay $\tau$.

The solution $x=x^{(\Delta)}$ during a pulse (which begins at $\Delta \in[0, \tilde{T})$ ) can now be described as follows: For $\Delta$ given, define $\phi=\tilde{x}_{\Delta}$ and then $\chi=S_{a}(\sigma, \phi)=x_{\sigma}^{a, \phi}$. We obtain

$$
\begin{aligned}
& x(t)=\tilde{x}(t) \text { for } t \leq \Delta, \\
& x(t)=x^{a, \phi}(t-\Delta) \text { for } \Delta \leq t \leq \Delta+\sigma, \\
& x(t)=x^{\chi}(t-(\Delta+\sigma)) \text { for } t \geq \Delta+\sigma,
\end{aligned}
$$

or equivalently,

$$
\begin{aligned}
& x_{t}=\tilde{x}_{t} \text { for } t \leq \Delta, \\
& x_{t}=S_{a}\left(t-\Delta, x_{\Delta}\right) \text { for } \Delta \leq t \leq \Delta+\sigma \\
& x_{t}=S\left(t-(\Delta+\sigma), x_{\Delta+\sigma}\right) \text { for } t \geq \Delta+\sigma
\end{aligned}
$$

Notice that for $x=x^{(\Delta)}$, all zeros are transversal and spaced at distances larger than the delay $\tau$. They form a strictly increasing sequence of points

$$
z_{\Delta, j}, \quad j \in \mathbb{Z}, \quad \text { with } \quad z_{\Delta, j}=\tilde{z}_{j} \text { for all integers } j \leq j_{\Delta}
$$

where $J=j_{\Delta} \in\{0,1,2\}$ is given by

$$
\tilde{z}_{J} \leq \Delta<\tilde{z}_{J+1}
$$

Using $\tilde{x}_{t}=S\left(t, \tilde{x}_{0}\right)$ for all $t \geq 0$ and the continuity of both semiflows we easily obtain from the previous representation of $x_{t}=x_{t}^{(\Delta)}$ that the map

$$
[0, \tilde{T}) \times[0, \infty) \ni(\Delta, t) \mapsto x_{t}^{(\Delta)} \in C
$$

is continuous, which in turn yields the continuity of the map

$$
[0, \tilde{T}) \times[0, \infty) \ni(\Delta, t) \mapsto x^{(\Delta)}(t) \in \mathbb{R}
$$

since $x^{(\Delta)}(t)=e v\left(x_{t}^{(\Delta)}\right)$ and the evaluation $e v: C \ni \phi \mapsto \phi(0) \in \mathbb{R}$ is continuous. Arguing as in the proof of Proposition 3.2 and using transversality of zeros we obtain the following results.

Proposition 4.1 For $\Delta_{0} \in[0, \tilde{T})$ and $z>0$ with $x^{(\Delta)}(z)=0$ and $\epsilon>0$ there exists $\delta>0$ such that for each $\Delta \in[0, \tilde{T})$ with $\left|\Delta-\Delta_{0}\right|<\delta$ there is $z^{\prime} \in(z-\epsilon, z+\epsilon)$ with $x^{\Delta}\left(z^{\prime}\right)=0$.

Corollary 4.1 Each map

$$
[0, \tilde{T}) \ni \Delta \mapsto z_{\Delta, j} \in(0, \infty), \quad j \in \mathbb{N}
$$

is continuous.

The proofs of the following results are provided in the appendix.
Proposition 4.2 For every $\Delta \in[0, \tilde{T})$ we have $T(\Delta)=z_{\Delta, J+2}-z_{\Delta, J}$ with $J=j_{\Delta}$.
Corollary 4.2 The cycle length map is continuous.
Proposition 4.3 The maps $[0, \tilde{T}) \ni \Delta \mapsto \bar{x}_{\Delta} \in \mathbb{R}$ and $[0, \tilde{T}) \ni \Delta \mapsto \underline{x}_{\Delta} \in \mathbb{R}$ are continuous.

## 5 Computation of the response

In the Sects. 5.1-5.4 below we keep the parameters $\tau>0, \beta_{L}>0>-\beta_{U}, a>$ $0, \sigma \in(0, \tau]$ fixed and require $-\beta_{U}+a<0$ as in the preceding section, and study the behaviour of $x^{(\Delta)}$ depending on the onset of the pulse (at $t=\Delta$ ) and on its termination (at $t=\Delta+\sigma$ ) relative to the zeros and extrema

$$
0<\tilde{z}_{1}<\tilde{z}_{1}+\tau=t_{\max }<\tilde{z}_{2}<\tilde{z}_{2}+\tau=\tilde{T}<\tilde{z}_{3}
$$

of the periodic solution $\tilde{x}$, on the sign of $x^{(\Delta)}(\Delta+\sigma)$, and on the position of $x^{(\Delta)}(\Delta+\sigma)$ relative to the level $\beta_{L}$.

The computations that follow in this section can become quite difficult to keep track of, and we therefore use what we hope is a simple and transparent nomenclature to aid the reader in following our progression. The reader may wish to consult Tables 1, 2 and 3 as a way of keeping track of the result.

Table 1 Summary of the effects of a positive pulse $(a>0, \underset{\tilde{\tau}}{\boldsymbol{\sim}}, \sigma)$ at different times during the limit cycle on the limit cycle minimum $(\underline{x})$, maximum $(\bar{x})$, and period $(\tilde{T})$

|  | Pulse location when $\delta_{2}<t_{\text {max }}$ with $\delta_{2}$ given by (5.11) |  |  |  |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :---: | :---: |
|  | RNRP | RPRP | RPFP | RPFN | FPFN | FNFN | FNRN |  |  |
| $\underline{x}_{\Delta}$ | U | U | U | $\uparrow \nearrow$ | $\uparrow \nearrow$ | $\uparrow$ | $\uparrow \searrow$ |  |  |
| $\bar{x}_{\Delta}$ | $\uparrow \nearrow$ | $\uparrow \nearrow$ | $\uparrow \searrow$ | $\uparrow \searrow$ | U | U | U |  |  |
| $T(\Delta)$ | $\nearrow$ | $\uparrow \nearrow$ | $\uparrow \nearrow$ | $\searrow$ | $\searrow$ | $\downarrow \searrow$ | $\downarrow \searrow$ |  |  |

"U" denotes unchanged, while $\uparrow$ means increased, $\downarrow$ means decreased, $\nearrow$ means that the value as a function of $\Delta$ is increasing, and $\searrow$ means that the values as a function of $\Delta$ is decreasing

Table 2 Summary of the effects of a positive pulse when $t_{\max } \leq \delta_{2}$

|  | Pulse location when $\delta_{2} \geq t_{\text {max }}$ |  |  |  |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :---: |
|  | RNRP | RPRP | RPFP | RPFN | FPFP | FPFN | FNFN | FNRN |  |
| $\underline{x}_{\Delta}$ | U | U | U | $\uparrow \nearrow$ | U | $\uparrow \nearrow$ | $\uparrow$ | $\uparrow \searrow$ |  |
| $\bar{x}_{\Delta}$ | $\uparrow \nearrow$ | $\uparrow \nearrow$ | $\uparrow \searrow$ | $\uparrow \searrow$ | U | U | U | U |  |
| $T(\Delta)$ | $\nearrow$ | $\uparrow \nearrow$ | $\uparrow \nearrow$ | $\searrow$ | $\uparrow \nearrow$ | $\searrow$ | $\downarrow \searrow$ | $\downarrow \searrow$ |  |

Table 3 Summary of the effects of a positive pulse in the remaining cases where $\delta_{1}$ is as in (5.6)

|  | $\delta_{1}>0$ <br> RNRN |
| :--- | :--- |
| $\underline{x}_{\Delta}$ | U |
| $\bar{x}_{\Delta}$ | U |
| $T(\Delta)$ | $\downarrow \searrow$ |
|  | $\delta_{1}<0$ |
|  | FNRP |
| $\underline{x}_{\Delta}$ | $\uparrow \nearrow$ |
| $\bar{x}_{\Delta}$ | $\uparrow \nearrow$ |
| $T(\Delta)$ | $\nearrow$ |

If the pulse starts at $\Delta \in\left[0, t_{\max }\right)$, where the periodic solution is increasing, then we say that we are in the rising phase and we use the letter $\mathbf{R}$. If it starts at $\Delta \in\left[t_{\max }, \tilde{T}\right)$, where the periodic solution is decreasing, then we say that we are in the falling phase and we use the letter $\mathbf{F}$. If $x^{(\Delta)}(t)$ is negative at the beginning of the pulse, i.e., $x^{(\Delta)}(\Delta)<0$, then we use the letter $\mathbf{N}$ (negative value at $\Delta$ ) and otherwise we write $\mathbf{P}$ (non-negative value at $\Delta$ ). We can thus say that we are in the subcase $\mathbf{R N}$ when we are at rising phase with a negative value at $\Delta$. Therefore the beginning of the pulse can be coded with two letters which gives four subcases: RN, RP, FN, FP. In the same way we can code the end of the pulse, namely if $\Delta+\sigma \in\left[t_{\max }, \tilde{T}\right)$, where the periodic solution is decreasing, then we are in the falling phase and we use the letter $\mathbf{F}$, otherwise we write R. If $x^{(\Delta)}(\Delta+\sigma)<0$, then we use the letter $\mathbf{N}$ and if $x^{(\Delta)}(\Delta+\sigma) \geq 0$ the letter $\mathbf{P}$. Here similarly, we can have four subcases and we combine them together to code each case with four letters. For example the case RNRN corresponds to the rising phase at $\Delta$ and at $\Delta+\sigma$ with negative values of $x^{(\Delta)}$ at $\Delta$ and at $\Delta+\sigma$.

There are three different periods of time that are important: before the pulse occurs, during the pulse, and after the pulse. We can easily write down the values $x^{(\Delta)}(t)$ before the pulse in each phase. If $\Delta \in\left[0, t_{\max }\right)$ then we have

$$
x^{(\Delta)}(t)=\beta_{L}+\left(\underline{x}-\beta_{L}\right) e^{-t}=\tilde{x}(t) \text { for } t \in[0, \Delta]
$$

and the value of $x^{(\Delta)}$ when the pulse turns on is

$$
x^{(\Delta)}(\Delta)=\beta_{L}+\left(\underline{x}-\beta_{L}\right) e^{-\Delta}=\tilde{x}(\Delta)
$$

From the definition of $\underline{x}$ it follows that $\beta_{L}-\underline{x}=\beta_{L} e^{\tilde{z}_{1}}$, which gives the following formula

$$
\begin{equation*}
x^{(\Delta)}(\Delta)=\beta_{L}-\beta_{L} e^{\tilde{z}_{1}-\Delta} \text { for } \Delta \in\left[0, t_{\max }\right) . \tag{5.1}
\end{equation*}
$$

If $\Delta \in\left[t_{\max }, \tilde{T}\right)$ then

$$
x^{(\Delta)}(t)=-\beta_{U}+\left(\bar{x}+\beta_{U}\right) e^{-\left(t-t_{\max }\right)}=\tilde{x}(t) \quad \text { for } t \in\left[t_{\max }, \Delta\right]
$$

and

$$
x^{(\Delta)}(\Delta)=-\beta_{U}+\left(\bar{x}+\beta_{U}\right) e^{-\left(\Delta-t_{\max }\right)}
$$

Since $\bar{x}+\beta_{U}=\beta_{U} e^{\tilde{z}_{2}-t_{\max }}$, we obtain

$$
\begin{equation*}
x^{(\Delta)}(\Delta)=-\beta_{U}+\beta_{U} e^{\tilde{z}_{2}-\Delta} \text { for } \Delta \in\left[t_{\max }, \tilde{T}\right) \tag{5.2}
\end{equation*}
$$

### 5.1 A pulse during the rising phase

We assume $0 \leq \Delta<\Delta+\sigma \leq t_{\max }$. Then $x^{(\Delta)}(\Delta)$ is given by (5.1). During the pulse,

$$
\begin{equation*}
x^{(\Delta)}(t)=\left(\beta_{L}+a\right)+\left(x^{(\Delta)}(\Delta)-\left(\beta_{L}+a\right)\right) e^{-(t-\Delta)} \quad \text { for } t \in[\Delta, \Delta+\sigma] \tag{5.3}
\end{equation*}
$$

and after the pulse,

$$
\begin{align*}
x^{(\Delta)}(t)= & \beta_{L}+\left(x^{(\Delta)}(\Delta+\sigma)-\beta_{L}\right) e^{-(t-(\Delta+\sigma))} \\
& \text { for } t>\Delta+\sigma \text { as long as } x^{(\Delta)}(t-\tau)<0, \tag{5.4}
\end{align*}
$$

with

$$
x^{(\Delta)}(\Delta+\sigma)=\beta_{L}+a+\left(x^{(\Delta)}(\Delta)-\left(\beta_{L}+a\right)\right) e^{-\sigma}
$$

Using (5.1) we have

$$
\begin{equation*}
x^{(\Delta)}(\Delta+\sigma)=\beta_{L}-\beta_{L} e^{\tilde{z}_{1}-\Delta-\sigma}+a\left(1-e^{-\sigma}\right) \tag{5.5}
\end{equation*}
$$

Case RNR The pulse starts before $\tilde{z}_{1}=t_{\max }-\tau, 0 \leq \Delta<\tilde{z}_{1}$.
Observe first that $\Delta \in\left[0, \tilde{z}_{1}\right)$ is such that $x^{(\Delta)}(\Delta+\sigma)<0$ if and only if

$$
\beta_{L}+a\left(1-e^{-\sigma}\right)<\beta_{L} e^{\tilde{z}_{1}-\Delta-\sigma}
$$

which in view of $\beta_{L}+a\left(1-e^{-\sigma}\right)>0$ is equivalent to

$$
e^{\Delta}<\frac{\beta_{L} e^{\tilde{z}_{1}-\sigma}}{\beta_{L}+a\left(1-e^{-\sigma}\right)}
$$

Let us define

$$
\begin{equation*}
\delta_{1}=\ln \frac{\beta_{L} e^{\tilde{z}_{1}-\sigma}}{\beta_{L}+a\left(1-e^{-\sigma}\right)}=\ln \frac{\beta_{L}+\beta_{U}\left(1-e^{-\tau}\right)}{\beta_{L} e^{\sigma}+a\left(e^{\sigma}-1\right)} . \tag{5.6}
\end{equation*}
$$

We have

$$
\delta_{1}=\tilde{z}_{1}-\sigma-\ln \frac{\beta_{L}+a\left(1-e^{-\sigma}\right)}{\beta_{L}}<\tilde{z}_{1} .
$$

If $\delta_{1}>0$ we obtain $x^{(\Delta)}(\Delta+\sigma)<0$ for $0 \leq \Delta<\delta_{1}$ and $x^{(\Delta)}(\Delta+\sigma) \geq 0$ for $\Delta \in\left[\delta_{1}, \tilde{z}_{1}\right)$ while if $\delta_{1} \leq 0$ we have $x^{(\Delta)}(\Delta+\sigma) \geq 0$ for all $\Delta \in\left[0, \tilde{z}_{1}\right)$. We consider three subcases.

Case RNRN The pulse parameters $(a, \Delta, \sigma)$ are such that $x^{(\Delta)}(t)$ remains negative during the pulse, $x^{(\Delta)}(\Delta+\sigma)<0$. Equivalently, $\delta_{1}>0$ and $\Delta \in\left[0, \delta_{1}\right)=$ $I_{R N R N}$.

Proposition 5.1 If $\Delta \in I_{R N R N}=\left[0, \delta_{1}\right)$ then $\underline{x}_{\Delta}=\underline{x}, \bar{x}_{\Delta}=\bar{x}$, and

$$
\begin{equation*}
T(\Delta)=\tilde{T}+\ln \left(1-\frac{a\left(e^{\sigma}-1\right)}{\beta_{L}} e^{\Delta-\tilde{z}_{1}}\right)<\tilde{T} \tag{5.7}
\end{equation*}
$$

In particular, the restriction of the map $T$ to $I_{R N R N}$ is strictly decreasing.
We now consider the case RNRP when $x^{(\Delta)}(\Delta+\sigma) \geq 0$. From Fig. 3 we expect that the first local maximum of $x^{(\Delta)}$ after $t=\Delta$ is achieved before $t=t_{\text {max }}$ and is not smaller than $\bar{x}$. In the following we prove this. Also we shall obtain a result about the cycle length $T(\Delta)$ of $x^{(\Delta)}$. However this time we can not conclude that either $T(\Delta)>\tilde{T}$ or $T(\Delta)<\tilde{T}$, see Figs. 3 and 4 , respectively.

From (5.3) in combination with $x^{(\Delta)}(\Delta)<0<\beta_{L}+a$ we see that $x^{(\Delta)}$ is strictly increasing on $[\Delta, \Delta+\sigma]$. So by $0 \leq x^{(\Delta)}(\Delta+\sigma)$ we obtain a first positive zero $z_{\Delta, 1}$ of $x^{(\Delta)}$, and

$$
\Delta<z_{\Delta, 1} \leq \Delta+\sigma
$$



Fig. 3 A schematic representation of the solution of the DDE when $0 \leq x^{(\Delta)}(\Delta+\sigma) \leq \beta_{L}$. The unperturbed periodic solution $\tilde{x}$ is the solid line and the solution $x^{(\Delta)}$ is the $\overline{\text { dashed line }}$


Fig. 4 A schematic representation of the solution of the DDE when $x^{(\Delta)}(\Delta+\sigma)>\beta_{L}$. As usual, the unperturbed periodic solution $\tilde{x}$ is the solid line and the solution $x^{(\Delta)}$ is the dashed line

Proposition 5.2 If $\Delta \in\left[\max \left\{0, \delta_{1}\right\}, \tilde{z}_{1}\right)$ then $T(\Delta)<\infty, \underline{x}_{\Delta}=\underline{x}, \bar{x}_{\Delta} \geq \bar{x}$, and

$$
\begin{equation*}
T(\Delta)=\tilde{T}+\ln \left(1+\frac{a\left(e^{\sigma}-1\right)}{\beta_{U}} e^{\Delta-\tilde{z}_{2}}+\frac{a\left(\beta_{L}+\beta_{U}\right) e^{\tau+\tilde{z}_{1}-\tilde{z}_{2}}}{\beta_{U}\left(\beta_{L}+a\right)}\left(e^{\Delta-\tilde{z}_{1}}-1\right)\right) \tag{5.8}
\end{equation*}
$$

Moreover, the maximal value $\bar{x}_{\Delta}$ is given by $\max \left\{x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right), x^{(\Delta)}(\Delta+\sigma)\right\}$ and is strictly increasing with respect to $\Delta \in\left[\max \left\{0, \delta_{1}\right\}, \tilde{z}_{1}\right)$.

Case RPRP The pulse occurs completely in the interval $\left[\tilde{z}_{1}, t_{\max }\right]$. This is equivalent to $\Delta \in\left[\tilde{z}_{1}, t_{\max }-\sigma\right]=I_{R P R P}$.

Proposition 5.3 If $\Delta \in\left[\tilde{z}_{1}, t_{\max }-\sigma\right]$ then $T(\Delta)<\infty, \underline{x}_{\Delta}=\underline{x}, \bar{x}_{\Delta}>\bar{x}$, and

$$
\begin{equation*}
T(\Delta)=\tilde{T}+\ln \left(1+\frac{a\left(e^{\sigma}-1\right)}{\beta_{U}} e^{\Delta-\tilde{z}_{2}}\right)>\tilde{T} \tag{5.9}
\end{equation*}
$$

Moreover, the map $\bar{x}_{\Delta}$ is strictly increasing on $I_{R P R P}=\left[\tilde{z}_{1}, t_{\max }-\sigma\right]$.
Remark 5.1 For $\delta_{1}>0$ and $\Delta$ close to $\delta_{1}$ we have $T(\Delta)<\tilde{T}$.

### 5.2 A pulse from the rising phase into the falling phase

Here we assume $\Delta \leq t_{\max }<\Delta+\sigma$. Then $x^{(\Delta)}(\Delta)$ is still given by (5.1). Since $\sigma \leq \tau$, we must have $\Delta>\tilde{z}_{1}$, thus $x^{(\Delta)}(\Delta)>0$. The largest zero of $\tilde{x}$ in $(-\infty, \Delta]$ is $\tilde{z}_{1}$. Also, with $t_{\max }=\tilde{z}_{1}+\tau$,

$$
\begin{aligned}
x^{(\Delta)}\left(t_{\max }\right) & =\beta_{L}+a+\left(x^{(\Delta)}(\Delta)-\left(\beta_{L}+a\right)\right) e^{-\left(t_{\max }-\Delta\right)} \\
& =\beta_{L}-\beta_{L} e^{-\tau}+a\left(1-e^{-\left(t_{\max }-\Delta\right)}\right) \\
& =\bar{x}+a\left(1-e^{\Delta-t_{\max }}\right) \geq \bar{x},
\end{aligned}
$$

and we can write

$$
x^{(\Delta)}(t)=-\beta_{U}+a+\left(x^{(\Delta)}\left(t_{\max }\right)+\beta_{U}-a\right) e^{-\left(t-t_{\max }\right)} \quad \text { for } t \in\left[t_{\max }, \Delta+\sigma\right]
$$

Observe that the function $\left[t_{\max }, \Delta+\sigma\right) \ni t \mapsto x^{(\Delta)}(t) \in \mathbb{R}$ is decreasing since $x^{(\Delta)}\left(t_{\max }\right)+\beta_{U}-a \geq 0$.

We have

$$
\begin{aligned}
x^{(\Delta)}(\Delta+\sigma) & =-\beta_{U}+a+\left(x^{(\Delta)}\left(t_{\max }\right)+\beta_{U}-a\right) e^{-\left(\Delta+\sigma-t_{\max }\right)} \\
& =-\beta_{U}+a+\left(\beta_{U} e^{\tilde{z_{2}}-t_{\max }}-a e^{\Delta-t_{\max }}\right) e^{-\left(\Delta+\sigma-t_{\max }\right)}
\end{aligned}
$$

which gives

$$
\begin{equation*}
x^{(\Delta)}(\Delta+\sigma)=-\beta_{U}+\beta_{U} e^{\tilde{z}_{2}-(\Delta+\sigma)}+a\left(1-e^{-\sigma}\right) \tag{5.10}
\end{equation*}
$$

## Using

$$
\tilde{z}_{1}+\tau<\Delta+\sigma \leq \tilde{z}_{1}+\tau+\sigma<\tilde{z}_{2}+\tau
$$

we have

$$
x^{(\Delta)}(t)=-\beta_{U}+\left(x^{(\Delta)}(\Delta+\sigma)+\beta_{U}\right) e^{-(t-(\Delta+\sigma))} \quad \text { for } \quad \Delta+\sigma \leq t \leq \tilde{z}_{2}+\tau
$$

and in particular

$$
x^{(\Delta)}(\Delta+\sigma)=\tilde{x}(\Delta+\sigma)+a\left(1-e^{-\sigma}\right)
$$

which yields $x^{(\Delta)}(\Delta+\sigma)>\tilde{x}(\Delta+\sigma)$.
We say that we are in
Case RPF We distinguish between the subcases

$$
\begin{aligned}
& \text { P } x^{(\Delta)}(\Delta+\sigma) \geq 0 \text {, see Fig. 5, } \\
& \mathbf{N} x^{(\Delta)}(\Delta+\sigma)<0 \text {, see Fig. } 6 .
\end{aligned}
$$

Note that $\Delta \in\left(t_{\max }-\sigma, t_{\max }\right]$ is such that $x^{(\Delta)}(\Delta+\sigma) \geq 0$ if and only if

$$
\beta_{U} e^{\tilde{z}_{2}-(\Delta+\sigma)} \geq \beta_{U}-a\left(1-e^{-\sigma}\right)
$$

Since $\beta_{U}>a \geq a\left(1-e^{-\sigma}\right)$, we have $\beta_{U}-a\left(1-e^{-\sigma}\right)>0$, which implies that $x^{(\Delta)}(\Delta+\sigma) \geq 0$ if and only if


Fig. 5 A schematic representation of the solution of the DDE for the case RPFP. The unperturbed limit cycle is the solid line while the solution with the pulse is the dashed line


Fig. 6 A schematic representation of the solution of the DDE for the case RPFN. The unperturbed limit cycle is the solid line while the solution with the pulse is the dashed line

$$
e^{\Delta} \leq \frac{\beta_{U} e^{\tilde{z}_{2}}}{\beta_{U} e^{\sigma}-a\left(e^{\sigma}-1\right)}
$$

Let us define $\delta_{2}$ by

$$
\begin{equation*}
\delta_{2}=\ln \frac{\beta_{U} e^{\tilde{z}_{2}-\sigma}}{\beta_{U}-a\left(1-e^{-\sigma}\right)} \tag{5.11}
\end{equation*}
$$

We conclude that the $\Delta$-intervals in the two subcases are of the form

$$
\begin{aligned}
& I_{R P F P}=\left(t_{\max }-\sigma, t_{\max }\right] \cap\left(-\infty, \delta_{2}\right] \quad \text { and } \\
& I_{R P F N}=\left(t_{\max }-\sigma, t_{\mathrm{max}}\right] \cap\left(\delta_{2}, \infty\right) .
\end{aligned}
$$

We have

$$
\delta_{2}=\tilde{z}_{2}-\sigma+\ln \frac{\beta_{U}}{\beta_{U}-a\left(1-e^{-\sigma}\right)} .
$$

Notice that $\delta_{2}>\tilde{z}_{2}-\sigma$, which implies that $\delta_{2}>t_{\max }-\sigma$. Consequently, if

$$
\beta_{U} e^{\sigma}-a\left(e^{\sigma}-1\right) \leq \beta_{U}+\beta_{L}\left(1-e^{-\tau}\right)
$$

which is equivalent to $\delta_{2} \geq t_{\max }$, then $x^{(\Delta)}(\Delta+\sigma) \geq 0$ for all $\Delta \in\left(t_{\max }-\sigma, t_{\max }\right]$. If the reverse inequality

$$
\begin{equation*}
\beta_{U}+\beta_{L}\left(1-e^{-\tau}\right)<\beta_{U} e^{\sigma}-a\left(e^{\sigma}-1\right) \tag{5.12}
\end{equation*}
$$

holds then $\delta_{2}<t_{\max }$ and $\delta_{2}$ is the maximal value of $\Delta \in\left(t_{\max }-\sigma, t_{\max }\right]$ such that $x^{(\Delta)}(t)$ changes the sign at $\Delta+\sigma$, from a positive value to a negative value, $x^{(\Delta)}(\Delta+\sigma) \geq 0$ for all $\Delta \in\left(t_{\max }-\sigma, \delta_{2}\right]$ and $x^{(\Delta)}(\Delta+\sigma)<0$ for all $\Delta \in\left(\delta_{2}, t_{\max }\right]$. Inequality (5.12) can be rewritten as

$$
\beta_{L}\left(1-e^{-\tau}\right)<\left(\beta_{U}-a\right)\left(e^{\sigma}-1\right)
$$

Proposition 5.4 If $\Delta \in I_{R P F P}$ then $T(\Delta)<\infty, \underline{x}_{\Delta}=\underline{x}, \bar{x}_{\Delta}>\bar{x}$, and $T(\Delta)$ is given by formula (5.9) as in Proposition 5.3. The map $I_{R P F P} \ni \Delta \mapsto \bar{x}_{\Delta} \in \mathbb{R}$ is strictly decreasing and $I_{R P F P} \ni \Delta \mapsto T(\Delta) \in \mathbb{R}$ is strictly increasing.

Proposition 5.5 If $\Delta \in I_{R P F N}$ then $T(\Delta)<\infty, \underline{x}_{\Delta}>\underline{x}$,

$$
\bar{x}_{\Delta}=\bar{x}+a\left(1-e^{\Delta-t_{\max }}\right) \geq \bar{x}
$$

and

$$
\begin{equation*}
T(\Delta)=\tilde{T}+\ln \left(1-\frac{a\left(e^{\sigma}-1\right)}{\beta_{L}} e^{\Delta-\tilde{z}_{1}-\tilde{T}}-\frac{a\left(\beta_{L}+\beta_{U}\right) e^{-\tilde{z}_{1}}}{\beta_{L}\left(\beta_{U}-a\right)}\left(e^{\Delta-\tilde{z}_{2}}-1\right)\right) \tag{5.13}
\end{equation*}
$$

Moreover, the map $I_{R P F N} \ni \Delta \mapsto \underline{x}_{\Delta} \in \mathbb{R}$ is strictly increasing and the map $I_{R P F N} \ni \Delta \mapsto T(\Delta) \in \mathbb{R}$ is strictly decreasing.
Remark 5.2 For $\Delta$ close to $\delta_{2}$ we have $T(\Delta)>\tilde{T}$.

### 5.3 A pulse during the falling phase

Suppose now that $x^{(\Delta)}(\Delta)$ is given by (5.2) for $\Delta \in\left[t_{\max }, \tilde{T}\right)$. We have

$$
x^{(\Delta)}(t)=-\beta_{U}+a+\left(x^{(\Delta)}(\Delta)+\beta_{U}-a\right) e^{-(t-\Delta)}
$$

for $\Delta \leq t \leq \Delta+\sigma$ as long as $x^{(\Delta)}(t-\tau)>0$.
We assume that $t_{\max } \leq \Delta<\Delta+\sigma \leq \tilde{z}_{2}+\tau=\tilde{T}$. Then the value of $x^{(\Delta)}$ at the end of the pulse is

$$
x^{(\Delta)}(\Delta+\sigma)=-\beta_{U}+a+\left(x^{(\Delta)}(\Delta)+\beta_{U}-a\right) e^{-\sigma}
$$

and, by (5.2), it is the same as in (5.10). We proceed in steps as before.
Case FPF The pulse starts in the interval $\left[t_{\max }, \tilde{z}_{2}\right]$. We distinguish between two subcases.
$\mathbf{P}$ The pulse parameters $(a, \Delta, \sigma)$ are such that $x^{(\Delta)}(t)$ remains non-negative during the pulse, and changes the sign after the pulse, $x^{(\Delta)}(\Delta+\sigma) \geq 0$.
$\mathbf{N}$ The pulse parameters $(a, \Delta, \sigma)$ are such that $x^{(\Delta)}$ changes the sign from positive to negative during the pulse (see Fig. 7), $t_{\max } \leq \Delta \leq \tilde{z}_{2}$ and $x^{(\Delta)}(\Delta+\sigma)<0$.

Since $\Delta \in\left[t_{\max }, \tilde{z}_{2}\right]$ it follows from (5.10) that $x^{(\Delta)}(\Delta+\sigma) \geq 0$ if and only if

$$
-\beta_{U}+\beta_{U} e^{\tilde{z}_{2}-\Delta-\sigma}+a\left(1-e^{-\sigma}\right) \geq 0
$$

or equivalently,

$$
\beta_{U} e^{\tilde{z}_{2}-\Delta} \geq \beta_{U} e^{\sigma}-a\left(e^{\sigma}-1\right) .
$$

The corresponding $\Delta$-intervals are of the form

$$
I_{F P F P}=\left[t_{\max }, \tilde{z}_{2}\right] \cap\left(-\infty, \delta_{2}\right]
$$

and

$$
I_{F P F N}=\left[t_{\max }, \tilde{z}_{2}\right] \cap\left(\delta_{2}, \infty\right)
$$

Proposition 5.6 If $\Delta \in I_{F P F P}$ then $T(\Delta)<\infty, \bar{x}_{\Delta}=\bar{x}, \underline{x}_{\Delta}=\underline{x}$, and $T(\Delta)$ is given by formula (5.9) in Proposition 5.3.

We turn to case FPFN. From Fig. 7 we expect $\tilde{z}_{2} \leq z_{\Delta, 2}$ and $\underline{x}_{\Delta}=x^{(\Delta)}\left(z_{\Delta, 2}+\tau\right)>$ $\tilde{x}\left(\tilde{z}_{2}+\tau\right)=\underline{x}$, that is, the minimum value of $x^{(\Delta)}$ is above the minimum value of $\tilde{x}$.

Proposition 5.7 If $\Delta \in I_{F P F N}$ then $T(\Delta)<\infty, \bar{x}_{\Delta}=\bar{x}, \underline{x}_{\Delta}>\underline{x}$, and both $\underline{x}_{\Delta}$ and $T(\Delta)$ are as in Proposition 5.5.


Fig. 7 A schematic representation of the solution of the DDE for case FPFN. The unperturbed periodic solution $\tilde{x}$ is the solid line while the solution $x^{(\Delta)}$ with the pulse is the dashed line


Fig. 8 A schematic representation of the solution for case FNFN. The unperturbed periodic solution $\tilde{x}$ is the solid line while the solution $x^{(\Delta)}$ with the pulse is the dashed line

Assume now that $\Delta \in\left(\tilde{z}_{2}, \tilde{z}_{2}+\tau-\sigma\right) . \operatorname{From} x^{(\Delta)}(\Delta)=-\beta_{U}+\left(0+\beta_{U}\right) e^{-\left(\Delta-\tilde{z}_{2}\right)}$ and

$$
x^{(\Delta)}(\Delta+\sigma)=-\beta_{U}+a+\left(x^{(\Delta)}(\Delta)+\beta_{U}-a\right) e^{-\sigma}
$$

we have $x^{(\Delta)}(\Delta)<0$ and $x^{(\Delta)}(\Delta+\sigma)<0$. Thus we consider the case
Case FNFN The pulse parameters $(a, \Delta, \sigma)$ are such that the pulse begins after $x^{(\Delta)}$ changes the sign from positive to negative and ends before the time $\tilde{T}=\tilde{z}_{2}+\tau$, $\tilde{z}_{2}<\Delta$ and $\Delta+\sigma<\tilde{z}_{2}+\tau$, see Fig. 8 .

From Fig. 8 we expect that the minimum value of $x^{(\Delta)}$ in $\left[\tilde{z}_{2}, \tilde{z}_{2}+\tau\right]$ is above the minimum value $\underline{x}$ of $\tilde{x}$, and the cycle length $T(\Delta)$ is below the minimal period $\tilde{T}$ of $\tilde{x}$. Observe that the function $[\Delta, \Delta+\sigma] \ni t \mapsto x^{(\Delta)}(t) \in \mathbb{R}$ is increasing (see Fig. 9) if and only if

$$
x^{(\Delta)}(\Delta)+\beta_{U}-a<0,
$$

which is equivalent to $\beta_{U} e^{\tilde{z}_{2}}<a e^{\Delta}$. Now, if $t \in[\Delta+\sigma, \tilde{T}]$ then $\tilde{z}_{1} \leq t-\tau \leq \tilde{z}_{2}$. Thus

$$
x^{(\Delta)}(t)=-\beta_{U}+\left(x^{(\Delta)}(\Delta+\sigma)+\beta_{U}\right) e^{-(t-(\Delta+\sigma))}
$$

and we have

$$
x^{(\Delta)}(\Delta+\sigma)+\beta_{U}=\beta_{U} e^{\tilde{z}_{2}-(\Delta+\sigma)}+a\left(1-e^{-\sigma}\right)
$$



Fig. 9 Graphs of two solutions for case FNFN, where the parameters are $\tau=1, \sigma=0.2, \beta_{U}=0.8$, $\beta_{L}=1, a=0.6$, and $\Delta \in\{2.2,2.7\}$, showing that the function $I_{F N F N} \ni \Delta \mapsto \underline{x}_{\Delta}$ is not given by one formula
which is always positive. Hence, the function $t \mapsto x^{(\Delta)}(t)$ is strictly decreasing on the interval $[\Delta+\sigma, \tilde{T}]$ and we have

$$
x^{(\Delta)}(\tilde{T})=-\beta_{U}+\left(x^{(\Delta)}(\Delta+\sigma)+\beta_{U}\right) e^{-(\tilde{T}-(\Delta+\sigma))}
$$

which becomes

$$
\begin{aligned}
x^{(\Delta)}(\tilde{T}) & =-\beta_{U}+\left(\beta_{U} e^{\tilde{z}_{2}-(\Delta+\sigma)}+a\left(1-e^{-\sigma}\right)\right) e^{-(\tilde{T}-(\Delta+\sigma))} \\
& =-\beta_{U}+\beta_{U} e^{-\tau}+a\left(e^{\sigma}-1\right) e^{\Delta-\tilde{T}}
\end{aligned}
$$

Define $\bar{\delta} \in \mathbb{R}$ by

$$
e^{\tilde{T}-\bar{\delta}}=\frac{\beta_{U}+\sqrt{\beta_{U}^{2}+4 a \beta_{U}\left(e^{\sigma}-1\right) e^{\tau}}}{2 \beta_{U}}
$$

Proposition 5.8 If $\Delta \in I_{F N F N}$ then $T(\Delta)<\infty, \bar{x}_{\Delta}=\bar{x}$,

$$
\underline{x}_{\Delta}=\min \left\{\tilde{x}(\Delta), x^{(\Delta)}(\tilde{T})\right\}>\underline{x},
$$

and

$$
\begin{equation*}
T(\Delta)=\tilde{T}+\ln \left(1-\frac{a\left(e^{\sigma}-1\right)}{\beta_{L}} e^{\Delta-\tilde{z}_{1}-\tilde{T}}\right)<\tilde{T} \tag{5.14}
\end{equation*}
$$

In case $I_{F N F N} \cap(-\infty, \bar{\delta}) \neq \emptyset$ the map

$$
I_{F N F N} \cap(-\infty, \bar{\delta}) \ni \Delta \mapsto \underline{x}_{\Delta} \in \mathbb{R}
$$

is strictly increasing while in case $I_{F N F N} \cap(\bar{\delta}, \infty) \neq \emptyset$ the map

$$
I_{F N F N} \cap(\bar{\delta}, \infty) \ni \Delta \mapsto \underline{x}_{\Delta} \in \mathbb{R}
$$

is strictly decreasing.


Fig. 10 Graphs of two solutions for case FNFP, where the parameters are $\tau=1, \sigma=0.4, \beta_{U}=0.6$, $\beta_{L}=0.3, a=0.95$, and $\Delta \in\{2.15,2.3\}$, showing that $T(\Delta)$ might not exists


Fig. 11 A graph of one solution for case FNFP, where the parameters are $\tau=1, \sigma=0.35, \beta_{U}=0.3$, $\beta_{L}=0.4, a=0.7$, and $\Delta=2.22$, suggesting that for certain values of the parameters the solution $x^{(\Delta)}$ might settle down on a rapidly oscillating periodic solution, in which case we would have $T(\Delta)=\infty$

Remark 5.3 If $\beta_{U} e^{\sigma-\tau}<a$ then $\bar{\delta} \in I_{F N F N}=\left(\tilde{z}_{2}, \tilde{T}-\sigma\right)$. If $\beta_{U} e^{\sigma-\tau} \geq a$ then $\bar{\delta} \geq \tilde{T}-\sigma$.

Remark 5.4 The following relations are equivalent:

$$
\begin{aligned}
a & <\beta_{U}, \\
a\left(e^{\sigma}-1\right) & <\beta_{U}\left(e^{\sigma}-1\right), \\
\beta_{U} & <\left(\beta_{U}-a\left(1-e^{-\sigma}\right)\right) e^{\sigma}, \\
\ln \frac{\beta_{U}}{\beta_{U}-a\left(1-e^{-\sigma}\right)} & <\sigma, \\
\delta_{2} & <\tilde{z}_{2} .
\end{aligned}
$$

Observe that we have $\tilde{z}_{2} \leq \delta_{2}<\tilde{T}-\sigma$ if and only if

$$
\begin{equation*}
\beta_{U} \leq a \quad \text { and } \quad a\left(1-e^{-\sigma}\right)<\beta_{U}\left(1-e^{-\tau}\right) \tag{5.15}
\end{equation*}
$$

which is excluded by our standing hypothesis $a<\beta_{U}$.
Let us briefly address what might happen if (5.15) holds. In that case if $\Delta \in$ $\left(\tilde{z}_{2}, \tilde{T}-\sigma\right)$ then $x^{(\Delta)}(\Delta+\sigma) \geq 0$ if and only if $\Delta \leq \delta_{2}$. We have $I_{F N F N}=\left(\delta_{2}, \tilde{T}-\sigma\right)$ and Proposition 5.8 remains true. The case FNFP is possible with $I_{F N F P}=\left(\tilde{z}_{2}, \delta_{2}\right]$. Figure 10 and in particular Fig. 11 indicate that in the interval $I_{F N F P}$ the cycle length may not be finite everywhere, due to a higher oscillation frequency of the solution on $[\Delta+\sigma, \infty)$.

### 5.4 A pulse from the falling phase into the rising phase

Here we suppose that $\tilde{z}_{2}<\Delta<\tilde{T} \leq \Delta+\sigma$ and that $x^{(\Delta)}(\Delta)$ is given by (5.2) for $\Delta \in[\tilde{T}-\sigma, \tilde{T})$. We have

$$
x^{(\Delta)}(t)=-\beta_{U}+a+\left(\tilde{x}(\Delta)+\beta_{U}-a\right) e^{-(t-\Delta)} \text { for } t \in(\Delta, \tilde{T}]
$$



Fig. 12 A schematic representation of the solution of the DDE for the case FNRN. The unperturbed limit cycle is the solid line while the solution with the pulse is the dashed line
and the function $[\Delta, \tilde{T}] \ni t \mapsto x^{(\Delta)}(t) \in \mathbb{R}$ is either strictly increasing, or decreasing. We have

$$
\begin{aligned}
x^{(\Delta)}(\tilde{T}) & =-\beta_{U}+a+\left(x^{(\Delta)}(\Delta)+\beta_{U}-a\right) e^{-(\tilde{T}-\Delta)} \\
& =-\beta_{U}+a+\left(-\beta_{U}+\beta_{U} e^{-\left(\Delta-\tilde{z}_{2}\right)}+\beta_{U}-a\right) e^{-(\tilde{T}-\Delta)} \\
& =-\beta_{U}+\beta_{U} e^{-\tau}+a\left(1-e^{-(\tilde{T}-\Delta)}\right) \\
& =\underline{x}+a\left(1-e^{\Delta-\tilde{T}}\right)>\underline{x},
\end{aligned}
$$

which shows that the map

$$
[\tilde{T}-\sigma, \tilde{T}) \ni \Delta \mapsto x^{(\Delta)}(\tilde{T}) \in \mathbb{R}
$$

is strictly decreasing.
For $t \in[\tilde{T}, \Delta+\sigma]$ we have

$$
x^{(\Delta)}(t)=\beta_{L}+a+\left(x^{(\Delta)}(\tilde{T})-\left(\beta_{L}+a\right)\right) e^{-(t-\tilde{T})}
$$

Since

$$
x^{(\Delta)}(\tilde{T})-\left(\beta_{L}+a\right)=-\left(\beta_{L}-\underline{x}\right)-a e^{\Delta-\tilde{T}}<0
$$

the function $[\tilde{T}, \Delta+\sigma] \ni t \mapsto x^{(\Delta)}(t) \in \mathbb{R}$ is strictly increasing. Also

$$
\begin{aligned}
x^{(\Delta)}(\Delta+\sigma) & =\beta_{L}+a+\left(x^{(\Delta)}(\tilde{T})-\left(\beta_{L}+a\right)\right) e^{-(\Delta+\sigma-\tilde{T})} \\
& =\beta_{L}+a+\left(\underline{x}-a e^{\Delta-\tilde{T}}-\beta_{L}\right) e^{-(\Delta+\sigma-\tilde{T})} \\
& =\beta_{L}+\left(\underline{x}-\beta_{L}\right) e^{-(\Delta+\sigma-\tilde{T})}+a\left(1-e^{-\sigma}\right) \quad(>\tilde{x}(\Delta+\sigma)) \\
& =\beta_{L}-\beta_{L} e^{\tilde{z}_{1}-(\Delta+\sigma-\tilde{T})}+a\left(1-e^{-\sigma}\right) .
\end{aligned}
$$

Case FNR We must distinguish between the two subcases
$\mathbf{N} x^{(\Delta)}(\Delta+\sigma)<0$, see Fig. 12,
P $x^{(\Delta)}(\Delta+\sigma) \geq 0$, see Fig. 13 .
We have $\tilde{T}-\sigma \leq \Delta<\tilde{T}$, and the condition for subcase FNRN, namely,

$$
0>x^{(\Delta)}(\Delta+\sigma)=\beta_{L}-\beta_{L} e^{\tilde{z}_{1}-(\Delta+\sigma-\tilde{T})}+a\left(1-e^{-\sigma}\right)
$$



Fig. 13 A schematic representation of the solution of the DDE for the case FNRP
is equivalent to

$$
\beta_{L}+a\left(1-e^{-\sigma}\right)<\beta_{L} e^{\tilde{z}_{1}-(\Delta+\sigma-\tilde{T})},
$$

or

$$
e^{\Delta}<e^{\tilde{T}} \frac{\beta_{L} e^{\tilde{z}_{1}}}{\beta_{L} e^{\sigma}+a\left(e^{\sigma}-1\right)}=e^{\tilde{T}+\delta_{1}} .
$$

Thus we are in case FNRN if and only if

$$
\Delta \in[\tilde{T}-\sigma, \tilde{T}) \cap\left(-\infty, \tilde{T}+\delta_{1}\right)=I_{F N R N}
$$

and we are in case FNRP if and only if

$$
\Delta \in[\tilde{T}-\sigma, \tilde{T}) \cap\left[\tilde{T}+\delta_{1}, \infty\right)
$$

Proposition 5.9 If $\Delta \in I_{F N R N}$ then $T(\Delta)<\infty, \bar{x}_{\Delta}=\bar{x}, \underline{x}_{\Delta}>\underline{x}$, and $T(\Delta)$ is as in (5.14) of Proposition 5.8. The map $I_{F N R N} \ni \Delta \mapsto \underline{x}_{\Delta} \in \mathbb{R}$ is strictly decreasing.

The interval $I_{F N R N}$ in Proposition 5.9 is nonempty if and only if $\delta_{1}>-\sigma$, which is always the case. To see this observe that

$$
e^{\delta_{1}}=e^{-\sigma} \frac{\beta_{L}-\underline{x}}{\beta_{L}+a\left(1-e^{-\sigma}\right)}>e^{-\sigma}
$$

since

$$
-\underline{x}=\beta_{U}\left(1-e^{-\tau}\right)>a\left(1-e^{-\sigma}\right) .
$$

We next proceed to the subcase FNRP. Note that the interval $I_{F N R P}$ is empty if and only if $\delta_{1} \geq 0$.

Proposition 5.10 If $\Delta \in I_{F N R P}$ then $T(\Delta)<\infty, \bar{x}_{\Delta}>\bar{x}, \underline{x}_{\Delta}>\underline{x}$, and

$$
\begin{equation*}
T(\Delta)=\tilde{T}+\ln \left(1+\frac{a\left(e^{\sigma}-1\right)}{\beta_{U}} e^{\Delta-\tilde{z}_{2}-\tilde{T}}+\frac{a\left(\beta_{L}+\beta_{U}\right) e^{\tau+\tilde{z}_{1}-\tilde{z}_{2}}}{\beta_{U}\left(\beta_{L}+a\right)}\left(e^{\Delta-\tilde{z}_{1}-\tilde{T}}-1\right)\right) \tag{5.16}
\end{equation*}
$$

The map $I_{F N R P} \ni \Delta \mapsto \underline{x}_{\Delta} \in \mathbb{R}$ is strictly decreasing as is the map $I_{F N R P} \ni \Delta \mapsto$ $\bar{x}_{\Delta} \in \mathbb{R}$.

## 6 The cycle length map

The behaviour of the cycle length map $[0, \tilde{T}) \ni \Delta \mapsto T(\Delta) \in \mathbb{R}$, illustrated in Figs. 14 and 15, is different in each of the (sub-) cases discussed in Sects. 5.1-5.4. Each of these cases corresponds to $\Delta$ varying in one of the subintervals $I_{R N R N}, \ldots, I_{F N R P}$ of $[0, \tilde{T})$. If $\Delta$ increases from 0 to $\tilde{T}$ then it travels through those subintervals which are not empty for the given parameter vector ( $\tau, \beta_{L}, \beta_{U}, a, \sigma$ ) (with $0<\tau,-\beta_{U}<0<$ $\left.\beta_{L}, 0<a<\beta_{U}, 0<\sigma \leq \tau\right)$. In other words, for each parameter vector we have a finite sequence of non-empty subintervals, ordered by, say, their left endpoints, whose union is $[0, \tilde{T})$. Below we describe the possible scenarios, in terms of sequences of cases and subcases. Recall also the Tables 1, 2, 3 in Sect. 5.

Before doing so it may be convenient to collect a few facts about the quantities $\delta_{1}, \delta_{2}$ which, together with $\tilde{z}_{1}, \tilde{z}_{2}, t_{\max }, \sigma, \tilde{T}$, determine the intervals $I_{R N R N}, \ldots, I_{F N R P}$. Recall that $\delta_{1}$ was defined by (5.6) and that $\delta_{2}$ was defined by (5.11). From $\beta_{U}>a$ and $\tau \geq \sigma$ we have $\delta_{1}>-\sigma$ and $\delta_{2}<\tilde{z}_{2}$. The condition $\delta_{1}>0$ is equivalent to

$$
\left(\beta_{L}+a\right)\left(e^{\sigma}-1\right)<\beta_{U}\left(1-e^{-\tau}\right)
$$

and the case when $\delta_{1} \in(-\sigma, 0]$ is described by

$$
a\left(1-e^{-\sigma}\right)<\beta_{U}\left(1-e^{-\tau}\right) \leq\left(\beta_{L}+a\right)\left(e^{\sigma}-1\right)
$$



Fig. 14 A graph of the function $[0, \tilde{T}] \ni \Delta \mapsto T(\Delta)$. The straight line represents the graph of $[0, \tilde{T}] \ni$ $\Delta \mapsto \tilde{T}$. We indicated with dots the values of the boundaries of all cases, and with lines the boundaries of subcases. Here the parameters are $\tau=1, \beta_{U}=0.8, \beta_{L}=0.4$, and $\sigma=0.4, a=0.2$, so that $\delta_{1}>0$ and $\delta_{2}<t_{\text {max }}$. The figure corresponds to the following sequence of subcases: RNRN, RNRP, RPRP, RPFP, RPFN, FPFN, FNFN, FNRN


Fig. 15 As in Fig. 14, but with $\beta_{L}=1.4$, which gives $\delta_{1} \in(-\sigma, 0), \delta_{2} \in\left(t_{\max }, \tilde{z}_{2}\right)$, and the following sequence of subcases: RNRP, RPRP, RPFP, FPFP, FPFN, FNFN, FNRN, FNRP

Next, the condition $\delta_{2}<t_{\text {max }}$ is equivalent to

$$
\beta_{L}\left(1-e^{-\tau}\right)<\left(\beta_{U}-a\right)\left(e^{\sigma}-1\right)
$$

while the condition $t_{\text {max }} \leq \delta_{2}<\tilde{z}_{2}$ is equivalent to

$$
0<\left(\beta_{U}-a\right)\left(e^{\sigma}-1\right) \leq \beta_{L}\left(1-e^{-\tau}\right)
$$

Since $\delta_{1}<\tilde{z}_{1}$ and $\delta_{2}>t_{\text {max }}-\sigma$, we conclude that the intervals

$$
I_{R N R P}=\left[\min \left\{0, \delta_{1}\right\}, \tilde{z}_{1}\right), \quad I_{R P F P}=\left(t_{\max }-\sigma, \min \left\{t_{\max }, \delta_{2}\right\}\right]
$$

are always nonempty. We have $I_{R P R P}=\left[\tilde{z}_{1}, t_{\max }-\sigma\right], t_{\max }=\tilde{z}_{1}+\tau$, and $\tau \geq \sigma$, thus each sequence of cases contains

RNRP, RPRP, and RPFP.

For $\delta_{1}>0$ each sequence starts with the case RNRN and ends with the case FNRN. For $\delta_{1}=0$ each sequence starts with the case RNRP and ends with the case FNRN.
For $\delta_{1}<0$ each sequence starts with the case RNRP and ends with the case FNRP.
If $\delta_{1}>0$ and $0 \leq \Delta<\delta_{1}$ then we are in case RNRN. If $\Delta$ grows from max $\left\{0, \delta_{1}\right\}$ to $\min \left\{t_{\text {max }}, \delta_{2}\right\}$ then we have the sequence of cases:

## RNRP, RPRP, and RPFP.

If $\Delta$ grows from $\min \left\{t_{\max }, \delta_{2}\right\}$ to $\min \left\{0, \delta_{1}\right\}+\tilde{T}$ and if $\delta_{2}<t_{\max }$ then we obtain the subsequent cases

## RPFN, FPFN, FNFN, FNRN.

In case $\delta_{2}=t_{\text {max }}$ we obtain
FPFP, FPFN, FNFN, FNRN.
The same sequence results in case $t_{\max }<\delta_{2}$. So we have two scenarios for $\Delta$ beyond the interval $I_{R P F P}$ and below $\min \left\{0, \delta_{1}\right\}+\tilde{T}$, which is the endpoint $\tilde{T}$ of the domain of the cycle length map if $\delta_{1} \geq 0$, while for $\delta_{1}<0$ the sequence of cases is completed by

## FNRP.

## 7 A therapy plan

We first describe the concept of the therapy plan, in case the evolution of the number of cells in the bloodstream (of a patient, without medical treatment) is governed by

Eq. (2.2), with the production function $f$ given by (2.1), and $b_{U}<\gamma \theta<b_{L}$. For convenience we shall work not with the original variables but with the transformed quantities from Sect. 2, namely, Eq. (2.5) with $f$ satisfying Eq. (2.6) for $-\beta_{U}<0<$ $\beta_{L}$. Then the variable $x(t)$ still represents the number of blood cells (of a certain type) in a patient, at time $t$.

The reader will find it helpful to consult Fig. 16 when following the argument below.
Suppose there is a critical level $x_{d}$, larger than the minimum $\underline{x}$ of the periodic oscillation in the patient without treatment. We want to find a therapy plan which consists of medication at certain times $t=t_{M}$ (which are to be determined) in such a way that the cell density in the patient never falls below the critical level.

Medication at a time $t=t_{M}$ results in the begin of the production of more (precursors of) cells in the bone marrow, and this increased production lasts for a time interval of duration $\sigma>0$, from the time $t_{M}$ until the time $t_{M}+\sigma$. As in Sect. 4 the effect of medication at $t=t_{M}$ can be expressed by a 'temporal change' of the production function $f$, for example, by replacing $f$ by a sum $f_{a}=f+a$ with $a>0$ as long as $t_{M} \leq t-\tau \leq t_{M}+\sigma$. (Alternatively, one might replace $f$ by a multiple $f_{a}=a f$ with $a>1$.) $a$ depends on the dose of the medication. Because of the delay $\tau$ due to the production process the number of cells in the bloodstream will begin to deviate from their number without treatment not earlier than $t \geq t_{M}+\tau$, where $t_{M}+\tau$ equals $\Delta$ in the cases studied in Sect. 5.

We begin with a simple situation and assume that for the time interval $[-\tau, 0]$ in the past the cell density in a patient is known, for example, by measurement, and

$$
0<x(t) \text { for }-\tau<t \leq 0 .
$$

$t=0$ stands for the present time. Then we use Eq. (2.5) in order to predict how the cell density would evolve in the patient without treatment: We compute the solution $y(t), t \geq 0$, of Eq. (2.5) with initial data $y(t)=x(t),-\tau \leq t \leq 0$. The solution $y(t)$ will be a translate of the unique slowly oscillating periodic solution $\tilde{x}$ of Eq. (2.5). There is a first zero $z_{1}=z_{1}(y(0))>0$ of $y$, and there exists a unique time $t_{d}$ between $z_{1}$ and $z_{1}+\tau$, at which $y$ reaches the critical level $x_{d}, y\left(t_{d}\right)=x_{d}$, from above. (At $z_{1}+\tau y$ attains its minimum value $\underline{x}<x_{d}$.)

Having predicted the time $t_{d}$ we define $t_{M}=t_{d}-\tau-\sigma$ as the time of medication. If $t_{M}$ is positive (is in the future), then it is not too late for medication. After medication


Fig. 16 A schematic representation of the ideas behind the arguments leading to a therapy plan. The unperturbed limit cycle is the solid line while the desired limit cycle due to a perturbation is the dashed line
at $t=t_{M}$ the cell density in the patient represented by $x(t)$ will equal $y(t)$ for $-\tau \leq t \leq t_{M}+\tau=t_{d}-\sigma$, because of the delay in Eq. (2.5). For $t_{d}-\sigma \leq t \leq t_{d}$ the release of cells into the circulation will be increased according to

$$
x^{\prime}(t)=f(x(t-\tau))+a-x(t)=-\beta_{U}+a-x(t)
$$

while for $t \geq t_{d}$, Eq. (2.5) holds once again.
The question is whether for a range of parameters $a>0$ this can be done in such a way that for $t_{d}-\sigma \leq t<z_{1}+\tau$ the solution $x(t)$ satisfies $x_{d} \leq x(t)<0$.

If yes then $x(t)$ would increase after time $z_{1}+\tau$ until there is a zero $z_{2}$, due to Eq. (2.5), and would coincide on $\left[z_{2}, z_{2}+\tau\right]$ with the piece of the unique slowly oscillating periodic solution $\tilde{x}$ of Eq. (2.5) before its maxima.

Upon that the whole process can be repeated. It would result in a periodic therapy plan and a periodic solution $x(t)$ which never falls below the critical value $x_{d}$ and has a period shorter than the period of $\tilde{x}$. (This latter property comes from $x\left(z_{1}+\tau\right) \geq$ $x_{d}>\underline{x}$.)

Below we show that there exist parameters $\beta_{L}, \beta_{U}, \tau, \sigma$ and $x_{d} \in(\underline{x}, 0)$ and $a>0$ for which the program just described can be carried out. As initial data we consider continuous functions $\phi:[-\tau, 0] \rightarrow \mathbb{R}$ with $0<\phi(t)$ for $-\tau<t \leq 0$. Then for some $q, 0<q<1, q \bar{x}<\phi(0)$. The solution $y$ of the initial value problem

$$
y^{\prime}(t)=f(y(t-\tau))-y(t) \text { for } t>0, y_{0}=\phi
$$

has a first zero at

$$
z_{1}=z_{1}(\phi(0))=\ln \frac{\phi(0)+\beta_{U}}{\beta_{U}}
$$

and strictly decreases on $\left[z_{1}, z_{1}+\tau\right]$ to the value $\underline{x}$. For $\underline{x}<x_{d}<0$ we find a unique time $t=t_{d}=t_{d}(\phi)$ in $\left(z_{1}, z_{1}+\tau\right)$ with $y\left(t_{d}\right)=x_{d}$, namely,

$$
t_{d}=z_{1}+\ln \frac{\beta_{U}}{x_{d}+\beta_{U}} .
$$

Next we show that for parameters $-\beta_{U}<0<\beta_{L}, \tau>0, q \in(0,1)$ and $\sigma>0$ sufficiently small, and $x_{d} \in(\underline{x}, 0)$ sufficiently close to $\underline{x}$ we have

$$
\begin{equation*}
t_{d}-\tau-\sigma>0 \tag{7.1}
\end{equation*}
$$

The inequality (7.1) is equivalent to

$$
\tau+\sigma<\ln \frac{\phi(0)+\beta_{U}}{\beta_{U}}+\ln \frac{\beta_{U}}{x_{d}+\beta_{U}}=\ln \frac{\phi(0)+\beta_{U}}{x_{d}+\beta_{U}},
$$

which follows from

$$
\tau+\sigma<\ln \frac{q \bar{x}+\beta_{U}}{x_{d}+\beta_{U}} .
$$

The preceding inequality can be achieved for $\sigma>0$ sufficiently small and $x_{d}>\underline{x}$ sufficiently close to $\underline{x}$ provided we have

$$
\begin{equation*}
\tau<\ln \frac{q \bar{x}+\beta_{U}}{\underline{x}+\beta_{U}} . \tag{7.2}
\end{equation*}
$$

We verify this: Using the equations for $\bar{x}$ and $\underline{x}$ from Corollary 3.2 we see that (7.2) is equivalent to

$$
\begin{aligned}
\tau & <\ln \frac{q \beta_{L}\left(1-e^{-\tau}\right)+\beta_{U}}{\beta_{U}-\beta_{U}\left(1-e^{-\tau}\right)} \\
& =\tau+\ln \frac{q \beta_{L}\left(1-e^{-\tau}\right)+\beta_{U}}{\beta_{U}}
\end{aligned}
$$

or,

$$
\beta_{U}<q \beta_{L}\left(1-e^{-\tau}\right)+\beta_{U}
$$

which is equivalent to

$$
\begin{equation*}
0<q \beta_{L}\left(1-e^{-\tau}\right) \tag{7.3}
\end{equation*}
$$

From now on assume that the parameters $-\beta_{U}<0<\beta_{L}, \tau>0, \sigma>0$, and $q \in(0,1), x_{d} \in(\underline{x}, 0)$ satisfy (7.1). Assume in addition for simplicity that $\sigma$ is so small that we have

$$
\begin{equation*}
z_{1}<t_{d}-\sigma \tag{7.4}
\end{equation*}
$$

Notice that $t_{d}-z_{1}=\ln \frac{\beta_{U}}{x_{d}+\beta_{U}}$ does not depend on $\phi$. We now define

$$
t_{M}=t_{M}(\phi)=t_{d}(\phi)-\tau-\sigma(>0)
$$

as the time of medication. For parameters $a>0$ we consider the continuous function $x:\left[-\tau, z_{1}+\tau\right] \rightarrow \mathbb{R}$ which coincides with $y(t)$ for $-\tau \leq t \leq t_{d}-\sigma$ and satisfies

$$
x^{\prime}(t)= \begin{cases}f(x(t-\tau))+a-x(t)=-\beta_{U}+a-x(t) & \text { for } t_{d}-\sigma<t<t_{d} \\ f(x(t-\tau))-x(t)=-\beta_{U}-x(t) & \text { for } t_{d}<t<z_{1}+\tau\end{cases}
$$

It follows that

$$
x\left(t_{d}-\sigma\right)=y\left(t_{d}-\sigma\right)=-\beta_{U}+\beta_{U} e^{-\left(t_{d}-\sigma-z_{1}\right)}=-\beta_{U}\left(1-e^{-\sigma} \frac{x_{d}+\beta_{U}}{\beta_{U}}\right) .
$$

Similarly we get for $t \in\left[t_{d}-\sigma, t_{d}\right]$ that

$$
\begin{aligned}
x(t) & =-\beta_{U}+a+\left(x\left(t_{d}-\sigma\right)-\left(-\beta_{U}+a\right)\right) e^{-\left(t-\left(t_{d}-\sigma\right)\right)} \\
& =-\beta_{U}+a+\left(y\left(t_{d}-\sigma\right)-\left(-\beta_{U}+a\right)\right) e^{-\left(t-\left(t_{d}-\sigma\right)\right)} \\
& =y(t)+a\left(1-e^{-\left(t-\left(t_{d}-\sigma\right)\right)}\right) \geq y(t) \geq x_{d}
\end{aligned}
$$

which shows that $x$ is monotone and above $x_{d}$ on this interval. Using

$$
x\left(t_{d}-\sigma\right)=y\left(t_{d}-\sigma\right)<0
$$

and monotonicity we conclude that we have $x(t)<0$ on $\left[t_{d}-\sigma, t_{d}\right]$ if and only if $x\left(t_{d}\right)<0$. Also,

$$
x\left(t_{d}\right)=y\left(t_{d}\right)+a\left(1-e^{-\left(t_{d}-\left(t_{d}-\sigma\right)\right)}\right)=x_{d}+a\left(1-e^{-\sigma}\right)
$$

which gives $x\left(t_{d}\right)<0$ if and only if

$$
\begin{equation*}
x_{d}+a\left(1-e^{-\sigma}\right)<0 \tag{7.5}
\end{equation*}
$$

We shall come back to this later, and turn to

$$
\begin{aligned}
x\left(z_{1}+\tau\right) & =-\beta_{U}+\left(x\left(t_{d}\right)+\beta_{U}\right) e^{-\left(z_{1}+\tau-t_{d}\right)} \\
& =-\beta_{U}+\left(y\left(t_{d}\right)+a\left(1-e^{-\sigma}\right)+\beta_{U}\right) e^{-\left(z_{1}+\tau-t_{d}\right)} \\
& =y\left(z_{1}+\tau\right)+a\left(1-e^{-\sigma}\right) e^{-\left(z_{1}+\tau-t_{d}\right)} \\
& =\underline{x}+a\left(1-e^{-\sigma}\right) e^{-\left(z_{1}+\tau-t_{d}\right)}>\underline{x} .
\end{aligned}
$$

It follows that there is a unique $a=a_{d}>0$ so that

$$
x\left(z_{1}+\tau\right)=x_{d}>\underline{x},
$$

namely,

$$
\begin{equation*}
a_{d}=\frac{x_{d}-\underline{x}}{\left(1-e^{-\sigma}\right) e^{-\left(z_{1}+\tau-t_{d}\right)}}=\frac{\left(x_{d}-\underline{x}\right) e^{\tau}\left(x_{d}+\beta_{U}\right)}{\beta_{U}\left(1-e^{-\sigma}\right)} . \tag{7.6}
\end{equation*}
$$

We would like to have $x(t)<0$ on $\left(z_{1}, z_{1}+\tau\right]$. This follows from $x\left(z_{1}+\tau\right)=x_{d}<0$ in combination with monotonicity provided we have $x\left(t_{d}\right)<0$, which was characterized by (7.5). So we ask under which conditions $a=a_{d}$ satisfies (7.5), or equivalently,

$$
\frac{\left(x_{d}-\underline{x}\right) e^{\tau}\left(x_{d}+\beta_{U}\right)}{\beta_{U}}=a_{d}\left(1-e^{-\sigma}\right)<-x_{d},
$$

which means

$$
\begin{equation*}
\left(x_{d}-\underline{x}\right) e^{\tau}\left(x_{d}+\beta_{U}\right)<-\beta_{U} x_{d} . \tag{7.7}
\end{equation*}
$$

Recall that $\underline{x}$ depends on $\tau$ and on $\beta_{U}$; given $\tau$ and $\beta_{U}$ the preceding inequality holds provided we consider $x_{d} \in(\underline{x}, 0)$ close enough to $\underline{x}$.

Assume from now on that $x_{d}$ is chosen so that (7.7) holds. If we follow the solution $x$ which started from $\phi$ further then we see from Eq. (2.5) and because of $x(t)<0$ on $\left(z_{1}, z_{1}+\tau\right]$ that $x$ begins to increase after $t=z_{1}+\tau$, has a first zero $z_{2}=z_{2}(\phi)>$ $z_{1}+\tau$, and coincides on $\left[z_{2}, z_{2}+\tau\right]$ with a translate of the periodic solution $\tilde{x}$ of Eq. (2.5) which has a zero $t=z_{2}$ and then increases to the value $\bar{x}$ at $t=z_{2}+\tau$. Notice that if we take this segment of $\tilde{x}$ as the initial value $\phi$ for the function $x$ then
$x_{z_{2}+\tau}=\phi=x_{0}$, and iteration of the whole procedure yields a periodic solution $x$. The inequalities $x(t)<0$ for $z_{1}<t \leq z_{1}+\tau$ and $y\left(z_{1}+\tau\right)=\underline{x}<x_{d}=x\left(z_{1}+\tau\right)$ in combination with Eq. (2.5) for $z_{1}+\tau \leq t$ imply that after the time $t=z_{1}+\tau$ the function $x$ reaches the zero level from below before $y$ does so, hence the period $z_{2}(\phi)+\tau$ of $x$ is shorter than the minimal period $\tilde{T}$ of $\tilde{x}$.

Remark 7.1 1. Crucial in the concept of the therapy plan is that the time $t_{d}$ at which the number $y(t)$ of blood cells in case of no medication would fall to the critical level $x_{d}$ is large enough for medication in the future (at some $t_{M}>0$ ) to become effective (during the time interval $\left[t_{M}+\tau, t_{M}+\tau+\sigma\right]$ ) before the time $t=t_{d}$. Necessary for this is

$$
0<t_{d}-\tau, \quad \text { or } \tau<t_{d}
$$

the stronger condition (7.1) which we used in the exposition above can be relaxed.
2. A practically useful version of the therapy plan would require an extension to more realistic model equations, probably with continuous production functions, in order to get reliable predictions of $t_{d}=t_{d}(\phi)$ for a large set of initial conditions, as they may arise from monitoring the number of blood cells in a patient.

## 8 Discussion

The computation of the response of the periodic solution of (2.5) to a perturbation of the form defined in Sect. 4 is complicated as we have seen in Sect. 5, and the response of the perturbed cycle length can be quite varied as shown in Sect. 6. However, our calculations have shown that in a variety of situations that are dependent on the timing of the pulse, the pulse has had no effect on subsequent minima of the model solution. This is in sharp contrast to what is noted in clinical situations where G-CSF is employed and in which both an amelioration as well as a worsening of neutropenia is clearly documented in response to the G-CSF, and the nature of the response is dependent on when G-CSF is given as well as the dosage. Although it seems to be a curious anomaly that a cytokine like G-CSF, which inhibits apoptosis, should actually make neutropenia worse, in this section we will show that in case the nonlinearity (production function) $f=f_{*}$ in Eq. (1.1) is piecewise constant with three values (see Fig. 17), say,

$$
f_{*}(\xi)=\beta_{L}>0 \text { for } \xi<0, f_{*}(\xi)=-\beta_{U}<0 \text { for } 0 \leq \xi<\xi_{*}
$$

and

$$
f_{*}(\xi)=-\beta_{*}<-\beta_{U} \text { for } \xi_{*} \leq \xi,
$$

then a pulse as in case RPRP can result in a subsequent minimum of the solution which is lower than the minimum of a periodic solution of the equation without a pulse. This will happen if the pulse pushes the state variable to high values where the negative feedback is so strong that after the delay time it drives the state down to very low values.


Fig. 17 The graph of the function $f_{*}$

In order to see that this actually happens for a suitable range of parameters, assume (in part for simplicity) that $\xi_{*}=\bar{x}$, and define solutions of the equation

$$
\begin{equation*}
x^{\prime}(t)=-x(t)+f_{*}(x(t-\tau)), \tag{8.1}
\end{equation*}
$$

i.e., of Eq. (1.1) with $f=f_{*}$, as in Sect. 3. Then our former periodic solution $\tilde{x}$ of Eq. (2.5) will also be a solution of Eq. (8.1). Consider a pulse which begins at $\tilde{z}_{1}$ and ends at $\tilde{z}_{1}+\tau$, that is, consider the function $x_{*}: \mathbb{R} \rightarrow \mathbb{R}$ which coincides with $\tilde{x}$ on $\left(-\infty, \tilde{z}_{1}\right]$, is given by

$$
x^{\prime}(t)=-x(t)+\beta_{L}+a \text { for } \tilde{z}_{1} \leq t \leq \tilde{z}_{1}+\tau
$$

and by Eq. (8.1) for $t \geq \tilde{z}_{1}+\tau$. We have

$$
x_{*}\left(\tilde{z}_{1}+\tau\right)=\left(\beta_{L}+a\right)\left(1-e^{-\tau}\right)>\beta_{L}\left(1-e^{-\tau}\right)=\tilde{x}\left(\tilde{z}_{1}+\tau\right)=\bar{x}=\xi_{*} .
$$

Incidentally, notice that $x_{*}\left(\tilde{z}_{1}+\tau\right) \rightarrow \beta_{L}+a$ as $\tau \rightarrow \infty$.
The first time $t_{*}>\tilde{z}_{1}$ at which $x_{*}$ crosses the level $\xi_{*}$ from below is given by

$$
\xi_{*}=\left(\beta_{L}+a\right)\left(1-e^{-\left(t_{*}-\tilde{z}_{1}\right)}\right),
$$

or equivalently,

$$
\beta_{L}\left(1-e^{-\tau}\right)=\left(\beta_{L}+a\right)\left(1-e^{-\left(t_{*}-\tilde{z}_{1}\right)}\right),
$$

hence

$$
e^{\tilde{z}_{1}-t_{*}}=\frac{a+\beta_{L} e^{-\tau}}{a+\beta_{L}}
$$

Notice here that

$$
e^{\tilde{z}_{1}-t_{*}} \rightarrow \frac{a}{a+\beta_{L}} \quad \text { as } \tau \rightarrow \infty .
$$

Next,

$$
\begin{aligned}
x_{*}\left(t_{*}+\tau\right) & =-\beta_{U}+\left(x_{*}\left(\tilde{z}_{1}+\tau\right)+\beta_{U}\right) e^{-\left(t_{*}+\tau-\left(\tilde{z}_{1}+\tau\right)\right)} \\
& =-\beta_{U}+\left(x_{*}\left(\tilde{z}_{1}+\tau\right)+\beta_{U}\right) e^{\tilde{z}_{1}-t_{*}}
\end{aligned}
$$

We observe that $x_{*}\left(t_{*}+\tau\right)$ has a limit as $\tau \rightarrow \infty$. It follows that

$$
\begin{aligned}
x_{*}\left(\tilde{z}_{1}+\tau+\tau\right) & =-\beta_{*}+\left(x_{*}\left(t_{*}+\tau\right)+\beta_{*}\right) e^{-\left(\tilde{z}_{1}+2 \tau-\left(t_{*}+\tau\right)\right)} \\
& =-\beta_{*}+\left(x_{*}\left(t_{*}+\tau\right)+\beta_{*}\right) e^{t_{*}-\tilde{z}_{1}} e^{-\tau}
\end{aligned}
$$

converges to $-\beta_{*}<-\beta_{U}<\underline{x}=\min \tilde{x}(\mathbb{R})$ as $\tau \rightarrow \infty$. So, given $\xi_{*}>0$ and $\beta_{L}>0>-\beta_{U}>-\beta_{*}$ and $a>0$ there exists $\tau_{0}>0$ so that for all $\tau \geq \tau_{0}$ the solution $x_{*}$ assumes values strictly less than $\underline{x}=\min \tilde{x}(\mathbb{R})$.

Our investigations in this paper have been confined to an examination of the response of the limit cycle solution of (2.5) to a single perturbation. However, in many situations of interest biologically (and certainly for the clinical questions that motivated this study) one is interested in the limiting behaviour of the limit cycle in response to periodic perturbations, c.f Winfree (1980), Guevara and Glass (1982), Glass and Winfree (1984), Krogh-Madsen et al. (2004) and Bodnar et al. (2013b) for representative examples. However, considerations of the response to the system we have studied to periodic perturbation is quite beyond the scope of this study as it would entail the development of completely different techniques than the ones that have proved so successful in the study of the response to single perturbations. We are of the opinion that deriving the phase response curve in the face of periodically delivered pulses will only be possible, in general, for certain limiting cases of the pulse parameters, namely $\sigma \simeq 0$ and, possibly, small values of the amplitude $a$. It is possible that techniques such as those employed in Kotani et al. (2012) may be useful in this regard.

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## 9 Appendix: Proofs of the results from Sects. 3 and 4

Proof of Proposition 3.1 1. We begin with continuity of the time- $\tau$-map

$$
S(\tau, \cdot): Z \ni \phi \mapsto x_{\tau}^{\phi} \in Z
$$

Observe that for $0 \leq t \leq \tau$,

$$
x^{\phi}(t)=e^{-t} \phi(0)+\int_{0}^{t} e^{-(t-s)} f(\phi(s-\tau)) d s
$$

For $\psi$ and $\phi$ in $Z$ and $0 \leq t \leq \tau$ we have

$$
\left|x^{\psi}(t)-x^{\phi}(t)\right| \leq|\phi(0)-\psi(0)|+\int_{-\tau}^{0}|f(\psi(s))-f(\phi(s))| d s
$$

where the integrand is nonzero only on the set

$$
N(\psi, \phi)=\{t \in[-\tau, 0]: \operatorname{sign}(\psi(t)) \neq \operatorname{sign}(\phi(t))\} .
$$

It follows that

$$
\left|x_{\tau}^{\psi}-x_{\tau}^{\phi}\right|_{C} \leq|\psi(0)-\phi(0)|+\beta \lambda(N(\psi, \phi))
$$

with the Lebesgue measure $\lambda$ and a positive constant $\beta$. It is easy to see that

$$
\lim _{Z \ni \psi \rightarrow \phi \in Z} \lambda(N(\psi, \phi))=0 .
$$

(Proof of this in case $\phi \in Z$ has zeros $z_{1}<z_{2}<\cdots<z_{J}$. Let $\epsilon>0$ be given. The complement of the set

$$
\bigcup_{j=1}^{J}\left(z_{j}-\frac{\epsilon}{2 J}, z_{j}+\frac{\epsilon}{2 J}\right)
$$

in $[-\tau, 0]$ is the finite union of compact intervals on each of which $\phi$ is either strictly positive, or strictly negative. There exists $\delta>0$ so that for every $\psi \in Z$ with $|\psi-\phi|<\delta$ the signs of $\psi(t)$ and $\phi(t)$ coincide on each of the compact intervals. This yields

$$
\left.\lambda(N(\psi, \phi)) \leq \sum_{j=1}^{J} 2 \frac{\epsilon}{2 J}=\epsilon .\right)
$$

Then it follows easily that

$$
\lim _{Z \ni \psi \rightarrow \phi \in Z}(S(\tau, \psi)-S(\tau, \phi))=\lim _{Z \ni \psi \rightarrow \phi \in Z}\left(x_{\tau}^{\psi}-x_{\tau}^{\phi}\right)=0 .
$$

2. Iterating we find that for every integer $n>0$ the time- $n \tau$-map $S(n \tau, \cdot)$ is continuous. Having this we obtain continuous dependence on initial data in the sense that for every $t \geq 0$ and $\phi \in Z$,

$$
\lim _{Z \ni \psi \rightarrow \phi \in Z} \max _{-\tau \leq s \leq t}\left|x^{\psi}(s)-x^{\phi}(s)\right|=0 .
$$

Finally, the continuity of $S$ at $(t, \phi) \in[0, \infty) \times Z$ follows by means of the estimate

$$
\begin{aligned}
|S(s, \psi)-S(t, \phi)| \leq & |S(s, \psi)-S(s, \phi)|+|S(s, \phi)-S(t, \phi)| \\
\leq & \max _{-\tau \leq v \leq t+1}\left|x^{\psi}(v)-x^{\phi}(v)\right| \\
& +\max _{-\tau \leq w \leq 0}\left|x^{\phi}(s+w)-x^{\phi}(t+w)\right|
\end{aligned}
$$

for $0 \leq s \leq t+1$ and $\psi \in Z$ from continuous dependence on initial data as before in combination with the uniform continuity of $x^{\phi}$ on $[-\tau, t+1]$.

Proof of Proposition 4.2 By definition the value of the cycle length map at $\Delta$ is $T(\Delta)=z-\tilde{z}_{J}=z-z_{\Delta, J}$ where $z$ is the smallest zero of $x^{(\Delta)}$ in $\left(\tilde{z}_{J}, \infty\right)$ such that $x^{(\Delta)}(z+t)=\tilde{x}\left(\tilde{z}_{J}+t\right)$ for all $t \geq 0$. We have

$$
\operatorname{sign}\left(\tilde{x}\left(\tilde{z}_{J}+t\right)\right)=\operatorname{sign}\left(x^{(\Delta)}\left(z_{\Delta, J}+t\right)\right)=-\operatorname{sign}\left(x^{(\Delta)}\left(z_{\Delta, J+1}+t\right)\right)
$$

for $0<t \leq \tau$ since $x^{(\Delta)}$ changes sign at each zero. We infer that $z>z_{\Delta, J+1}$. Notice that the definition of $J=j_{\Delta}$ implies $\Delta<z_{\Delta, J+1}$. Hence the next zero $z_{\Delta, J+2}$ satisfies $z_{\Delta, J+2}>z_{\Delta, J+1}+\tau>\Delta+\tau \geq \Delta+\sigma$. Therefore on $\left(z_{\Delta, J+2}, \infty\right)$ the function $x^{(\Delta)}$ is given by Eq. (2.5), and satisfies

$$
\operatorname{sign}\left(x^{(\Delta)}\left(z_{\Delta, J+2}+t\right)\right)=-\operatorname{sign}\left(x^{(\Delta)}\left(z_{\Delta, J+1}+t\right)\right)=\operatorname{sign}\left(\tilde{x}\left(\tilde{z}_{J}+t\right)\right)
$$

for $0<t \leq \tau$. This yields $x^{(\Delta)}\left(z_{\Delta, J+2}+t\right)=\tilde{x}\left(\tilde{z}_{J}+t\right)$ for all $t \geq 0$.
Proposition 9.1 For $\Delta=\tilde{z}_{J}, J=1$ or $J=2$, we have $T(\Delta)=z_{\Delta, J+1}-z_{\Delta, J-1}$.
Proof From $\Delta=\tilde{z}_{J}$ we obtain $\Delta+\sigma \leq \Delta+\tau<z_{\Delta, J+1}$. This implies that for $t \geq z_{\Delta, J+1}$ the function $x^{(\Delta)}$ satisfies Eq. (2.5). Using this and the fact that $x^{(\Delta)}$ and $\tilde{x}$ change sign at $\tilde{z}_{J-1}=z_{\Delta, J-1}$ and at $\tilde{z}_{J}=z_{\Delta, J}$ respectively we infer that for all $t \geq z_{\Delta, J+1}$ we have $x^{(\Delta)}\left(z_{\Delta, J+1}+t\right)=\tilde{x}\left(\tilde{z}_{J-1}+t\right)$. It follows that

$$
z_{\Delta, J+2}=z_{\Delta, J+1}+\left(\tilde{z}_{J}-\tilde{z}_{J-1}\right)
$$

Combining this with Proposition 4.2 we find

$$
T(\Delta)=z_{\Delta, J+2}-\tilde{z}_{J}=z_{\Delta, J+1}-\tilde{z}_{J-1}=z_{\Delta, J+1}-z_{\Delta, J-1}
$$

Proof of Corollary 4.2 Let $\Delta_{0} \in[0, \tilde{T})$ be given and set $J=j\left(\Delta_{0}\right)$. Then $\Delta_{0}<$ $z_{\Delta_{0}, J+1}$. Corollary 4.1 yields a neighbourhood $N$ of $\Delta_{0}$ in $[0, \tilde{T})$ such that for all $\Delta \in N$ we have $\Delta<z_{\Delta, J+1}$.

1. The case $\tilde{z}_{J}<\Delta_{0}$. Then by Corollary 4.1, $\tilde{z}_{J}<\Delta<z_{\Delta, J+1}$ for all $\Delta$ in a neighbourhood $V \subset N$ of $\Delta_{0}$ in $[0, \tilde{T})$. For $\Delta \in V$ we get $j(\Delta)=J$, hence $T(\Delta)=z_{\Delta, J+2}-\tilde{z}_{J}$, and Corollary 4.1 yields continuity at $\Delta_{0}$.
2. The case $\tilde{z}_{J}=\Delta_{0}$. There is a neighbourhood $U \subset N$ of $\Delta_{0}$ in [0, $\left.\tilde{T}\right)$ with $\tilde{z}_{J-1}<\Delta$ for all $\Delta \in U$. For all $\Delta \in U$ with $\Delta<\tilde{z}_{J}$ this yields $j(\Delta)=J-1$ and $T(\Delta)=z_{\Delta, J-1+2}-\tilde{z}_{J-1}$. At $\Delta=\Delta_{0}$ we have

$$
\begin{aligned}
T(\Delta)\left(\Delta_{0}\right) & =z_{\Delta_{0}, J+2}-\tilde{z}_{J} \\
& =z_{\Delta_{0}, J+1}-z_{\Delta_{0}, J-1} \quad \text { (see Proposition 9.1) } \\
& =z_{\Delta_{0}, J+1}-\tilde{z}_{J-1}
\end{aligned}
$$

The continuity of the map $\Delta \mapsto z_{\Delta, J+1}$ due to Corollary 4.1 now shows that the restriction of the cycle length map to the set $\left[0, \tilde{z}_{J}\right] \cap U$ is continuous. For $\tilde{z}_{J} \leq \Delta \in U \subset N$ we have $\tilde{z}_{J} \leq \Delta<z_{\Delta, J+1}$, hence $j(\Delta)=J$, and thereby $T(\Delta)=z_{\Delta, J+2}-\tilde{z}_{J}$. The continuity of the map $\Delta \mapsto z_{\Delta, J+2}$ due to Corollary 4.1 shows that the restriction of the cycle length map to the set $U \cap\left[\tilde{z}_{J}, \tilde{T}\right)$ is continuous. As both restrictions coincide at $\tilde{z}_{J}=\Delta_{0}$ we obtain continuity of the cycle length map at $\Delta_{0}$.

Proof of Proposition 4.3 1. Let $\Delta_{0} \in[0, \tilde{T})$ be given. Set $J=j(\Delta)$. Then $\tilde{z}_{J} \leq$ $\Delta_{0}<z_{\Delta_{0}, J+1}$. Using Corollary 4.1 we find a neighbourhood $N$ of $\Delta_{0}$ in [0, $\left.\tilde{T}\right)$ such that for every $\Delta \in N$ we have

$$
\tilde{z}_{J-1}<\Delta<z_{\Delta, J+1} .
$$

In the following we show continuity of the map $[0, \tilde{T}) \ni \Delta \mapsto \bar{x}_{\Delta} \in \mathbb{R}$. The proof for the other map is analogous.
2. For $\Delta \in N \cap\left[\tilde{z}_{J}, \infty\right)$ we have $J=j(\Delta)$, hence

$$
\bar{x}_{\Delta}=\max _{\tilde{z}_{J \leq t \leq z_{\Delta}, J+2}} x^{(\Delta)}(t) .
$$

Using this, the uniform continuity of the map

$$
[0, \tilde{T}) \times[0, \infty) \ni(\Delta, t) \mapsto x^{(\Delta)}(t) \in \mathbb{R}
$$

on compact sets, and the continuity of the map $\Delta \mapsto z_{\Delta, J+2}$ (see Corollary 4.1) one can easily show that the map

$$
N \cap\left[\tilde{z}_{J}, \infty\right) \ni \Delta \mapsto \bar{x}_{\Delta} \in \mathbb{R}
$$

is continuous.
Similarly we have for $\Delta \in N \cap\left(-\infty, \tilde{z}_{J}\right)$ that $J-1=j(\Delta)$, hence

$$
\bar{x}_{\Delta}=\max _{\tilde{z}_{J-1} \leq t \leq z_{\Delta, J+1}} x^{(\Delta)}(t)
$$

As before one can then easily show that the map

$$
N \cap\left(-\infty, \tilde{z}_{J}\right) \ni \Delta \mapsto \bar{x}_{\Delta} \in \mathbb{R}
$$

is continuous.
3. It remains to prove that in case $\Delta_{0}=\tilde{z}_{J}$ (where $N \cap\left(-\infty, \tilde{z}_{J}\right) \neq \emptyset$ ) we have $\bar{x}_{\Delta} \rightarrow \bar{x}_{\Delta_{0}}$ as $\Delta \nearrow \Delta_{0}$.
(a) The case $\Delta_{0}=\tilde{z}_{J}$ and $\tilde{x}^{\prime}\left(\tilde{z}_{J}\right)<0$. Using the fact that $x^{\left(\Delta_{0}\right)}$ changes sign at each zero we obtain $x^{\left(\Delta_{0}\right)}(t) \leq 0$ on $\left[\tilde{z}_{j}, z_{\Delta_{0}, J+1}\right]=\left[\Delta_{0}, z_{\Delta_{0}, J+1}\right]$, and

$$
\begin{aligned}
\bar{x}_{\Delta_{0}} & =\max _{\tilde{z}_{J \leq t \leq z_{\Delta_{0}, J+2}} x^{\left(\Delta_{0}\right)}(t)} \\
& =\max _{z_{\Delta_{0}, J+1 \leq t \leq z_{\Delta_{0}, J+2}}} x^{\left(\Delta_{0}\right)}(t) \\
& =\bar{x}
\end{aligned}
$$

where the last equation holds because $\Delta_{0}=\tilde{z}_{J}$ implies $\Delta_{0}+\sigma<\Delta_{0}+\tau<$ $z_{\Delta_{0}, J+1}$ and thereby $x^{\left(\Delta_{0}\right)}\left(z_{\Delta_{0}, J+1}+t\right)=\tilde{x}\left(\tilde{z}_{J-1}+t\right)$ for all $t \geq 0$.
Using continuity as in Part 2 above we find a neighbourhood $U \subset N$ of $\Delta_{0}$ in $[0, \tilde{T})$ such that for each $\Delta \in U$ we have

$$
x^{(\Delta)}(t)<\frac{1}{2} \bar{x} \quad \text { on } \quad\left[\Delta, z_{\Delta, J+1}\right]
$$

and $\tilde{z}_{J-1}+\tau<\Delta$. For $\Delta \in U$ with $\Delta<\tilde{z}_{J}$, we have $J-1=j(\Delta)$, and the preceding inequality yields

$$
x^{(\Delta)}\left(\tilde{z}_{J-1}+\tau\right)=\tilde{x}\left(\tilde{z}_{J-1}+\tau\right)=\bar{x} \quad\left(>\frac{1}{2} \bar{x}\right) .
$$

It follows that

$$
\begin{aligned}
\bar{x}_{\Delta} & =\max _{\tilde{z}_{J-1 \leq t \leq z_{\Delta, J+1}}} x^{(\Delta)}(t) \\
& =\max _{\tilde{z}_{J-1 \leq t \leq \Delta}} x^{(\Delta)}(t) \\
& =\bar{x},
\end{aligned}
$$

so the map $U \cap\left(-\infty, \tilde{z}_{J}\right) \ni \Delta \rightarrow \bar{x}_{\Delta} \in \mathbb{R}$ is constant with value $\bar{x}=\bar{x}_{\Delta_{0}}$.
(b) The case $\Delta_{0}=\tilde{z}_{J}$ and $\tilde{x}^{\prime}\left(\tilde{z}_{J}\right)>0$. Then $x^{\left(\Delta_{0}\right)}$ is negative on $\left(\tilde{z}_{J-1}, \tilde{z}_{J}\right)$, positive on ( $\left.\tilde{z}_{J}, z_{\Delta_{0}, J+1}\right)$ and negative on $\left(z_{\Delta_{0}, J+1}, z_{\Delta_{0}, J+2}\right)$, and

$$
\bar{x}_{\Delta_{0}}=\max _{\tilde{z}_{J \leq t \leq z_{\Delta_{0}}, J+2}} x^{(\Delta)}(t)=\max _{\tilde{z}_{J \leq t \leq z_{\Delta_{0}}, J+1}} x^{(\Delta)}(t)>0 .
$$

Choose $t_{0} \in\left(\tilde{z}_{J}, z_{\Delta_{0}, J+1}\right)$ with

$$
x^{(\Delta)}\left(t_{0}\right)=\bar{x}_{\Delta_{0}}>0 .
$$

By continuity there exists a neighbourhood $V \subset N$ of $\Delta_{0}$ in $[0, \tilde{T})$ such that for every $\Delta \in V$ we have

$$
\begin{aligned}
x^{(\Delta)}(t) & <\frac{1}{2} \bar{x}_{\Delta_{0}} \quad \text { on } \quad\left[\tilde{z}_{J-1}, z_{\Delta, J}\right], \\
z_{\Delta, J} & <t_{0}<z_{\Delta, J+1}, \\
\frac{1}{2} \bar{x}_{\Delta_{0}} & <x^{(\Delta)}\left(t_{0}\right) .
\end{aligned}
$$

For $\Delta \in V$ with $\Delta<\tilde{z}_{J}$ we have $j(\Delta)=J-1$, and we conclude that

$$
\begin{aligned}
\bar{x}_{\Delta} & =\max _{\tilde{z}_{J-1} \leq t \leq z_{\Delta, J+1}} x^{(\Delta)}(t) \\
& =\max _{z_{\Delta, J} \leq t \leq z_{\Delta, J+1}} x^{(\Delta)}(t)
\end{aligned}
$$

where by continuity the last term converges to

$$
\max _{z_{\Delta_{0}, J} \leq t \leq z_{\Delta_{0}, J+1}} x^{\left(\Delta_{0}\right)}(t)=\bar{x}_{\Delta_{0}}
$$

as $V \ni \Delta \nearrow \Delta_{0}$.

## Proofs of the results from Sect. 5

Proof of Proposition 5.1 First we show that $x^{(\Delta)}$ has a first positive zero $z_{\Delta, 1}<\tilde{z}_{1}$. We have $x^{(\Delta)}(t)<0$ for $t \in(-\tau, \Delta+\sigma)$. For $t \in[\Delta+\sigma, \infty), x^{(\Delta)}(t)$ is given by (5.4) as long as $x^{(\Delta)}(t-\tau)<0$. Compute

$$
z_{\Delta, 1}=\Delta+\sigma+\ln \frac{\beta_{L}-x^{(\Delta)}(\Delta+\sigma)}{\beta_{L}} \geq \Delta+\sigma
$$

from the condition $x^{(\Delta)}\left(z_{\Delta, 1}\right)=0$. Similarly since $\tilde{x}\left(\tilde{z}_{1}\right)=0$ we obtain

$$
\tilde{z}_{1}=\Delta+\sigma+\ln \frac{\beta_{L}-\tilde{x}(\Delta+\sigma)}{\beta_{L}} .
$$

Since $\tilde{x}(\Delta+\sigma)<x^{(\Delta)}(\Delta+\sigma)$ we have $z_{\Delta, 1}<\tilde{z}_{1}$.
The largest zero of $\tilde{x}$ on $(-\infty, \Delta]$ is $\tilde{z}_{0}=-\tau$. Hence the minimal value of $x^{(\Delta)}$ on $\left[-\tau, z_{\Delta, 1}\right]$ is equal to $\underline{x}$. On the interval $\left[z_{\Delta, 1}, z_{\Delta, 1}+\tau\right], x^{(\Delta)}(t)$ is given by (5.4). This yields $x^{(\Delta)}\left(t+z_{\Delta, 1}\right)=\tilde{x}\left(t+\tilde{z}_{1}\right)$ for all $t \geq 0$. It follows that $\bar{x}_{\Delta}=\bar{x}$ and
$\underline{x}_{\Delta}=\underline{x}$ and

$$
\begin{aligned}
T(\Delta) & =\left(z_{\Delta, 1}+\left(\tilde{z}_{2}-\tilde{z}_{1}\right)\right)-\tilde{z}_{0}=z_{\Delta, 1}-\tilde{z}_{1}+\tilde{T} \\
& =\tilde{T}+\ln \frac{\beta_{L}-x^{(\Delta)}(\Delta+\sigma)}{\beta_{L}}-\ln \frac{\beta_{L}-\tilde{x}(\Delta+\sigma)}{\beta_{L}} \\
& =\tilde{T}+\ln \frac{\beta_{L}-x^{(\Delta)}(\Delta+\sigma)}{\beta_{L}-\tilde{x}(\Delta+\sigma)}
\end{aligned}
$$

The formula (5.5) yields

$$
x^{(\Delta)}(\Delta+\sigma)-\beta_{L}=-\beta_{L} e^{\tilde{z}_{1}-\Delta-\sigma}+a\left(1-e^{-\sigma}\right)
$$

and with

$$
\tilde{x}(\Delta+\sigma)-\beta_{L}=\left(\underline{x}-\beta_{L}\right) e^{-(\Delta+\sigma)}=-\beta_{L} e^{\tilde{z}_{1}-(\Delta+\sigma)}
$$

we find

$$
\begin{aligned}
T(\Delta) & =\tilde{T}+\ln \frac{\beta_{L} e^{\tilde{z}_{1}-\Delta-\sigma}-a\left(1-e^{-\sigma}\right)}{\beta_{L} e^{\tilde{z}_{1}-(\Delta+\sigma)}} \\
& =\tilde{T}+\ln \left(1-\frac{a\left(1-e^{-\sigma}\right)}{\beta_{L}} e^{\Delta+\sigma-\tilde{z}_{1}}\right)
\end{aligned}
$$

Thus, $T(\Delta)<\tilde{T}$, and the restriction of $T(\Delta)$ to $I_{R N R N}$ is strictly decreasing.
From (5.4) it follows that $x^{(\Delta)}$ is strictly decreasing right after $t=\Delta+\sigma$ if and only if $x^{(\Delta)}(\Delta+\sigma)-\beta_{L}>0$. We have, by (5.5),

$$
x^{(\Delta)}(\Delta+\sigma)-\beta_{L}=-\beta_{L} e^{\tilde{z}_{1}-\Delta-\sigma}+a\left(1-e^{-\sigma}\right)
$$

Hence, $x^{(\Delta)}(\Delta+\sigma)-\beta_{L}>0$ if and only if

$$
e^{\Delta}>\frac{\beta_{L} e^{\tilde{z}_{1}-\sigma}}{a\left(1-e^{-\sigma}\right)}
$$

Let us define

$$
\begin{equation*}
\hat{\delta}_{1}=\ln \frac{\beta_{L} e^{\tilde{z}_{1}-\sigma}}{a\left(1-e^{-\sigma}\right)} \tag{9.1}
\end{equation*}
$$

We have

$$
\hat{\delta}_{1}=\tilde{z}_{1}-\sigma+\ln \frac{\beta_{L}}{a\left(1-e^{-\sigma}\right)}>\delta_{1}
$$

Thus we divide the case RNRP when $x^{(\Delta)}(\Delta+\sigma) \geq 0$ into the two subcases $0 \leq$ $x^{(\Delta)}(\Delta+\sigma) \leq \beta_{L}$ and $x^{(\Delta)}(\Delta+\sigma)>\beta_{L}$ and we consider these two subcases
separately as RNRP1 and RNRP2, see Figs. 3 and 4, respectively. The $\Delta$-intervals are of the form
$I_{R N R P 1}=\left[\max \left\{0, \delta_{1}\right\}, \tilde{z}_{1}\right) \cap\left[\delta_{1}, \hat{\delta}_{1}\right]$ and $I_{R N R P 2}=\left[\max \left\{0, \delta_{1}\right\}, \tilde{z}_{1}\right) \cap\left(\hat{\delta}_{1}, \infty\right)$.
Case RNRP1 The pulse parameters $(a, \Delta, \sigma)$ are such that $\Delta<\tilde{z}_{1}$ and $x^{(\Delta)}(\Delta+$ $\sigma) \in\left[0, \beta_{L}\right]$. Then $x^{(\Delta)}$ is increasing right after the pulse.
Case RNRP2 The pulse parameters $(a, \Delta, \sigma)$ are such that $\Delta<\tilde{z}_{1}$ and $x^{(\Delta)}(\Delta+$ $\sigma)>\beta_{L}$. Then $x^{(\Delta)}$ is strictly decreasing right after the pulse.

## Proof of Proposition 5.2 The equation

$$
0=x^{(\Delta)}\left(z_{\Delta, 1}\right)=\beta_{L}+a+\left(x^{(\Delta)}(\Delta)-\left(\beta_{L}+a\right)\right) e^{-\left(z_{\Delta, 1}-\Delta\right)}
$$

together with (5.1) yields

$$
\begin{align*}
\left(\beta_{L}+a\right) e^{z_{\Delta, 1}} & =\left(\beta_{L}+a-x^{(\Delta)}(\Delta)\right) e^{\Delta} \\
& =\beta_{L} e^{\tilde{z}_{1}}+a e^{\Delta} . \tag{9.2}
\end{align*}
$$

Use $\Delta<z_{\Delta, 1}$ and $a>0, \beta_{L}>0$ to obtain

$$
e^{z_{\Delta, 1}}-e^{\tilde{z}_{1}}=\frac{a}{\beta_{L}}\left(e^{\Delta}-e^{z_{\Delta, 1}}\right)<0
$$

and to conclude that $z_{\Delta, 1}<\tilde{z}_{1}$.
We have

$$
x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right)=\beta_{L}+\left(x^{(\Delta)}(\Delta+\sigma)-\beta_{L}\right) e^{-\left(z_{\Delta, 1}+\tau-(\Delta+\sigma)\right)}
$$

which by (5.5) can be rewritten as

$$
\begin{aligned}
x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right) & =\beta_{L}+\left(-\beta_{L} e^{\tilde{z}_{1}-\Delta-\sigma}+a\left(1-e^{-\sigma}\right)\right) e^{-\left(z_{\Delta, 1}+\tau-(\Delta+\sigma)\right)} \\
& =\beta_{L}-\left(\beta_{L} e^{\tilde{z}_{1}}+a e^{\Delta}\right) e^{-\left(z_{\Delta, 1}+\tau\right)}+a e^{-\left(z_{\Delta, 1}+\tau-(\Delta+\sigma)\right)} .
\end{aligned}
$$

From (9.2) it follows that

$$
\begin{equation*}
x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right)=\beta_{L}-\left(\beta_{L}+a\right) e^{-\tau}+a e^{-\tau} e^{\sigma+\Delta-z_{\Delta, 1}} \tag{9.3}
\end{equation*}
$$

and that the function $\left[\max \left\{0, \delta_{1}\right\}, \tilde{z}_{1}\right) \ni \Delta \mapsto e^{z_{\Delta, 1}-\Delta} \in \mathbb{R}$ is strictly decreasing, which shows that $x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right)$ is strictly increasing with respect to $\Delta \in$ $\left[\max \left\{0, \delta_{1}\right\}, \tilde{z}_{1}\right)$. Since $\Delta+\sigma \geq z_{\Delta, 1}$ and $a e^{-\tau}>0$, we get

$$
x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right) \geq \beta_{L}-\beta_{L} e^{-\tau}=\bar{x}
$$

Also the function $\left[\max \left\{0, \delta_{1}\right\}, \tilde{z}_{1}\right) \ni \Delta \mapsto x^{(\Delta)}(\Delta+\sigma) \in \mathbb{R}$ is increasing.

In subcase RNRP1 the function $x^{(\Delta)}$ is increasing on [ $\left.\Delta+\sigma, z_{\Delta, 1}+\tau\right]$ while in subcase RNRP2 it is decreasing on that interval. It follows that in subcase RNRP1,

$$
\max _{z_{\Delta, 1} \leq t \leq z_{\Delta, 1}+\tau} x^{(\Delta)}(t)=x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right) \geq \bar{x}
$$

while in subcase RNRP2,

$$
\max _{z_{\Delta, 1} \leq t \leq z_{\Delta, 1}+\tau} x^{(\Delta)}(t)=x^{(\Delta)}(\Delta+\sigma) \geq x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right) \geq \bar{x} .
$$

In both subcases $\max _{z_{\Delta, 1} \leq t \leq z_{\Delta, 1}+\tau} x^{(\Delta)}(t)$ is increasing with respect to $\Delta$.
Also in both subcases we have $0<x^{(\Delta)}(t)$ for $z_{\Delta, 1}<t \leq z_{\Delta, 1}+\tau$. It follows that after $t=z_{\Delta, 1}+\tau \geq \Delta+\sigma$ the function $x^{(\Delta)}$ is given by

$$
x^{(\Delta)}(t)=-\beta_{U}+\left(x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right)+\beta_{U}\right) e^{-\left(t-\left(z_{\Delta, 1}+\tau\right)\right)}
$$

as long as $x^{(\Delta)}(t-\tau)>0$. We obtain a first zero $z_{\Delta, 2}$ of $x^{(\Delta)}$ in $\left(z_{\Delta, 1}+\tau, \infty\right)$, and for all $t \geq 0, x^{(\Delta)}\left(z_{\Delta, 2}+t\right)=\tilde{x}\left(\tilde{z}_{0}+t\right)\left(\right.$ recall $\left.\tilde{z}_{0}=-\tau\right)$. Then

$$
T(\Delta)=z_{\Delta, 2}-\tilde{z}_{0}=z_{\Delta, 2}+\tau
$$

Moreover,

$$
\begin{aligned}
\bar{x}_{\Delta} & =\max _{-\tau \leq t \leq z_{\Delta, 2}} x^{(\Delta)}(t)=\max _{z_{\Delta, 1} \leq t \leq z_{\Delta, 1}+\tau} x^{(\Delta)}(t) \\
& =\max \left\{x^{(\Delta)}(\Delta+\sigma), x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right)\right\} \\
& \geq \bar{x}
\end{aligned}
$$

is increasing with respect to $\Delta$ in both subcases.
Now the equation

$$
0=x^{(\Delta)}\left(z_{\Delta, 2}\right)=-\beta_{U}+\left(x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right)+\beta_{U}\right) e^{z_{\Delta, 1}+\tau-z_{\Delta, 2}}
$$

yields

$$
\beta_{U} e^{z_{\Delta, 2}}=\left(x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right)+\beta_{U}\right) e^{z_{\Delta, 1}+\tau}
$$

Also, from (9.2) and (9.3) we obtain

$$
\beta_{U} e^{z_{\Delta, 2}}=\left(\beta_{L}+\beta_{U}\right) e^{z_{\Delta, 1}+\tau}-\beta_{L} e^{\tilde{z}_{1}}+a\left(1-e^{-\sigma}\right) e^{\Delta+\sigma} .
$$

Since

$$
\beta_{U}+\beta_{L}-\beta_{L} e^{-\tau}=\beta_{U}+\bar{x}=\beta_{U} e^{\tilde{z}_{2}-t_{\max }}
$$

and $\tilde{z}_{2}-t_{\text {max }}+\tau=\tilde{z}_{2}-\tilde{z}_{1}$, we arrive at

$$
\begin{equation*}
\beta_{L} e^{\tilde{z}_{1}}+\beta_{U} e^{\tilde{z}_{2}}=\left(\beta_{L}+\beta_{U}\right) e^{\tau+\tilde{z}_{1}} \tag{9.4}
\end{equation*}
$$

Thus

$$
\begin{aligned}
\beta_{U} e^{z_{\Delta, 2}} & =\beta_{U} e^{\tilde{z}_{2}}+\left(\beta_{L}+\beta_{U}\right) e^{\tau}\left(e^{z_{\Delta, 1}}-e^{\tilde{z}_{1}}\right)+a\left(e^{\sigma}-1\right) e^{\Delta} \\
& =\beta_{U} e^{\tilde{z}_{2}}+\frac{a\left(\beta_{L}+\beta_{U}\right) e^{\tau}}{\beta_{L}+a}\left(e^{\Delta}-e^{\tilde{z}_{1}}\right)+a\left(e^{\sigma}-1\right) e^{\Delta} \quad(\text { with }(9.2)),
\end{aligned}
$$

which implies the formula for $T(\Delta)=z_{\Delta, 2}+\tau$, since $\tilde{z}_{2}+\tau=\tilde{T}$.
In the case RPRP we always have $x^{(\Delta)}(\Delta+\sigma)>0$, but in the following proof we need to distinguish between two cases

Case RPRP1 $0<x^{(\Delta)}(\Delta+\sigma) \leq \beta_{L}$, and
Case RPRP2 $x^{(\Delta)}(\Delta+\sigma)>\beta_{L}$.
Proof of Proposition 5.3 Using (5.4) for $\Delta+\sigma \leq t \leq z_{\Delta, 1}+\tau=\tilde{z}_{1}+\tau=t_{\max }$ and (5.5) we obtain

$$
\begin{aligned}
x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right) & =\beta_{L}+\left(x^{(\Delta)}(\Delta+\sigma)-\beta_{L}\right) e^{-\left(\tilde{z}_{1}+\tau-(\Delta+\sigma)\right)} \\
& =\beta_{L}-\beta_{L} e^{-\tau}+a\left(1-e^{-\sigma}\right) e^{-\left(\tilde{z}_{1}+\tau-(\Delta+\sigma)\right)} \\
& =\bar{x}+a\left(e^{\sigma}-1\right) e^{-\left(\tilde{z}_{1}+\tau\right)+\Delta}>\bar{x},
\end{aligned}
$$

and $x^{(\Delta)}\left(z_{\Delta, 1}+\tau\right)$ is increasing as a function of $\Delta \in I_{R P R P}$. Also $x^{(\Delta)}(\Delta+\sigma)$ is increasing as a function of $\Delta \in I_{R P R P}$. In case RPRP1 the function $x^{(\Delta)}$ is increasing on $\left[\tilde{z}_{1}, \tilde{z}_{1}+\tau\right]$ with $x^{(\Delta)}(t) \leq \beta_{L}$ on this interval, hence $\max _{\tilde{z}_{1} \leq t \leq \tilde{z}_{1}+\tau} x^{(\Delta)}(t)=$ $x^{(\Delta)}\left(\tilde{z}_{1}+\tau\right) \leq \beta_{L}$. In case RPRP2 the function $x^{(\Delta)}$ is increasing on $\left[\tilde{z}_{1}, \Delta+\right.$ $\sigma$ ] and decreasing on [ $\Delta+\sigma, \tilde{z}_{1}+\tau$ ], with $x^{(\Delta)}(t)>\beta_{L}$ on this interval, hence $\max _{\tilde{z}_{1} \leq t \leq \tilde{z}_{1}+\tau} x^{(\Delta)}(t)=x^{(\Delta)}(\Delta+\sigma)>\beta_{L}$. In both subcases,

$$
\max _{\tilde{z}_{1} \leq t \leq \tilde{z}_{1}+\tau} x^{(\Delta)}(t)=\max \left\{x^{(\Delta)}(\Delta+\sigma), x^{(\Delta)}\left(\tilde{z}_{1}+\tau\right)\right\} \geq x^{(\Delta)}\left(\tilde{z}_{1}+\tau\right)>\bar{x}
$$

and $\max _{\tilde{z}_{1} \leq t \leq \tilde{z}_{1}+\tau} x^{(\Delta)}(t)$ is increasing as a function of $\Delta \in I_{R P R P}$.
As $x^{(\Delta)}(t)>0$ on $\left(\tilde{z}_{1}, \tilde{z}_{1}+\tau\right]$ we have

$$
x^{(\Delta)}(t)=-\beta_{U}+\left(x^{(\Delta)}\left(\tilde{z}_{1}+\tau\right)+\beta_{U}\right) e^{-\left(t-\left(\tilde{z}_{1}+\tau\right)\right)}
$$

for $t \geq \tilde{z}_{1}+\tau$ as long as $x^{(\Delta)}(t-\tau)>0$. It follows that there is a smallest zero $z_{\Delta, 2}$ of $x^{(\bar{\Delta})}$ in $\left(\tilde{z}_{1}+\tau, \infty\right)$, and

$$
x^{(\Delta)}\left(z_{\Delta, 2}+t\right)=\tilde{x}\left(\tilde{z}_{0}+t\right) \text { for all } t \geq 0
$$

This yields

$$
T(\Delta)=\left(z_{\Delta, 2}+\left(\tilde{z}_{1}-\tilde{z}_{0}\right)\right)-\tilde{z}_{1}=z_{\Delta, 2}+\tau
$$

Moreover,

$$
\begin{aligned}
\bar{x}_{\Delta} & =\max _{\tilde{z}_{1} \leq t \leq z_{\Delta, 2}+\left(\tilde{z}_{1}-\tilde{z}_{0}\right)} x^{(\Delta)}(t) \\
& =\max _{\tilde{z}_{1} \leq t \leq \tilde{z}_{1}+\tau} x^{(\Delta)}(t) \\
& =\max \left\{x^{(\Delta)}(\Delta+\sigma), x^{(\Delta)}\left(\tilde{z}_{1}+\tau\right)\right\} \geq \bar{x}
\end{aligned}
$$

is increasing as a function of $\Delta \in I_{R P R P}$, and $\underline{x}_{\Delta}=\underline{x}$.
Recall $t_{\max }=\tilde{z}_{1}+\tau$. From

$$
0=x^{(\Delta)}\left(z_{\Delta, 2}\right)=-\beta_{U}+\left(x^{(\Delta)}\left(t_{\max }\right)+\beta_{U}\right) e^{-\left(z_{\Delta, 2}-t_{\max }\right)}
$$

we get

$$
\beta_{U} e^{z_{\Delta, 2}}=\left(x^{(\Delta)}\left(t_{\max }\right)+\beta_{U}\right) e^{t_{\max }}=\left(\bar{x}+\beta_{U}\right) e^{t_{\max }}+a\left(e^{\sigma}-1\right) e^{\Delta}
$$

Since $\left(\bar{x}+\beta_{U}\right) e^{t_{\max }}=\beta_{U} e^{\tilde{z}_{2}}$, we conclude that

$$
\beta_{U} e^{z_{\Delta, 2}}=\beta_{U} e^{\tilde{z}_{2}}+a\left(e^{\sigma}-1\right) e^{\Delta}
$$

For the cycle length we obtain $T(\Delta)=z_{\Delta, 2}+\tau>\tilde{z}_{2}+\tau=\tilde{T}$ and the formula for $T(\Delta)$ follows.

Proof of Remark 5.1 1. Let $\delta_{1}>0$. We first show that the expressions defining $T(\Delta)$
in Proposition 5.1 and in Proposition 5.2 yield the same value for $\Delta=\delta_{1}$.
Consider the argument of $\ln$ in (5.8). We have

$$
\begin{aligned}
& \frac{a\left(e^{\sigma}-1\right)}{\beta_{U}} e^{\Delta-\tilde{z}_{2}}+\frac{a\left(\beta_{L}+\beta_{U}\right) e^{\tau+\tilde{z}_{1}-\tilde{z}_{2}}}{\beta_{U}\left(\beta_{L}+a\right)}\left(e^{\Delta-\tilde{z}_{1}}-1\right) \\
& \quad=\frac{a e^{-\tilde{z}_{2}}}{\beta_{U}\left(\beta_{L}+a\right)}\left(\left(\left(\beta_{L}+a\right)\left(e^{\sigma}-1\right)+\left(\beta_{L}+\beta_{U}\right) e^{\tau}\right) e^{\Delta}-\left(\beta_{L}+\beta_{U}\right) e^{\tau+\tilde{z}_{1}}\right)
\end{aligned}
$$

From (5.6) for $\delta_{1}$ it follows that

$$
\beta_{L} e^{\sigma}+a\left(e^{\sigma}-1\right)=\beta_{L} e^{\tilde{z}_{1}-\delta_{1}}
$$

which gives

$$
\left(\beta_{L}+a\right)\left(e^{\sigma}-1\right)=\beta_{L} e^{\tilde{z}_{1}-\delta_{1}}-\beta_{L}
$$

Since

$$
-\beta_{L}+\left(\beta_{L}+\beta_{U}\right) e^{\tau}=e^{\tau}\left(\beta_{U}+\bar{x}\right)=\beta_{U} e^{\tilde{z}_{2}-\tilde{z}_{1}}
$$

we obtain

$$
\begin{aligned}
& \frac{a\left(e^{\sigma}-1\right)}{\beta_{U}} e^{\Delta-\tilde{z}_{2}}+\frac{a\left(\beta_{L}+\beta_{U}\right) e^{\tau+\tilde{z}_{1}-\tilde{z}_{2}}}{\beta_{U}\left(\beta_{L}+a\right)}\left(e^{\Delta-\tilde{z}_{1}}-1\right) \\
& \quad=\frac{a e^{-\tilde{z}_{2}}}{\beta_{U}\left(\beta_{L}+a\right)}\left(\left(\beta_{L} e^{\tilde{z}_{1}-\delta_{1}}+\beta_{U} e^{\tilde{z}_{2}-\tilde{z}_{1}}\right) e^{\Delta}-\left(\beta_{L} e^{\tilde{z}_{1}}+\beta_{U} e^{\tilde{z}_{2}}\right)\right)
\end{aligned}
$$

which for $\Delta=\delta_{1}$ becomes

$$
\frac{a e^{-\tilde{z}_{2}}}{\beta_{U}\left(\beta_{L}+a\right)}\left(\left(\beta_{L} e^{\tilde{z}_{1}-\delta_{1}}+\beta_{U} e^{\tilde{z}_{2}-\tilde{z}_{1}}\right) e^{\Delta}-\left(\beta_{L} e^{\tilde{z}_{1}}+\beta_{U} e^{\tilde{z}_{2}}\right)\right)=\frac{a\left(e^{\delta_{1}-\tilde{z}_{1}}-1\right)}{\beta_{L}+a}
$$

We have

$$
e^{\delta_{1}-\tilde{z}_{1}}-1=\frac{\beta_{L}}{\beta_{L} e^{\sigma}+a\left(e^{\sigma}-1\right)}-1=-\frac{\left(\beta_{L}+a\right)\left(e^{\sigma}-1\right)}{\beta_{L} e^{\sigma}+a\left(e^{\sigma}-1\right)},
$$

which leads to

$$
\frac{a\left(e^{\delta_{1}-\tilde{z}_{1}}-1\right)}{\beta_{L}+a}=-\frac{a\left(e^{\sigma}-1\right)}{\beta_{L} e^{\sigma}+a\left(e^{\sigma}-1\right)}=-\frac{a\left(e^{\sigma}-1\right)}{\beta_{L}} e^{\delta_{1}-\tilde{z}_{1}}
$$

and shows that the formulae for $T(\Delta)$ from Propositions 5.1 and 5.2 yield the same value for $\Delta=\delta_{1}>0$. From (5.7) of Proposition 5.1, this value is strictly less than $\tilde{T}$.
2. By continuity, we infer $T(\Delta)<\tilde{T}$ for $\Delta$ close to $\delta_{1}>0$.

Proof of Proposition 5.4 Since $x^{(\Delta)}(t)>0$ for $\tilde{z}_{1}<t<\Delta+\sigma$ and $\tilde{z}_{1}+\tau<\Delta+\sigma$ we obtain that on $[\Delta+\sigma, \infty)$,

$$
x^{(\Delta)}(t)=-\beta_{U}+\left(x^{(\Delta)}(\Delta+\sigma)+\beta_{U}\right) e^{-(t-(\Delta+\sigma))} \quad \text { as long as } \quad 0<x^{(\Delta)}(t-\tau)
$$

As $-\beta_{U}<0$ there is a smallest zero $z_{\Delta, 2}$ of $x^{(\Delta)}$ in $[\Delta+\sigma, \infty)$, and

$$
x^{(\Delta)}\left(z_{\Delta, 2}+t\right)=\tilde{x}\left(\tilde{z}_{2}+t\right) \text { for all } t \geq 0
$$

It follows that

$$
T(\Delta)=z_{\Delta, 2}+\left(\tilde{z}_{3}-\tilde{z}_{2}\right)-\tilde{z}_{1}=z_{\Delta, 2}+\tilde{T}-\tilde{z}_{2}=z_{\Delta, 2}-\tilde{z}_{0}=z_{\Delta, 2}+\tau
$$

Moreover, $\bar{x}_{\Delta}=x^{(\Delta)}\left(t_{\max }\right)$ if $\Delta \leq \hat{\delta}_{2}$ (in which case $x^{(\Delta)}$ is decreasing on $\left[t_{\max }, \Delta+\right.$ $\sigma]$ ), while for $\Delta>\hat{\delta}_{2}$ the function $x^{(\Delta)}$ is increasing on $\left[t_{\max }, \Delta+\sigma\right]$ and $\bar{x}_{\Delta}=$
$x^{(\Delta)}(\Delta+\sigma)>x^{(\Delta)}\left(t_{\max }\right) \geq \bar{x}$. Hence $\bar{x}_{\Delta} \geq \bar{x}$. Obviously, $\underline{x}_{\Delta}=\underline{x}$. Also, $\bar{x}_{\Delta}$ is strictly decreasing as a function of $\Delta \in I_{R P F P}$, see (5.10).

From

$$
0=x^{(\Delta)}\left(z_{\Delta, 2}\right)=-\beta_{U}+\left(x^{(\Delta)}(\Delta+\sigma)+\beta_{U}\right) e^{-\left(z_{\Delta, 2}-(\Delta+\sigma)\right)}
$$

it follows that

$$
\begin{aligned}
\beta_{U} e^{z_{\Delta, 2}} & =\left(x^{(\Delta)}(\Delta+\sigma)+\beta_{U}\right) e^{\Delta+\sigma} \\
& =\left(\beta_{U} e^{\tilde{z}_{2}-(\Delta+\sigma)}+a\left(1-e^{-\sigma}\right)\right) e^{\Delta+\sigma} \quad(\text { see }(5.10)) \\
& =\beta_{U} e^{\tilde{z}_{2}}+a\left(e^{\sigma}-1\right) e^{\Delta}
\end{aligned}
$$

Hence $z_{\Delta, 2}>\tilde{z}_{2}$, and thereby

$$
T(\Delta)=z_{\Delta, 2}+\tau>\tilde{z}_{2}+\tau=\tilde{T}
$$

Furthermore,

$$
\begin{aligned}
T(\Delta) & =z_{\Delta, 2}+\tau=\tilde{T}+z_{\Delta, 2}-\tilde{z}_{2} \\
& =\tilde{T}+\ln \left(1+\frac{a\left(e^{\sigma}-1\right) e^{\Delta-\tilde{z}_{2}}}{\beta_{U}}\right),
\end{aligned}
$$

and the map

$$
I_{R P F P} \ni \Delta \mapsto T(\Delta) \in \mathbb{R}
$$

is strictly increasing.
Proof of Proposition 5.5 From $x^{(\Delta)}(\Delta+\sigma)<0<x^{(\Delta)}\left(t_{\max }\right)$ we know that $x^{(\Delta)}$ is strictly decreasing on $\left[t_{\max }, \Delta+\sigma\right]$. It follows that there is a single zero in this interval, which is given by

$$
0=x^{(\Delta)}\left(z_{\Delta, 2}\right)=-\beta_{U}+a+\left(x^{(\Delta)}\left(t_{\max }\right)+\beta_{U}-a\right) e^{-\left(z_{\Delta, 2}-t_{\max }\right)}
$$

or equivalently,

$$
\left(\beta_{U}-a\right) e^{z_{\Delta, 2}}=\left(x^{(\Delta)}\left(t_{\max }\right)+\beta_{U}-a\right) e^{t_{\max }}
$$

Also,

$$
\begin{aligned}
\left(\beta_{U}-a\right) e^{z_{\Delta, 2}} & =\left(\bar{x}+a\left(1-e^{\Delta-t_{\max }}\right)+\beta_{U}-a\right) e^{t_{\max }} \\
& =\left(\bar{x}+\beta_{U}\right) e^{t_{\max }}-a e^{\Delta}
\end{aligned}
$$

and we arrive at

$$
\begin{equation*}
\left(\beta_{U}-a\right) e^{z_{\Delta, 2}}=\beta_{U} e^{\tilde{z}_{2}}-a e^{\Delta} \tag{9.5}
\end{equation*}
$$

Since $\beta_{U}>a$ we infer that the map $I_{R P F N} \ni \Delta \mapsto e^{z \Delta, 2-\Delta} \in \mathbb{R}$ is strictly decreasing.
We have $z_{\Delta, 2}<\Delta+\sigma \leq z_{\Delta, 2}+\tau$, and on $\left[\Delta+\sigma, z_{\Delta, 2}+\tau\right]$,

$$
x^{(\Delta)}(t)=-\beta_{U}+\left(x^{(\Delta)}(\Delta+\sigma)+\beta_{U}\right) e^{-(t-(\Delta+\tau))}
$$

is strictly decreasing and negative because we have $x^{(\Delta)}(\Delta+\sigma)+\beta_{U}>0$ from (5.10). It follows that $x^{(\Delta)}$ is strictly increasing on $\left[z_{\Delta, 2}+\tau, \infty\right)$ as long as $x^{(\Delta)}(t-\tau)<0$. There is a smallest zero $z_{\Delta, 3}$ of $x^{(\Delta)}$ in this interval, and

$$
\begin{aligned}
x^{(\Delta)}\left(z_{\Delta, 3}+t\right) & =\tilde{x}\left(\tilde{z}_{1}+t\right) \text { for all } t \geq 0, \\
T(\Delta) & =z_{\Delta, 3}-\tilde{z}_{1}, \\
\bar{x}_{\Delta} & =x^{(\Delta)}\left(t_{\max }\right) \geq \bar{x}, \\
\underline{x}_{\Delta} & =x^{(\Delta)}\left(z_{\Delta, 2}+\tau\right) .
\end{aligned}
$$

We compute

$$
\underline{x}_{\Delta}=x^{(\Delta)}\left(z_{\Delta, 2}+\tau\right)=-\beta_{U}+\left(x^{(\Delta)}(\Delta+\sigma)+\beta_{U}\right) e^{-\left(z_{\Delta, 2}+\tau-(\Delta+\sigma)\right)}
$$

and use

$$
x^{(\Delta)}(\Delta+\sigma)+\beta_{U}=\beta_{U} e^{\tilde{z}_{2}-(\Delta+\sigma)}+a\left(1-e^{-\sigma}\right)>0
$$

from (5.10). This gives

$$
\underline{x}_{\Delta}=-\beta_{U}+\left(\beta_{U} e^{\tilde{z}_{2}}-a e^{\Delta}\right) e^{-\left(z_{\Delta, 2}+\tau\right)}+a e^{-\left(z_{\Delta, 2}+\tau-(\Delta+\sigma)\right)} .
$$

With (9.5) we obtain

$$
\begin{aligned}
\underline{x}_{\Delta} & =-\beta_{U}+\left(\beta_{U}-a\right) e^{-\tau}+a e^{-\left(z_{\Delta, 2}+\tau-(\Delta+\sigma)\right)} \\
& =-\beta_{U}+\beta_{U} e^{-\tau}+a e^{-\tau}\left(e^{-z_{\Delta, 2}+\Delta+\sigma}-1\right) .
\end{aligned}
$$

Since $\Delta+\sigma>z_{\Delta, 2}$ we conclude that

$$
\underline{x}_{\Delta}>-\beta_{U}+\beta_{U} e^{-\tau}=\underline{x} .
$$

We turn to the cycle length $T(\Delta)=z_{\Delta, 3}-\tilde{z}_{1}$. The equation for $z_{\Delta, 3}$, namely,

$$
\begin{aligned}
0 & =x^{(\Delta)}\left(z_{\Delta, 3}\right)=\beta_{L}+\left(x^{(\Delta)}\left(z_{\Delta, 2}+\tau\right)-\beta_{L}\right) e^{-\left(z_{\Delta, 3}-\left(z_{\Delta, 2}+\tau\right)\right)} \\
& =\beta_{L}+\left(\underline{x}_{\Delta}-\beta_{L}\right) e^{-\left(z_{\Delta, 3}-\left(z_{\Delta, 2}+\tau\right)\right)}
\end{aligned}
$$

yields

$$
\begin{aligned}
\beta_{L} e^{z_{\Delta, 3}}= & \left(\beta_{L}-\underline{x}_{\Delta}\right) e^{z_{\Delta, 2}+\tau} \\
= & \left(\beta_{L}+\beta_{U}\right) e^{z_{\Delta, 2}+\tau}-\left(\beta_{U} e^{\tilde{z}_{2}}-a e^{\Delta}\right)-a e^{\Delta+\sigma} \\
& \left(\text { with the formula for } \underline{x}_{\Delta}\right. \text { and (9.5)) } \\
= & \left(\beta_{L}+\beta_{U}\right) e^{z_{\Delta, 2}+\tau}-\beta_{U} e^{\tilde{z}_{2}}-a\left(e^{\sigma}-1\right) e^{\Delta} .
\end{aligned}
$$

Since $\tilde{z}_{3}=\tilde{z}_{2}+\tau+\tilde{z}_{1}$, we have

$$
\begin{aligned}
\beta_{L} e^{\tilde{z}_{3}}+\beta_{U} e^{\tilde{z}_{2}} & =\left(\beta_{L} e^{\tilde{z}_{1}}+\beta_{U} e^{-\tau}\right) e^{\tilde{z}_{2}+\tau} \\
& =\left(\beta_{L}+\beta_{U}\right) e^{\tilde{z}_{2}+\tau}
\end{aligned}
$$

which gives

$$
\beta_{L} e^{z_{\Delta, 3}}=\beta_{L} e^{\tilde{z}_{3}}+\left(\beta_{L}+\beta_{U}\right) e^{\tau}\left(e^{z_{\Delta, 2}}-e^{\tilde{z}_{2}}\right)-a\left(e^{\sigma}-1\right) e^{\Delta} .
$$

Now use (9.5) again to obtain

$$
\beta_{L} e^{z_{\Delta, 3}}=\beta_{L} e^{\tilde{z}_{3}}-\frac{a\left(\beta_{L}+\beta_{U}\right) e^{\tau}}{\beta_{U}-a}\left(e^{\Delta}-e^{\tilde{z}_{2}}\right)-a\left(e^{\sigma}-1\right) e^{\Delta},
$$

which implies the formula for $T(\Delta)$. We have

$$
\frac{\left(\beta_{L}+\beta_{U}\right)}{\beta_{U}-a} e^{\tau}+e^{\sigma}-1=\frac{\beta_{L} e^{\tilde{z}_{1}+\tau}+\left(\beta_{U}-a\right) e^{\sigma}+a}{\beta_{U}-a} .
$$

Thus, the map $I_{R P F N} \ni \Delta \mapsto T(\Delta) \in \mathbb{R}$ is strictly decreasing, since $\beta_{U}>a$.
Proof of Remark 5.2 1. We first show that the expressions defining $T(\Delta)$ in Proposition 5.3 and in Proposition 5.5 (Eqs. 5.9 and 5.13 respectively) yield the same value for $\Delta=\delta_{2}$.
Consider the argument of $\ln$ in (5.13) of Proposition 5.5. We have

$$
\begin{aligned}
& \frac{a\left(e^{\sigma}-1\right)}{\beta_{L}} e^{\Delta-\tilde{z}_{1}-\tilde{T}}+\frac{a\left(\beta_{L}+\beta_{U}\right) e^{-\tilde{z}_{1}}}{\beta_{L}\left(\beta_{U}-a\right)} e^{\Delta-\tilde{z}_{2}} \\
& \quad=\frac{a e^{\Delta-\tilde{z}_{1}-\tilde{T}}}{\beta_{L}\left(\beta_{U}-a\right)}\left(\left(\beta_{U}-a\right)\left(e^{\sigma}-1\right)+\left(\beta_{L}+\beta_{U}\right) e^{\tau}\right)
\end{aligned}
$$

Note that from the definition of $\delta_{2}$ it follows that

$$
\beta_{U} e^{\sigma}-a\left(e^{\sigma}-1\right)=\beta_{U} e^{\tilde{z}_{2}-\delta_{2}}
$$

which gives

$$
\begin{aligned}
\left(\beta_{U}-a\right)\left(e^{\sigma}-1\right)+\left(\beta_{L}+\beta_{U}\right) e^{\tau} & =\beta_{U} e^{\sigma}-a\left(e^{\sigma}-1\right)+\left(\beta_{L}-\underline{x}\right) e^{\tau} \\
& =\beta_{U} e^{\tilde{z}_{2}-\delta_{2}}+\beta_{L} e^{t_{\max }}
\end{aligned}
$$

and leads to

$$
\begin{aligned}
& \frac{a\left(e^{\sigma}-1\right)}{\beta_{L}} e^{\Delta-\tilde{z}_{1}-\tilde{T}}+\frac{a\left(\beta_{L}+\beta_{U}\right) e^{-\tilde{z}_{1}}}{\beta_{L}\left(\beta_{U}-a\right)}\left(e^{\Delta-\tilde{z}_{2}}-1\right) \\
& \quad=\frac{a e^{-\tilde{z}_{1}}}{\beta_{L}\left(\beta_{U}-a\right)}\left(\left(\beta_{U} e^{-\tau-\delta_{2}}+\beta_{L} e^{\tilde{z}_{1}-\tilde{z}_{2}}\right) e^{\Delta}-\left(\beta_{L}+\beta_{U}\right)\right) .
\end{aligned}
$$

Observe that for $\Delta=\delta_{2}$ we have

$$
\left(\beta_{U} e^{-\tau-\delta_{2}}+\beta_{L} e^{\tilde{z}_{1}-\tilde{z}_{2}}\right) e^{\Delta}=\beta_{U} e^{-\tau}+\beta_{L} e^{\tilde{z}_{1}+\delta_{2}-\tilde{z}_{2}}
$$

Using this we obtain that for $\Delta=\delta_{2}$,

$$
\frac{a\left(e^{\sigma}-1\right)}{\beta_{L}} e^{\Delta-\tilde{z}_{1}-\tilde{T}}+\frac{a\left(\beta_{L}+\beta_{U}\right) e^{-\tilde{z}_{1}}}{\beta_{L}\left(\beta_{U}-a\right)}\left(e^{\Delta-\tilde{z}_{2}}-1\right)=\frac{a}{\left(\beta_{U}-a\right)}\left(e^{\delta_{2}-\tilde{z}_{2}}-1\right) .
$$

We have

$$
\begin{aligned}
\frac{a}{\left(\beta_{U}-a\right)}\left(e^{\delta_{2}-\tilde{z}_{2}}-1\right) & =\frac{a}{\left(\beta_{U}-a\right)}\left(\frac{\beta_{U}}{\beta_{U} e^{\sigma}-a\left(e^{\sigma}-1\right)}-1\right) \\
& =-\frac{a\left(e^{\sigma}-1\right)}{\beta_{U} e^{\sigma}-a\left(e^{\sigma}-1\right)} \\
& =-\frac{a\left(e^{\sigma}-1\right)}{\beta_{U}} e^{\delta_{2}-\tilde{z}_{2}} .
\end{aligned}
$$

Thus, the formulae (5.9) and (5.13) from Propositions 5.3 and 5.5 are the same for $\Delta=\delta_{2}$. Since

$$
-\frac{a\left(e^{\sigma}-1\right)}{\beta_{U}} e^{\delta_{2}-\tilde{z}_{2}}<0
$$

we also deduce that the value given by both equations for $\Delta=\delta_{2}$ is strictly larger than $\tilde{T}$.
2. By continuity, we infer $T(\Delta)>\tilde{T}$ for $\Delta$ close to $\delta_{2}$.

Proof of Proposition 5.6 From $0 \leq x^{(\Delta)}(t)$ on $\left[t_{\max }, \Delta+\sigma\right]$ we infer that $x^{(\Delta)}$ decreases on $[\Delta+\sigma, \infty)$ as long as $x^{(\Delta)}(t-\tau)>0$. This yields the existence of a smallest zero $z_{\Delta, 2}$ in $[\Delta+\sigma, \infty)$, and $x^{(\Delta)}\left(t+z_{\Delta, 2}\right)=\tilde{x}\left(t+\tilde{z}_{2}\right)$ for all $t \geq 0$.

Hence

$$
T(\Delta)=z_{\Delta, 2}+\tilde{z}_{3}-\tilde{z}_{2}-\tilde{z}_{1}=z_{\Delta, 2}+\tilde{T}-\tilde{z}_{2}=z_{\Delta, 2}+\tau
$$

We have

$$
-\beta_{U}+\left(x^{(\Delta)}(\Delta+\sigma)+\beta_{U}\right) e^{-\left(z_{\Delta, 2}-(\Delta+\sigma)\right)}=0
$$

and conclude that

$$
\begin{aligned}
\beta_{U} e^{z_{\Delta, 2}} & =\left(x^{(\Delta)}(\Delta+\sigma)+\beta_{U}\right) e^{\Delta+\sigma} \\
& =\left(\beta_{U} e^{\tilde{z}_{2}-\sigma-\Delta}+a\left(1-e^{-\sigma}\right)\right) e^{\Delta+\sigma} \\
& =\beta_{U} e^{\tilde{z}_{2}}+a\left(e^{\sigma}-1\right) e^{\Delta},
\end{aligned}
$$

which implies (5.9), the desired formula for $T(\Delta)$.
Proof of Proposition 5.7 Since $[\Delta, \Delta+\sigma] \subset\left(t_{\max }, \tilde{T}\right)$ we have

$$
x^{(\Delta)}(t)=-\beta_{U}+a+\left(x^{(\Delta)}(\Delta)+\beta_{U}-a\right) e^{-(t-\Delta)} \quad \text { on } \quad[\Delta, \Delta+\sigma] .
$$

Using $x^{(\Delta)}(\Delta)+\beta_{U}-a>0$ we see that $x^{(\Delta)}$ is strictly decreasing on $[\Delta, \Delta+\sigma]$, and has a unique zero $z_{\Delta, 2}$ in $[\Delta, \Delta+\sigma)$, which is given implicitly by

$$
0=x^{(\Delta)}\left(z_{\Delta, 2}\right)=-\beta_{U}+a+\left(x^{(\Delta)}(\Delta)+\beta_{U}-a\right) e^{-\left(z_{\Delta, 2}-\Delta\right)}
$$

Combining this with (5.2) gives

$$
\left(\beta_{U}-a\right) e^{z_{\Delta, 2}}=\beta_{U} e^{\tilde{z}_{2}}-a e^{\Delta}
$$

which is the same as (9.5). As $z_{\Delta, 2} \geq \Delta$ we infer

$$
\beta_{U}\left(e^{z_{\Delta, 2}}-e^{\tilde{z}_{2}}\right)=a\left(e^{z_{\Delta, 2}}-e^{\Delta}\right) \geq 0
$$

Hence $z_{\Delta, 2} \geq \tilde{z}_{2}$. Since $x^{(\Delta)}(\Delta) \geq 0$ and $\beta_{U}>a$, the map $[\Delta, \Delta+\sigma] \ni t \mapsto x^{(\Delta)}(t)$ is strictly decreasing. The rest of the proof is the same as the proof of Proposition 5.5 starting after (9.5).

Proof of Proposition 5.8 1. We have $x^{(\Delta)}\left(\tilde{z}_{2}\right)=0$, and the function $x^{(\Delta)}$ is strictly decreasing on $\left[\tilde{z}_{2}, \Delta\right]$, monotone on $[\Delta, \Delta+\sigma]$ with $x^{(\Delta)}(t)<0$ for $\Delta \leq t<$ $\Delta+\sigma$, and strictly decreasing on $[\Delta+\sigma, \tilde{T}]$. It follows that $x^{(\Delta)}$ is strictly increasing on $[\tilde{T}, \infty)$ as long as $x^{(\Delta)}(t-\tau)<0$. This yields a first zero $z_{\Delta, 3}$ of $x^{(\Delta)}$ in $[\tilde{T}, \infty)$ and $x^{(\Delta)}\left(z_{\Delta, 3}+t\right)=\tilde{x}\left(\tilde{z}_{1}+t\right)$ for all $t \geq 0$, and thus

$$
\begin{aligned}
T(\Delta) & =z_{\Delta, 3}+\left(\tilde{z}_{2}-\tilde{z}_{1}\right)-\tilde{z}_{2}=z_{\Delta, 3}-\tilde{z}_{1}, \\
\underline{x}_{\Delta} & =\min \left\{x^{(\Delta)}(\Delta), x^{(\Delta)}(\tilde{T})\right\},
\end{aligned}
$$

and

$$
\bar{x}_{\Delta}=\bar{x} .
$$

Observe that

$$
\begin{aligned}
x^{(\Delta)}(\tilde{T}) & =-\beta_{U}+\beta_{U} e^{-\tau}+a\left(e^{\sigma}-1\right) e^{\Delta-\tilde{T}} \\
& =\underline{x}+a\left(e^{\sigma}-1\right) e^{\Delta-\tilde{T}}>\underline{x} .
\end{aligned}
$$

Using this and $x^{(\Delta)}(\Delta)=\tilde{x}(\Delta)>\underline{x}$ we have $\underline{x}_{\Delta}>\underline{x}$. The equation

$$
0=x^{(\Delta)}\left(z_{\Delta, 3}\right)=\beta_{L}+\left(x^{(\Delta)}(\tilde{T})-\beta_{L}\right) e^{-\left(z_{\Delta, 3}-\tilde{T}\right)}
$$

is equivalent to

$$
\beta_{L} e^{z_{\Delta, 3}}=\left(\beta_{L}-x^{(\Delta)}(\tilde{T})\right) e^{\tilde{T}}
$$

which gives

$$
\begin{aligned}
\beta_{L} e^{z_{\Delta, 3}} & =\left(\beta_{L}+\beta_{U}\right) e^{\tilde{T}}-\left(\beta_{U} e^{-\tau}+a\left(e^{\sigma}-1\right) e^{\Delta-\tilde{T}}\right) e^{\tilde{T}} \\
& =\left(\beta_{L}+\beta_{U}-\beta_{U} e^{-\tau}\right) e^{\tilde{T}}-a\left(e^{\sigma}-1\right) e^{\Delta} .
\end{aligned}
$$

We have $\beta_{L}+\beta_{U}-\beta_{U} e^{-\tau}=\beta_{L} e^{\tilde{z}_{1}}$ and we conclude that

$$
\beta_{L} e^{z_{\Delta, 3}}=\beta_{L} e^{\tilde{z}_{1}+\tau+\tilde{z}_{2}}-a\left(e^{\sigma}-1\right) e^{\Delta},
$$

from which we obtain (5.14) for $T(\Delta)=z_{\Delta, 3}-\tilde{z}_{1}$.
2. Observe that for every $\Delta \in I_{F N F N}$ we have

$$
x^{(\Delta)}(\tilde{T})-\tilde{x}(\Delta)=g(\Delta)
$$

with the strictly increasing function

$$
g: \mathbb{R} \ni \Delta \mapsto a\left(e^{\sigma}-1\right) e^{\Delta-\tilde{T}}+\beta_{U} e^{-\tau}-\beta_{U} e^{\tilde{z}_{2}-\Delta} \in \mathbb{R}
$$

which has a single zero at $\Delta=\bar{\delta}$ since $g(\Delta) \rightarrow-\infty$ as $\Delta \rightarrow-\infty, g(\Delta) \rightarrow \infty$ as $\Delta \rightarrow \infty$, and $g(\Delta)=0$ if and only if

$$
\beta_{U}\left(e^{\tilde{T}-\Delta}\right)^{2}-\beta_{U} e^{\tilde{T}-\Delta}-a\left(e^{\sigma}-1\right) e^{\tau}=0
$$

or equivalently,

$$
\begin{aligned}
e^{\tilde{T}-\Delta} & =\frac{1}{2}+\sqrt{\frac{1}{4}+\frac{a\left(e^{\sigma}-1\right) e^{\tau}}{\beta_{U}}} \\
& =\frac{\beta_{U}+\sqrt{\beta_{U}^{2}+4 a \beta_{U}\left(e^{\sigma}-1\right) e^{\tau}}}{2 \beta_{U}}
\end{aligned}
$$

For $\Delta<\bar{\delta}$ in $I_{F N F N}$ we have $g(\Delta)<0$, hence

$$
\underline{x}_{\Delta}=\min \left\{\tilde{x}(\Delta), x^{(\Delta)}(\tilde{T})\right\}=x^{(\Delta)}(\tilde{T}),
$$

and the formula (5.10) for $x^{(\Delta)}(\tilde{T})$ in Part 1 above shows that in case $I_{F N F N} \cap$ $(-\infty, \bar{\delta}) \neq \emptyset$ the map

$$
I_{F N F N} \cap(-\infty, \bar{\delta}) \ni \Delta \mapsto \underline{x}_{\Delta} \in \mathbb{R}
$$

is strictly increasing. For $\Delta>\bar{\delta}$ in $I_{F N F N}$ we get $0<g(\Delta)$, hence

$$
\underline{x}_{\Delta}=\min \left\{\tilde{x}(\Delta), x^{(\Delta)}(\tilde{T})\right\}=\tilde{x}(\Delta),
$$

and we see that in case $I_{F N F N} \cap(\bar{\delta}, \infty) \neq \emptyset$ the map

$$
I_{F N F N} \cap(\bar{\delta}, \infty) \ni \Delta \mapsto \underline{x}_{\Delta} \in \mathbb{R}
$$

is strictly decreasing.
Proof of Remark 5.3 The strictly increasing function $g$ from the previous proof satisfies

$$
\begin{aligned}
g\left(\tilde{z}_{2}\right) & =a\left(e^{\sigma}-1\right) e^{-\tau}+\beta_{U}\left(e^{-\tau}-1\right) \\
& <\beta_{U}\left(\left(e^{\sigma}-1\right) e^{-\tau}+e^{-\tau}-1\right) \quad\left(\text { with } a<\beta_{U}\right) \\
& =\beta_{U}\left(e^{\sigma-\tau}-1\right) \\
& \leq 0 \quad(\text { since } \sigma \leq \tau)
\end{aligned}
$$

and

$$
g(\tilde{T}-\sigma)=a\left(e^{\sigma}-1\right) e^{-\sigma}+\beta_{U} e^{-\tau}-\beta_{U} e^{\tilde{z}_{2}-\tilde{T}+\sigma}>0
$$

if and only if $\beta_{U} e^{\sigma-\tau}<a$. By the intermediate value theorem the only zero $\bar{\delta}$ of $g$ belongs to the interval $\left(\tilde{z}_{2}, \tilde{T}-\sigma\right)$ in case $\beta_{U} e^{\sigma-\tau}<a$. Otherwise $\bar{\delta} \geq \tilde{T}-\sigma$.

Proof of Proposition 5.9 From $0>x^{(\Delta)}(\Delta+\sigma)>x^{(\Delta)}(\tilde{T})$ we obtain $x^{(\Delta)}(t)<0$ on ( $\left.\tilde{z}_{2}, \Delta+\sigma\right]$. It follows that on [ $\Delta+\sigma, \infty$ ) the function $x^{(\Delta)}$ is strictly increasing as long as $x^{(\Delta)}(t-\tau)<0$, and there is a smallest zero $z_{\Delta, 3}$ of $x^{(\Delta)}$ in $[\Delta+\sigma, \infty)$. Moreover, $x^{(\Delta)}\left(z_{\Delta, 3}+t\right)=\tilde{x}\left(\tilde{z}_{1}+t\right)$ for all $t \geq 0$, hence

$$
T(\Delta)=z_{\Delta, 3}+\left(\tilde{z}_{2}-\tilde{z}_{1}\right)-\tilde{z}_{2}=z_{\Delta, 3}-\tilde{z}_{1}
$$

and

$$
\underline{x}_{\Delta}=x^{(\Delta)}(\tilde{T})>\underline{x},
$$

and $\bar{x}_{\Delta}=\bar{x}$. As the map $[\tilde{T}-\sigma, \tilde{T}) \ni \Delta \mapsto x^{(\Delta)}(\tilde{T}) \in \mathbb{R}$ is strictly decreasing we infer that also the map $I_{F N R N} \ni \Delta \mapsto \underline{x}_{\Delta} \in \mathbb{R}$ is strictly decreasing.

From

$$
0=x^{(\Delta)}\left(z_{\Delta, 3}\right)=\beta_{L}+\left(x^{(\Delta)}(\Delta+\sigma)-\beta_{L}\right) e^{-\left(z_{\Delta, 3}-(\Delta+\sigma)\right)}
$$

we obtain

$$
\begin{aligned}
\beta_{L} e^{z_{\Delta, 3}} & =\left(\beta_{L}-x^{(\Delta)}(\Delta+\sigma)\right) e^{\Delta+\sigma} \\
& =\left(\beta_{L}-\beta_{L}+\beta_{L} e^{\tilde{z}_{1}-(\Delta+\sigma-\tilde{T})}-a\left(1-e^{-\sigma}\right)\right) e^{\Delta+\sigma} \\
& =\beta_{L} e^{\tilde{z}_{1}+\tilde{T}}-a\left(e^{\sigma}-1\right) e^{\Delta}
\end{aligned}
$$

and the formula for $T(\Delta)=z_{\Delta, 3}-\tilde{z}_{1}$ follows.
Proof of Proposition 5.10 We have $x^{(\Delta)}(t)<0$ on $\left(\tilde{z}_{2}, \tilde{T}\right]$, and there is a first zero of $x^{(\Delta)}$ in $(\tilde{T}, \Delta+\sigma]$, which is given by

$$
0=x^{(\Delta)}\left(z_{\Delta, 3}\right)=\beta_{L}+a+\left(x^{(\Delta)}(\tilde{T})-\left(\beta_{L}+a\right)\right) e^{-\left(z_{\Delta, 3}-\tilde{T}\right)}
$$

or equivalently,

$$
\begin{aligned}
\left(\beta_{L}+a\right) e^{z_{\Delta, 3}} & =\left(\beta_{L}+a-x^{(\Delta)}(\tilde{T})\right) e^{\tilde{T}} \\
& =\left(\beta_{L}+a-\left(\underline{x}+a-a e^{\Delta-\tilde{T}}\right)\right) e^{\tilde{T}} \\
& =\left(\beta_{L}-\underline{x}\right) e^{\tilde{T}}+a e^{\Delta} \\
& =\beta_{L} e^{\tilde{z}_{1}+\tilde{T}}+a e^{\Delta} .
\end{aligned}
$$

Incidentally, this shows that the map $I_{F N R P} \ni \Delta \mapsto z_{\Delta, 3}-\Delta \in \mathbb{R}$ is strictly decreasing.

On $\left[z_{\Delta, 3}, \Delta+\sigma\right]$ the function $x^{(\Delta)}$ is strictly increasing. Since $x^{(\Delta)}(t-\tau)<0$ on ( $\tilde{z}_{2}, z_{\Delta, 3}$ ), we infer that $x^{(\Delta)}$ is strictly increasing on $\left[z_{\Delta, 3}, z_{\Delta, 3}+\tau\right]$. Hence $0<x^{(\Delta)}(t)$ on $\left(z_{\Delta, 3}, z_{\Delta, 3}+\tau\right]$. This implies that $x^{(\Delta)}$ is strictly decreasing on $\left[z_{\Delta, 3}+\tau, \infty\right)$ as long as $x^{(\Delta)}(t-\tau)>0$, and that there is a first zero $z_{\Delta, 4}$ in $\left(z_{\Delta, 3}+\tau, \infty\right)$. We obtain $x^{(\Delta)}\left(z_{\Delta, 4}+t\right)=\tilde{x}\left(\tilde{z}_{2}+t\right)$ for all $t \geq 0$, which yields

$$
T(\Delta)=z_{\Delta, 4}-\tilde{z}_{2}
$$

$\underline{x}_{\Delta}=\min \left\{\tilde{x}(\Delta), x^{(\Delta)}(\tilde{T})\right\}>\underline{x}$, and

$$
\bar{x}_{\Delta}=x^{(\Delta)}\left(z_{\Delta, 3}+\tau\right)
$$

As in the proof of Proposition 5.9 we infer that the map $I_{F N R P} \ni \Delta \mapsto \underline{x}_{\Delta} \in \mathbb{R}$ is strictly decreasing.

From the formula for $x^{(\Delta)}(\Delta+\sigma)$ we obtain

$$
\begin{aligned}
x^{(\Delta)}\left(z_{\Delta, 3}+\tau\right) & =\beta_{L}+\left(x^{(\Delta)}(\Delta+\sigma)-\beta_{L}\right) e^{-\left(z_{\Delta, 3}+\tau-(\Delta+\sigma)\right)} \\
& =\beta_{L}-\beta_{L} e^{\tilde{z}_{1}+\tilde{T}-\left(z_{\Delta, 3}+\tau\right)}+a\left(e^{\sigma}-1\right) e^{\Delta} e^{-\left(z_{\Delta, 3}+\tau\right)}
\end{aligned}
$$

which can be rewritten as

$$
\bar{x}_{\Delta}=\beta_{L}-\left(\beta_{L} e^{\tilde{z}_{1}+\tilde{T}}+a e^{\Delta}\right) e^{-\left(z_{\Delta, 3}+\tau\right)}+a e^{\sigma+\Delta} e^{-\left(z_{\Delta, 3}+\tau\right)} .
$$

Hence,

$$
\begin{aligned}
\bar{x}_{\Delta} & =\beta_{L}-\left(\beta_{L}+a\right) e^{-\tau}+a e^{\sigma+\Delta} e^{-\left(z_{\Delta, 3}+\tau\right)} \\
& =\bar{x}+a e^{-\tau}\left(e^{\sigma+\Delta-z_{\Delta, 3}}-1\right) \\
& >\bar{x} .
\end{aligned}
$$

We also know that the map $I_{F N R P} \ni \Delta \mapsto \bar{x}_{\Delta} \in \mathbb{R}$ is strictly increasing.
From

$$
0=x^{(\Delta)}\left(z_{\Delta, 4}\right)=-\beta_{U}+\left(x^{(\Delta)}\left(z_{\Delta, 3}+\tau\right)+\beta_{U}\right) e^{-\left(z_{\Delta, 4}-\left(z_{\Delta, 3}+\tau\right)\right)}
$$

we have

$$
\beta_{U} e^{z_{\Delta, 4}}=\left(\beta_{L}+\beta_{U}\right) e^{z_{\Delta, 3}+\tau}-\beta_{L} e^{\tilde{z}_{1}+\tilde{T}}+a\left(e^{\sigma}-1\right) e^{\Delta}
$$

Use (9.4) to obtain

$$
\beta_{L} e^{\tilde{z}_{1}+\tilde{T}}+\beta_{U} e^{\tilde{z}_{2}+\tilde{T}}=\left(\beta_{L}+\beta_{U}\right) e^{\tau+\tilde{z}_{1}+\tilde{T}}
$$

which, combined with the previous equation, gives

$$
\beta_{U} e^{z_{\Delta, 4}}=\beta_{U} e^{\tilde{z}_{2}+\tilde{T}}+\left(\beta_{L}+\beta_{U}\right) e^{\tau}\left(e^{z_{\Delta, 3}}-e^{\tilde{z}_{1}+\tilde{T}}\right)+a\left(e^{\sigma}-1\right) e^{\Delta}
$$

Now the formula for $z_{\Delta, 3}$ leads to

$$
\beta_{U} e^{z_{\Delta, 4}}=\beta_{U} e^{\tilde{z}_{2}+\tilde{T}}+\frac{a\left(\beta_{L}+\beta_{U}\right) e^{\tau}}{\beta_{L}+a}\left(e^{\Delta}-e^{\tilde{z}_{1}+\tilde{T}}\right)+a\left(e^{\sigma}-1\right) e^{\Delta}
$$

which yields the formula for $T(\Delta)=z_{\Delta, 4}-\tilde{z}_{2}$.

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