The Role of Surface Infiltration in Hydromechanical Coupling Effects in an Unsaturated Porous Medium of Semi-Infinite Extent

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Rainfall infiltration into an unsaturated region of the earth’s surface is a pervasive natural phenomenon. During the rainfall-induced seepage process, the soil skeleton can deform and the permeability can change with the water content in the unsaturated porous medium. A coupled water infiltration and deformation formulation is used to examine a problem related to the mechanics of a two-dimensional region of semi-infinite extent. The van Genuchten model is used to represent the soil-water characteristic curve. The model, incorporating coupled infiltration and deformation, was developed to resolve the coupled problem in a semi-infinite domain based on numerical methods. The numerical solution is verified by the analytical solution when the coupled effects in an unsaturated medium of semi-infinite extent are considered. The computational results show that a numerical procedure can be employed to examine the semi-infinite unsaturated seepage incorporating coupled water infiltration and deformation. The analysis indicates that the coupling effect is significantly influenced by the boundary conditions of the problem and varies with the duration of water infiltration.

1. Introduction

Coupling between water infiltration and mechanical deformation in partially saturated porous media is central to many natural and man-made systems in civil and environmental engineering. During water infiltration the pore-water pressure is redistributed, on one hand by the hydraulic properties of the partially saturated soil including retention characteristics and permeability and on the other hand by the external loading due to climate conditions (rainfall intensity, duration, and evapotranspiration rate). Changes in the pore-water pressure are generated by infiltration, which in turn modifies the hydraulic domain and induces deformations of the partially saturated porous medium. Alternatively, any variation in the mechanical loading can exert an effect on the infiltration process. It is indeed the hydromechanical coupled response of a partially saturated porous medium that is responsible for the most common instabilities associated with water infiltration: landslides and excessive settlements, due to collapse or shear strength reduction [1, 2].

The prediction of the hydromechanical behavior of geomaterials due to rainfall infiltration plays an important role in addressing geotechnical and geoenvironmental issues related to ground movement or landslides and other geologic disasters. Examples of the effects of saturation variation induced by wetting on the soil deformation have been reported at both laboratory and in situ scales [1, 3, 4].

Analytical solutions play an important role in the examination of water infiltration in saturated and unsaturated heterogeneous porous media [5–13]. Srivastava and Yeh [5] derived an analytical approach to describe one-dimensional rainfall infiltration towards the water table. In recent years, analytical solutions to the one-dimensional problem incorporating rainfall-induced hydromechanical coupling in unsaturated porous media have been developed [11, 14, 15].
The interactions between infiltration and deformation in porous media have been modelled computationally using various multiphysics schemes. Numerical techniques for analyzing coupled water infiltration and deformation in partially saturated porous media are being increasingly utilized because rainfall infiltration involves complicated geometries, soil heterogeneity, and complex spatial and temporal boundary conditions [16]. Analytical approaches to such practical situations are difficult [17]. Many numerical approaches that incorporate the coupled seepage and deformation behaviors involved in water infiltration into a partially saturated porous medium have been presented in the literature in both the soil mechanics and geotechnical disciplines [18]. Numerical models have also been developed to investigate the coupled hydromechanical behavior of partially saturated soils where ground deformation occurs as a result of the water table fluctuation and surface loading [19–21] or the development of soil slope instability [22, 23]. Multiphase, coupled elastoviscoplastic finite element analysis has been used to examine the evolution of pore-water pressure and deformations due to rainwater infiltration into a one-dimensional soil column [17].

A semi-infinite unsaturated extension to the 2D coupled problem is seldom considered when the effects of both seepage and deformation in a partially saturated porous media are being investigated. The numerical method presented in this paper examines rainfall infiltration into a two-dimensional semi-infinite region that also incorporates the coupled seepage and deformation in the unsaturated porous medium. The two-dimensional computational model was developed to resolve the semi-infinite infiltration coupled with deformation in partially saturated porous media. The numerical solution is verified by an analytical solution for the 1D coupled infiltration, and the effect of semi-infinite boundaries on the coupling is analyzed. We then use these numerical approaches to examine the hydromechanical coupling in a 2D semi-infinite unsaturated porous medium.

2. Governing Equations for Coupled Seepage and Deformation in Unsaturated Soils

In the development of a coupled unsaturated seepage and deformation model we use the following assumptions: (i) the soil fabric is deformable, the grains are nondeformable, and the water in the pore space is incompressible; (ii) hysteresis of the soil-water characteristic curve is neglected; (iii) the pore-air pressure within the porous medium is assumed to be constant.

2.1. Static Equilibrium Equations. The equation of equilibrium for a soil mass can be written as

\[ \sigma_{ij,j} + f_i = 0, \]  

where \( \sigma_{ij} \) is the total stress tensor and \( f_i \) is the body force vector.

A thermodynamically consistent effective stress equation has been developed in terms of the effective stress tensor \( \sigma'_{ij} \) in the unsaturated porous medium [24, 25]. On the basis of Bishop's effective stress, the effective stress for a partially saturated porous medium, related to the skeletal deformation, can be given as [26]

\[ \sigma'_{ij} = \sigma_{ij} - [u_w S_r + u_a (1 - S_r)] \delta_{ij}, \]  

where \( u_a \) is the pore-air pressure; \( u_w \) is the pore-water pressure; \( S_r \) is the degree of saturation; \( \delta_{ij} \) is Kronecker's delta function.

The soil skeleton is assumed to exhibit linear isotropic elastic behavior; a constitutive equation can be written as [27]

\[ \varepsilon_{ij} = \frac{\sigma'_{ij}}{2G} - \frac{\lambda}{2G(3\lambda + 2G)} \delta_{ij}, \]  

where \( \varepsilon_{ij} \) is the strain tensor; \( G = \mu = E/(2(1+\nu)) \); \( \lambda = E\nu/(1+\nu)(1-2\nu) \) where \( E \) is the elastic modulus; \( \nu \) is Poisson's ratio of the soil, and \( E \) is the elastic modulus; \( \mu \) and \( \lambda \) are Lamé's constants.

The linearized strain-displacement relations are

\[ \varepsilon_{ij} = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \]  

in which \( u_i, u_j \) are displacement components.

We restrict our attention to the class of problems where the deformation corresponds to a state of plane strain. From assumption (iii), the pore-air pressure is a constant and therefore the partial derivative of the pore-air pressure with respect to time is zero. The external stress applied to the soil mass is assumed to remain constant with time (no loading or the applied loading is unchanged). Making use of (2) to (3), the partial derivative of the volumetric strain can be obtained as

\[ \frac{\partial \varepsilon_v}{\partial t} = -2 \frac{\partial (S_r u_w)}{\partial t} \]  

in which \( \varepsilon_v \) is the volumetric strain and \( C = E/(1+\nu)(1-2\nu) \).

2.2. Soil-Water Characteristic Curve and Hydraulic Conductivity. If the van Genuchten model [28] is used to describe the soil-water characteristic curve (SWCC), we can use the following expression for the volumetric moisture content:

\[ \theta = \theta_r + \frac{\theta_s - \theta_r}{[1 + (\alpha h)^n]^{1/m}}, \]  

where \( h \) is the pressure head, the pore-water pressure \( u_w = \gamma_w h, \gamma_w \) is the unit weight of water, \( \alpha, n, \) and \( m \) are fitting parameters; \( m = 1 - 1/n, \theta_r \) is the residual volumetric moisture content, and \( \theta_s \) is the saturated moisture content.

The unsaturated hydraulic conductivity can be expressed as

\[ k = k_s S_e^{1/2} \left[ 1 - \left( 1 - S_e^{1/m} \right)^m \right]^2, \]  

in which \( k_s \) is the coefficient of infiltration at saturation and \( S_e = 1/[1 + (\alpha h)^n] \) is the effective saturation.
2.3. Seeage Equations. The two-dimensional Richards’ equation is employed, which, when combined with Darcy’s Law for an unsaturated medium, gives

$$\frac{\partial}{\partial x} \left( k \frac{\partial h}{\partial x} \right) + \frac{\partial}{\partial z} \left( k \frac{\partial h}{\partial z} - 1 \right) = \frac{\partial \theta}{\partial t},$$  \hspace{1cm} (8)

where $\theta$ is the volumetric water content; $k$ is the coefficients of hydraulic conductivity in the principal directions $x$ and $z$; $t$ is the time variable. (Implicit in (8) is the assumption of isotropy of the porous medium.) The hydraulic properties of the porous medium are assumed to be isotropic, and the elastic deformability characteristics of the porous medium are considered as isotropic.

The change in the volumetric water content in an unsaturated porous medium is related to the normal stress or strain and the suction of the partially saturated porous medium. Dakshanamurthy et al. [29] proposed that the constitutive equation for an unsaturated porous medium be expressed as

$$\theta = \beta e + \omega (u_a - u_w)$$ \hspace{1cm} (9)

in which $(u_a - u_w)$ is the matric suction and the coefficients $\beta$ and $\omega$ are used to express the volume change modulus $\beta = E/H(1 - 2\nu)$ and $\omega = 1/R - 3\beta/H$, where $H$ is the elastic modulus for the water phase with respect to the effective stress [24]; $R$ is the elastic modulus for the water phase with respect to matric suction $(u_a - u_w)$ [29]; $1/R$ is the envelope of the soil-water characteristic curve.

Assuming that $H$ is a constant, we obtain from the van Genuchten model [28]

$$\frac{1}{R} = \frac{\partial \theta}{\partial \psi} \frac{\partial h}{\partial \psi} = \left( -\frac{1}{\rho_w g} \right) \left( \theta_i - \theta_s \right) \frac{1}{\left( 1 + 1/m \right)} \frac{\partial S_e}{\partial \psi} \left( -\alpha \right)^{n-1}$$ \hspace{1cm} (10)

where $\psi$ is the matric suction.

We now take the partial derivative of (10) with respect to time $t$ and substitute the result into (6), to obtain the following equation:

$$\frac{\partial \theta}{\partial t} = -\rho_w g \left[ \frac{2\beta}{C} \left( S_e + h \frac{\partial S_e}{\partial h} \right) + \omega + h \frac{\partial (1/R)}{\partial h} \right] \frac{\partial h}{\partial t}.$$ \hspace{1cm} (11)

Thus, the right hand side of (11) is a function of the pressure head.

Referring to a Cartesian coordinate system and substituting (11) into (9), we obtain the equation governing the transient hydromechanical process in a partially saturated porous medium as follows:

$$\frac{\partial}{\partial x} \left( k_x \frac{\partial h}{\partial x} \right) + \frac{\partial}{\partial z} \left[ k_z \left( \frac{\partial h}{\partial z} - 1 \right) \right] = -\rho_w g \left[ \frac{2\beta}{C} \left( S_e + h \frac{\partial S_e}{\partial h} \right) + \omega + h \frac{\partial (1/R)}{\partial h} \right] \frac{\partial h}{\partial t}.$$ \hspace{1cm} (12)

If the porous soil mass is considered to be rigid $(C \rightarrow \infty)$, (12) reduces to the unsaturated flow equation without coupling between infiltration and deformation in the unsaturated porous medium.

3. Analytical Solution to the Coupled Infiltration in a Semi-Infinite Region

The effective analysis of unsaturated seepage coupled with deformation in a semi-infinite unsaturated domain is based on the following assumptions: (a) the soil is homogeneous and behaves as an isotropic elastic material; (b) the soil structure is deformable, but the pore water is not compressible; (c) the volume change of the soil is due to the wetting or drying of the soil only; (d) the coefficient of permeability at full saturation is kept constant regardless of ground deformation; (e) the pore-air pressure in the soil remains constant.

Based on Darcy’s Law and conservation of the 1D fluid mass, combined with the stress-strain relationship applicable to normal strains proposed by Fredlund and Rahardjo [24], the 1D governing equation for partially saturated infiltration coupled with deformation in a semi-infinite domain, where the water compressibility is neglected, can be given as [11, 14]

$$\frac{\partial}{\partial z} \left[ k (h) \frac{\partial (h - z)}{\partial z} \right] = \left[ \frac{\partial S_e}{\partial h} + S_e \frac{(1 + \nu)}{(1 - \nu)} - \frac{1}{R} \right] \frac{\partial h}{\partial t},$$ \hspace{1cm} (13)

where $S_e$ is the unit weight of water; $n$ is porosity; $\alpha_s = 1 - K/K_s$ is the hydromechanical coupling coefficient ($0 \leq \alpha_s \leq 1$), where $K$ is the bulk modulus of the solid skeleton and $K_s$ is the bulk modulus of the soil solids; and $\nu$ is Poisson’s ratio.

The constitutive behavior of unsaturated soils that allow fluid infiltration is bound to have advanced concepts involving plasticity and dilatant plasticity in their descriptions. If such processes take place, then the alterations in the fluid transport characteristics should also be accounted for. It seems unlikely that an elastoplastic solution could be developed through analytical solutions to address these issues. The linearization of the modelling including elastic linearized model is a useful first step towards understanding the coupling effects. Ideally, developments involving both saturated and unsaturated geomaterials should take into account the irreversible deformations of the porous skeleton, which can include both elastoplasticity and creep [30, 31]. The latter processes are, however, less likely to be significant in the type of problems discussed in this paper. Elastoplasticity and poromechanical effects can be important and in order to examine such effects the problem must be formulated using a three-dimensional incremental approach with specified failure criteria and flow rules [32, 33]. The elementary elasticity approach is not perfect but will provide an opportunity to assess whether further extensions to other constitutive modelling will be warranted. The approach will also provide a bound for the behavior of the soil under rainfall influx.

The van Genuchten model [28] is complex and highly nonlinear. Hence the model is difficult to apply to analytical methods. The exponential variation of hydraulic conductivity with pore-water pressure was applied by Gardner [34]. In many analytical solutions, the generally simple exponential hydraulic conductivity function of Gardner [34] is employed [12]. The relationship between pore-water pressure and hydraulic conductivity is nonlinear. Combining this result with Kirchhoff’s transformation and observing that pore-water pressure varies linearly with hydraulic conductivity after a series of transformations allow the linearization of Richards’
equation for transient partially saturated flow in order to achieve an analytical solution to coupled infiltration and deformation in unsaturated soils. Therefore, \( k(h) = k_s e^{\alpha h} \) and \( \theta(h) = \theta_s e^{\alpha h} \) are employed to describe the hydraulic conductivity of an unsaturated porous medium and the relationship between the volumetric moisture and matric suction. \( \alpha \) is the desaturation coefficient.

Substituting the exponential forms into (8), we obtain the 1D governing equation:

\[
\frac{\partial}{\partial z} \left[ k_s e^{\alpha h} \frac{\partial (h - z)}{\partial z} \right] = P \cdot e^{\alpha h} \frac{\partial h}{\partial t}
\]

(14)
in which \( P = \theta_s \alpha + \gamma_w \alpha / (1 + \nu)H \).

If coupling is not considered, (14) becomes

\[
\frac{\partial}{\partial z} \left[ k_s e^{\alpha h} \frac{\partial (h - z)}{\partial z} \right] = \alpha \theta_s e^{\alpha h} \frac{\partial h}{\partial t}.
\]

(15)

The infiltration equations for a semi-infinite domain can be solved with a specified surface flux \( q(t) \). The initial water pressure is described by \( h(z,0) = h_r \), and a boundary condition is \( h(z,t)|_{z \to \infty} = h_r \).

Dimensional variables of \( Z = \alpha z \), \( T = \alpha^2 k_s t/P \), and \( K(Z, T) = e^{\alpha h} \cdot e^{T/4 - T/2} \) are employed. Based on (14) and (15), we obtain

\[
\frac{\partial^2 K}{\partial Z^2} = \frac{\partial K}{\partial T}
\]

(16)

\[
K(Z,0) = e^{\alpha h} \cdot e^{-Z/2} = f(Z)
\]

(17a)

\[
-K(Z,T)|_{Z \to \infty} = e^{\alpha h} \cdot e^{T/4}.
\]

(17c)

The analytical solution to (16) is given by [35]

\[
K(Z,T) = \int_{\beta=0}^{\infty} M(\beta, X) \cdot e^{\beta Z} \left[ \bar{F}(\beta) + \int_{t'=0}^{T} e^{\beta t'} A(\beta, t') dt' \right] d\beta
\]

(18)

in which

\[
\bar{F}(\beta) = \int_{z'=0}^{\infty} M(\beta, z') f(z') dz'
\]

(19a)

\[A(\beta, t') = M(\beta, z)|_{z=0} g(t')\]

(19b)

\[M(\beta, Z) = \sqrt{\frac{2 \beta \cos \beta Z + 0.5 \sin \beta Z}{\sqrt{\beta^2 + 0.25}}},\]

(19c)

where \( f(z') \) and \( g(t') \) represent the initial and surface boundary conditions, respectively; \( M(\beta, Z) \) is the kernel of the equation.

### 4. Computational Modelling of Water Infiltration in a Semi-Infinite Unsaturated Region

Figure 1 shows a two-dimensional semi-infinite region where there is rainfall influx over a symmetric finite region. The 2D semi-infinite region is represented by a symmetric model of finite extent shown in Figure 2; in this model the region \( A \to 0 \) represents the width of the rainfall region with different rainfall intensities and the region \( A \to B \) is an impervious boundary. The two sides 0D and BC are controlled by a zero water flow rate in the horizontal direction. The lower boundary CD is controlled by \( \nabla h = 0 \), and the entire domain is imposed with zero normal displacement boundary conditions on the base of the region and along the vertical sides. An external stress (200 kPa) is vertically applied at the surface boundary 0B in Figure 2.

The domain shown in Figure 2 is discretized into 478 triangular finite elements using an adaptive mesh feature available in COMSOL Multiphysics. For the computational simulations, the model dimensions are set to \( w \) in length by \( d \) in height, and the width of the rainfall region is \( a \), as shown in Figure 2. In Figure 2, \( u \) denotes displacement, and \( \tau \) is the shear stress. Since the COMSOL mesh generation is adaptive, stabilized finite element methods are necessary to develop numerical results to examine the two-dimensional symmetric coupled seepage and deformation problem related to a partially saturated porous medium [36]. In the multiphysics code COMSOL the mesh refinement is automatically adjusted when the wetting front caused by water infiltration moves into the partially saturated porous medium [36].

In the model, the two lateral boundaries shown in Figure 2 are subjected to fluid flow boundary condition corresponding to a zero flux and constant head boundary condition as follows:

\[
\left[ k(h) \frac{\partial h}{\partial x} \right]_{x=0} = 0 \quad 0 < z < d
\]

(20a)

\[k h|_{x=w} = h_r \quad 0 < z < d.\]

(20b)
At the plane of symmetry the fluid flow boundary conditions consist of the lower and upper boundaries, which can be given as

\[ h(x, 0, 0) = h_r \]  
\[ \left[ k(h) \frac{\partial h}{\partial z} - k(h) \right]_{z=d} = q(x, t) \quad 0 < x < w, \]  

where \( h_r \) is the pore-water pressure head that corresponds to the residual volumetric water content and \( q \) is a function which varies with position \( x \) and time \( t \). In addition, the boundary conditions corresponding to traction and displacements applicable to the porous medium are

\[ u(x, z, t)|_{x=0, z \in [0, d]} = 0 \]  
\[ u(x, z, t)|_{x=w, z \in [0, d]} = 0 \]  
\[ u(x, z, t)|_{z=d, x \in [0, w]} = 0. \]

Initially the medium is assumed to be undeformed; the only free boundary is at the ground surface, and the other three boundaries are controlled by a fixed constraint. The initial pressure head distribution (i.e., the initial condition) is written as

\[ h(x, z, 0) = h_i(x, z), \]

where the initial pressure head \( h_i \) is assumed to be constant.

The model of a semi-infinite domain is depicted in Figure 1. The top flux is specified when \( z \) approaches infinity, and the pore-water pressure head is constant \( (h_r) \). The initial pore-water pressure is \( h_r \). The parameters used are listed in Table 1. Figure 3 compares the 1D numerical and analytical solutions for the finite unsaturated domain considering the coupling between infiltration and deformation. The 1D numerical solution is obtained by substituting \( k(h) = k_s e^{\alpha h} \) and \( \theta(h) = \theta_s e^{\alpha h} \) into (8), which, when combined with (4), can be resolved in COMSOL. The difference observed between the analytical and numerical solutions to the 1D infiltration with coupled effects is small.

The parameters for the partially saturated porous medium used in this study are listed in Tables 2 and 3. Table 2 gives the parameters used in the calculations involving the partially saturated porous medium. The basic soil properties are presented in Table 3.

Figure 4 shows the variation of pressure head with duration of rainfall under the conditions of coupled water infiltration and deformation in an unsaturated porous medium. In Figure 4, \( w = 10 \) m, \( d = 10 \) m, and \( a = 1 \) m. The changes in the pressure head occur only in the shallow layers of the partially saturated porous medium. Noticeable changes in the pore-water pressure head occur close to the water infiltration region.

As shown in Figures 5 and 6, the pressure head and displacement at the upper boundary 0B change with the duration of rainfall. The parameters are the same as in Figure 4. At the start of water infiltration, the pressure head and the ground settlement is uniform. As rainfall continues, water infiltrates into the deeper regions of the unsaturated porous medium and uneven displacement in the shallow layers of the unsaturated soil begins to occur. After \( t = 10 \) h, we observe different displacements at the ground surface of the unsaturated region, as shown in Figure 6. With continued rainfall, differential settlements occur at the surface of the partially saturated porous medium.

Figure 7 shows how changing the infiltration width 0B affects the pore-water pressure head at the boundary 0B shown in Figure 2. In Figure 7 the infiltration width is altered \((w = 5, 10, 15, 20, 25, 30 \) m\) but the other dimensions...
Table 1: The parameters used in analytical solutions to 1D unsaturated infiltration.

<table>
<thead>
<tr>
<th>Saturated volumetric water content ($\theta_s$)</th>
<th>Coefficient of permeability at saturation ($k_s$)</th>
<th>Rainfall intensity ($q$)</th>
<th>Desaturation coefficient ($\alpha$)</th>
<th>Elastic modulus with respect to suction change</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.4</td>
<td>$1 \times 10^{-5}$ m/s</td>
<td>$10^{-5}$ m/s</td>
<td>0.01 cm$^{-1}$</td>
<td>10 MPa</td>
</tr>
</tbody>
</table>

Table 2: The parameters used in the calculations involving SWCC.

<table>
<thead>
<tr>
<th>Saturated volumetric water content ($\theta_s$)</th>
<th>Residual volumetric water content ($\theta_r$)</th>
<th>Initial porosity ($n_0$)</th>
<th>$a$</th>
<th>$n$</th>
<th>$m$</th>
<th>Rainfall intensity ($q$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.4</td>
<td>0.01</td>
<td>0.35</td>
<td>0.02</td>
<td>2</td>
<td>0.5</td>
<td>$0.8 \times 10^{-5}$ m/s</td>
</tr>
</tbody>
</table>

Table 3: The basic soil properties of an unsaturated porous medium.

<table>
<thead>
<tr>
<th>Elastic modulus</th>
<th>Poisson’s ratio</th>
<th>Unit weight of soil mass</th>
<th>Unit weight of water</th>
</tr>
</thead>
<tbody>
<tr>
<td>10 MPa</td>
<td>0.3</td>
<td>18.0 kN/m$^3$</td>
<td>10.0 kN/m$^3$</td>
</tr>
</tbody>
</table>

Figure 4: The pressure head profile variation with infiltration time: (a) $t = 5$ h; (b) $t = 10$ h; (c) $t = 30$ h; (d) $t = 50$ h.
remain unchanged. The boundary effect caused by different infiltration distances varies with the infiltration time. As shown in Figure 7, compared with $w = 30\, m$, the difference in the pressure head is 13% for an infiltration width of 10 m and 20% for a 5 m width at $t = 10\, h$. However, the difference in the pressure head is only 2% when the width of the infiltration area is 20 m. In comparison with the 30 m width, the difference in the pressure head at $t = 40\, h$ is 17% for an infiltration width of 10 m and 25% for a 5 m width. At $t = 200\, h$ the difference in the pore-water pressure head for a 5 m width is 13% and 25% for an infiltration width of 3 m. However, the difference in the pressure head is only 3% when the width of the infiltration area is 20 m at $t = 50, 200\, h$. Figure 7 indicates that the boundary dimensions in the model play a significant role in the pore-water pressure head movement in the 2D semi-finite domain. It is recommended that in the computational model a width of 20 m (dimensionless length of 20) be used to analyze the 2D semi-infinite coupling problem.

Figure 8 shows the settlement at point 0 on the ground surface with the same rainfall duration. There is obvious settlement at the early stage of rainfall infiltration and the settlement rate gradually slows over time. When $k_s = 1 \times 10^{-3}\, m/s$ representing for a typical sandy soil, after $t = 9\, h$, the settlement of the ground surface appears to have stabilized.
When $k_s = 1 \times 10^{-5}$ m/s for a typical silty soil, after $t = 80$ h, the settlement of the ground surface appears to have stabilized. When $k_s = 1 \times 10^{-6}$ m/s (clay), after $t = 550$ h, the settlement of the ground surface appears to have stabilized. With a decrease of $k_s$, the time required for stabilization becomes larger. When $k_s \geq 1 \times 10^{-3}$ m/s, the settlement at point 0 follows the same trend as $k_s = 1 \times 10^{-4}$ m/s. However, the settlement for $k_s \geq 1 \times 10^{-3}$ m/s is much less that for $k_s < 1 \times 10^{-3}$ m/s.

5. Conclusions

A numerical method is employed to examine coupled infiltration and deformation in a semi-infinite unsaturated porous medium. The hydromechanical coupling in the semi-infinite unsaturated porous medium region could be transferred to a symmetric finite domain problem by setting up an effective investigation area. The results indicate that the numerical solutions could effectively resolve the semi-infinite
unsaturated infiltration when the coupled problem is examined. The different widths and depths considered are of significance when examining the hydromechanical coupling during rainfall infiltration into a semi-infinite domain. The results show that the coupled effect is significantly affected by the boundary size, and the coupling is related to the duration of infiltration. The rate of change in the pore-water pressure head and settlement in an axisymmetric semi-infinite unsaturated porous medium gradually slows over the infiltration duration. This offers a first step necessary to develop more complex models describing irreversible processes of unsaturated porous media.

Conflicts of Interest

The authors declare that there are no conflicts of interest regarding the publication of this paper.

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