

CORRELATION FUNCTIONS AND DENSITIES

13 NOVEMBER, 1992

FILE: CORREL.TEX

MICHAEL C. MACKEY

Centre for Nonlinear Dynamics in Physiology and Medicine
and
Departments of Physiology, Physics and Mathematics
McGill University
Montreal, Quebec, Canada

We first express correlations for ergodic transformations in terms of phase space averages. Let $S : [0, 1] \rightarrow [0, 1]$ be an ergodic transformation with unique stationary density f_* , and $F, G : [0, 1] \rightarrow R$ be integrable functions. Then the correlation of F with G is given by

$$\begin{aligned} C_{F,G}(\tau) &\equiv \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{t=1}^N F(x_{t+\tau})G(x_t) \\ &= \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{t=1}^N F(S^{t+\tau}(x))G(S^t(x)) \\ &= \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{t=1}^N F(S^\tau \circ S^t(x))G(S^t(x)) \\ &= \int_0^1 F(S^\tau(x))G(x)f_*(x)dx, \end{aligned}$$

where the last line follows by the assumed ergodicity and the Birkhoff individual ergodic theorem.

If we identify $F = G = x$ so $C_{xx}(\tau)$ is the autocorrelation of x , and assume $S^\tau(x)$ and $f_*(x)$ are even functions about the point $x = \frac{1}{2}$ for $\tau \geq 1$, then

$$\begin{aligned} C_{xx}(\tau) &= \int_0^1 xS^\tau(x)f_*(x)dx \\ &= \int_0^1 \left(x - \frac{1}{2}\right) S^\tau(x)f_*(x)dx + \frac{1}{2} \int_0^1 S^\tau(x)f_*(x)dx \\ &= \begin{cases} \langle x^2 \rangle & \tau = 0 \\ \frac{1}{2} \langle x \rangle & \tau \geq 1. \end{cases} \end{aligned}$$

Since f_* was assumed to be even about $x = \frac{1}{2}$, we have $\langle x \rangle = \frac{1}{2}$ so $\frac{1}{2} \langle x \rangle \equiv \langle x \rangle^2$. Now let $\rho_{xx}(\tau)$ be the covariance of x : $\rho_{xx}(\tau) \equiv C_{xx}(\tau) - \langle x \rangle^2$. Then

$$\rho_{xx}(\tau) = \begin{cases} \langle x^2 \rangle - \langle x \rangle^2 & \tau = 0 \\ 0 & \tau \geq 1. \end{cases}$$

Independence of η . The auto-correlation of η_t^i defined by (19) is

$$C_{\eta^2}^i(\tau) = \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{t=1}^N \eta_t^i \eta_{t+\tau}^i \quad \tau = 0, 1, \dots \quad (20)$$

If it is the case that the covariance $\rho_{\eta^2}^i \equiv 0$ for all $\tau \geq 1$, then the η_t^i can be treated as independent variables. If further for t sufficiently large the η_t^i have achieved a stationary density g , then we can consider the η_t^i to be not only independent but also identically distributed (i.i.d.) random variables.

Inserting η from (19) we may write (20) explicitly as

$$C_{\eta^2}^i(\tau) = \frac{\epsilon^2}{4} \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{t=1}^N \{S(x_t^{i-1})S(x_{t+\tau}^{i-1}) + S(x_t^{i+1})S(x_{t+\tau}^{i+1}) \\ + S(x_t^{i-1})S(x_{t+\tau}^{i+1}) + S(x_t^{i+1})S(x_{t+\tau}^{i-1})\} \quad (21)$$

Note that the first two terms of $C_{\eta^2}^i$ are the auto-correlation functions of the map S at lattice sites $(i-1)$ and $(i+1)$ respectively, $C_{S^2}^{i-1}$ and $C_{S^2}^{i+1}$. This and the symmetry in the lattice allows us to write (21) in the form

$$C_{\eta^2}^i(\tau) = \frac{\epsilon^2}{2} \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{t=1}^N S(x_t^{i-1})S(x_{t+\tau}^{i+1}) + \frac{\epsilon^2}{2} \begin{cases} \langle x^2 \rangle & \tau = 0 \\ \frac{1}{2} \langle x \rangle & \tau \geq 1. \end{cases} \\ = \frac{\epsilon^2}{2} \int_0^1 \int_0^1 x^{i-1} S^\tau(x^{i+1}) f_*(x^{i-1}, x^{i+1}) dx^{i-1} dx^{i+1} + \frac{\epsilon^2}{2} \begin{cases} \langle x^2 \rangle & \tau = 0 \\ \frac{1}{2} \langle x \rangle & \tau \geq 1. \end{cases}$$

For ϵ small we expect that x^{i-1} and x^{i+1} should be approximately independent so the integral may be evaluated to give

$$C_{\eta^2}^i = \frac{\epsilon^2}{2} \begin{cases} \langle x^2 \rangle + \langle x \rangle^2 & \tau = 0 \\ \frac{1}{2} \langle x \rangle + \langle x \rangle^2 & \tau \geq 1, \end{cases}$$

and thus $\rho_{\eta^2}(\tau)$ takes the form

$$\rho_{\eta^2}(\tau) \simeq \frac{\epsilon^2}{2} \begin{cases} \langle x^2 \rangle - \langle x \rangle^2 & \tau = 0 \\ 0 & \tau \geq 1. \end{cases}$$

Thus for small ϵ the covariance of η is of $\mathcal{O}(\epsilon^2)$ for $\tau = 0$ and approximately 0 for $\tau \geq 1$, and we can take the η_t to be independent of one another. Numerical calculations indicate that for $\epsilon \leq 10^{-1}$ our estimate of $\rho_{\eta^2}(\tau)$ is accurate.

Density of the Distribution of η . For small ϵ we expect that g should be given by the convolution of the stationary density of each of the terms making up η :

$$g(z) = \int_0^\epsilon f_*(z-y) f_*(y) dy.$$

If $\text{supp } f_* = [0, 1]$ then $\text{supp } g = [0, \epsilon]$.

Cross Correlation of T with η . The nearest neighbor spatial-averaged cross-correlation of T with η is defined by

$$C_{T\eta}(t) = \frac{1}{L} \sum_{i=1}^L T(x_t^i) \eta_t^i \quad (22)$$

If it can be shown that $C_{T\eta}(t)$ is negligible for large t , then T and η can be taken as approximately independent as was assumed in the derivation of (18).

To estimate the cross-correlation $C_{T\eta}(t)$ between T and η , we insert the expressions for T and η , so equation (22) can be written in the explicit form

$$C_{T\eta}(t) = \epsilon(1-\epsilon) \left\{ \frac{1}{L} \sum_{i=1}^L S(x_t^i) S(x_t^{i+1}) \right\} \quad (23)$$

because of the cyclicity of the lattice. If the lattice is large so $1 \ll L$, then the bracketed term in (23) should approximate $\langle x^2 \rangle$. Therefore

$$C_{T\eta}(t) \simeq \langle x^2 \rangle \epsilon(1-\epsilon) \quad (24)$$

Numerical computations of $C_{T\eta}(t)$ indicate that (24) accurately predicts its value.